

## Systematic Dependence of $R_4$ on $N_p N_n$ Product for Light and Medium Mass Nuclei

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### Introduction

The study of nuclear structure with  $N$ ,  $Z$ ,  $N_B$  and  $N_p N_n$  provide a deep understanding of nuclear interactions involved. Various studies had been carried to study the systematic dependence of various nuclear properties on  $N_p N_n$ . Casten et al [1] presented a review on the evaluation of nuclear structure on the bases of  $N_p N_n$ ; this phenomenon has been called the  $N_p N_n$  scheme. Gupta et al [2] presented a systematic dependence of the  $\gamma - g$   $B(E2)$  ratios on the  $N_p N_n$  in different parts of the major shell space  $Z = 50 - 82$ ,  $N < 82$  and  $N = 82 - 126$  and demonstrated that the interband  $B(E2)$  ratios were smooth functions of  $N_p N_n$ . Further, Gupta et al [3] pointed out the limitations of the F-spin and  $N_p N_n$  scheme in reproducing the overall  $E_{2g^+}$  systematic in the major shell space  $Z = 50 - 82$ ,  $N = 82 - 126$  into four quadrants and predicted the position of the  $\beta$ -stability line on the  $N$ - $Z$  chart to explain the existence of isotonic (isotopic) multiplets.

Recently, most work related to the  $N_p N_n$  scheme mainly concentrate on P- factor [ $=N_p N_n / (N_p + N_n)$ ] [ref. 1], and  $\beta_2$  [ref. 4, 5, 6], systematic law of  $E_{2g^+}$  for heavy nuclei [7], excitation energy of first  $2^+$  state of all even-even nuclei [8, 9] and odd-even staggering [10]. The systematic dependence of  $R_4$  on  $N_p N_n$  in major shell- space has not been studied sufficiently. The present study of dependence of  $R_4$  on  $N_p N_n$  is interesting to investigate new insight. Recent experimental data have been taken ref. [11]. In the present work, we focus on the systematic dependence of  $R_4$  on  $N_p N_n$  in even - even nuclei for  $Z = 50 - 82$ ,  $N = 82 - 126$  region, by dividing the whole space into four quadrants.

### Result and Discussion

In this work, we adopt a grouping based on the valance particle and hole pairs consideration [12]. Thus the  $Z = 50 - 82$ ,  $N = 82 - 126$ , major shell space is partitioned into four quadrants. The quadrant- I has  $Z = 50 - 66$ ,  $N = 82 - 104$  (p-p), in quadrant- II,  $Z = 66 - 82$ ,  $N = 82 - 104$  (p-h), quadrant - III,  $Z = 66 - 82$ ,  $N = 104 - 126$  (h-h), and quadrant- IV,  $Z = 50 - 66$ ,  $N = 104 - 126$ , (p-h), where, p = valance particle, proton or neutron and h = hole. The quadrant-IV does not has any data point, thus it is empty.

### Dependence of $R_4$

In the IBM-1[13], a useful measure of collectivity is  $R_4$  for rotational nuclei this ratio  $R_4 = 3.33$  and the region  $3 \leq R_4 \leq 3.33$  is called the rotational region and for the vibrational nuclei  $R_4 = 2$  and the region  $2 \leq R_4 \leq 2.4$  is called the vibrational region while the region  $2.4 \leq R_4 \leq 3$  is called the transition region. The transition region contains nuclei with structure intermediate between vibration and rotational. For simplicity, we divide the whole data of ratio  $R_4$  of as function of  $N_p N_n$  into four quadrants as above. First we consider quadrant-I, Fig. 1 shows the ratio  $R_4$  for even-even nuclei as function of  $N_p N_n$ . In this fig. for Xe -Dy nuclei with  $N=82-104$  (p-p region). The  $R_4 < 1.5$  for  $N_p N_n = 0$  (i.e. magic nuclei) and  $R_4$  varies smoothly with  $N_p N_n \leq 30$ . These nuclei have a shape change vibrational to transitional because  $1.9 \leq R_4 \leq 3$  and finally  $R_4$  remains unchanged when  $N_p N_n \geq 30$  these nuclei are called rational nuclei because  $R_4 = 3.33$ .

From Fig 2 for Dy -Pt nuclei with  $N = 82-104$  (h-p boson), i.e. quadrant-II here most of the data

points lie on a smooth curve that rises with increasing NpNn product. In this region most of nuclei have rotational nature because  $3 \leq R_4 \leq 3.33$  except few data points of Er, Dy and Pt.

From fig. 3, for Yb -Hg nuclei with N = 104-126 (h-h boson region) i.e. quadrant III, in this region the dependence of  $R_4$  on NpNn is again smooth for Yb -Os these nuclei are in rotational region because  $3 \leq R_4 \leq 3.33$  and the data point of Pt and Hg indicates a new signature of vibrational nature because  $R_4 = 2.5$ .

This systematic dependence of  $R_4$  for ground state in all three regions on the NpNn indicates its close relationship to the shape deformation of the nuclear core.

**Conclusions**

The present study reflects that the ratio  $R_4$  depends upon NpNn. The variation of  $R_4$  shows the average dependence on NpNn for shape deformation in all three quadrants except Pt and Hg isotopes in quadrant-III. Here a complexity of nuclear structure is exhibited between NpNn and  $R_4$ .

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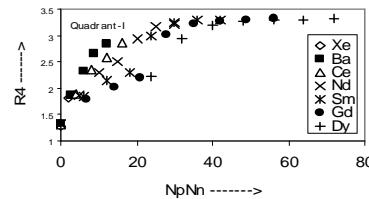


Fig. 1. Experimental data for  $R_4$  for even – even nuclei as a function of NpNn product for Z = 50 - 66, N = 82 - 104 i.e., quadrant – I.

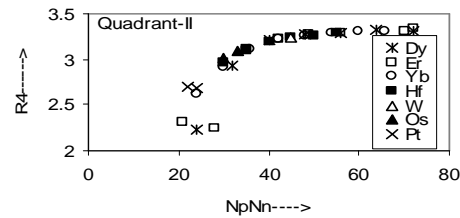


Fig.2. Same as fig.1 for Z= 66 - 82, N = 82 - 104, i.e, quadrant –II.

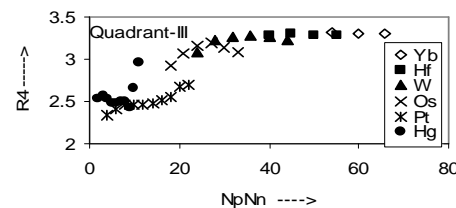


Fig. 3. Same as fig. 1 for Z = 66 - 82, N = 104 - 126, i.e., quadrant –III.