# Study of triaxial rotor model and correlation of deformation parameter $\beta$ and $\gamma_0$ in A=120-200 mass region

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## Introduction

One basic property of the nucleus is its geometric shape. To describe the nuclear structure we turns into the two well known parameter asymmetric parameter  $(\gamma_0)$  and deformation parameter  $(\beta)$ , which can be obtained by one of the simplest model for the nuclear collectivity, the rigid triaxial rotor model of Davydov and Fillipov [1]. The triaxiality is the key property of the ground state, as well as the excited states of nucleus. To understand the triaxiality features the gamma degree of freedom plays an important role in the description of these nuclei. The  $\beta$  and  $\gamma_0$  of the collective model are the basic descriptor of the nuclear equilibrium shape and structure. The anharmonic vibrator (AHV) expression [2],

$$E(J) = aJ(J+1) + bJ \tag{1}$$

However, this two-parameter expression yields a linear relation of  $R_{6/2}(=E(6_g^+)/E(2_g^+)$  with  $R_{4/2}$  but the experimental data deviate from this, showing the need of a 3rd term viz

$$E(J) = aJ(J+1) + bJ + cJ^2(J+1)$$
 (2)

which gives a non-linear relationship of  $R_{6/2}$ with  $R_{4/2}$ . The variation of  $\text{ROTE}/E(2_g^+)$ with asymmetric parameter  $(\gamma_0)$  in all the four quadrants is already studied by [3]. We extended this work here, to understand the shape phase transition when mixing angle  $(\tau)$  is added to the  $\gamma_0$ . We also study the variation of  $\gamma_0$  and  $\beta$ .

## Calculation and Results

According to the J. L Wood [4] the Hamiltonian for the model is

$$H = A_1 \tilde{I}_1^2 + A_2 \tilde{I}_2^2 + A_3 \tilde{I}_3^2 \tag{3}$$

The Hamiltonian can be written as

$$H = A\tilde{I}^2 + F\tilde{I}_3^2 + G(\tilde{I}_+^2 + \tilde{I}_-^2)$$
(4)

where

$$A = \frac{1}{2}(A_1 + A_2), \ F = A_3 - A, \ G = \frac{1}{4}(A_1 - A_2)$$
(5)

and for I = 2

$$\begin{pmatrix} 6A & 4\sqrt{3}G\\ 4\sqrt{3}G & 6A+4F \end{pmatrix} \tag{6}$$

therefore

t

$$\tan \tau = \frac{\sqrt{F^2 + 12G^2} - F}{2\sqrt{3G}} \tag{7}$$

Therefore moment of inertia is directly equal to the

$$J_k = 4B\beta^2 \sin^2(\gamma_0 - k\frac{2\pi}{3}), k = 1, 2, 3... \quad (8)$$

where B is mass parameter, therefore [4]

$$\tau = -\frac{1}{2}\cos^{-1}\left(\frac{\cos 4\gamma_0 + 2\cos 2\gamma_0}{\sqrt{9 - 8\sin^2 3\gamma_0}}\right) \quad (9)$$

From (9) we calculate the value of mixing angle. To calculate the  $\gamma_0$  there are various methods, but the  $\gamma_0$  derived from  $R_{\gamma}(=E(2^+_{\gamma})/E(2^+_g))$  is more relevant [5]

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$$\gamma_0 = \frac{1}{3} \sin^{-1} \left[ \frac{9}{8} \left( 1 - \left( \frac{R_\gamma - 1}{R_\gamma + 1} \right)^2 \right) \right]^{1/2}$$
(10)

Here  $\beta$  is calculated from [6]

$$\beta = \beta_G \left( \frac{9 - 8\sqrt{81 - 72\sin^2 3\gamma_0}}{4\sin^2(3\gamma_0)} \right) \quad (11)$$

where

$$\beta_G^2 \cong \frac{1224}{E_{2^+}A^{7/3}}.$$
 (12)

For above calculation the experimental data is taken from [7].



FIG. 1: Correlation of  $\gamma_0$  with  $\beta$  for N=68-126 region nuclei.

By counting of particle (p) and hole (h) boson partitions the major shell space of

Z=50-82, N=66-82, N = 82-126 are divided in four quadrants (Q) viz. quadrant I and III for p-p and h-h bosons, and II, IV for p-h and h-p bosons respectively [8]. In Fig. 1 we see the variation of asymmetric parameter  $\gamma_0$ with deformation parameter  $\beta$  and these two parameter shows excellent agreement with each other. From this we observe that  $\beta$  is constant with  $\gamma_0$  for soft nuclei as compared to the well deformed nuclei.

## Conclusion

Above study reveals some important description for A=120-200 mass region nuclei with the help of Davydov Filippov model. The present work provides new insight into the triaxial rotor model, which is widely used as a basic description of nuclear collectivity. The correlation of asymmetric parameter  $\gamma_0$ and deformation parameter  $\beta$  gives extremely well agreement with each other.

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