

## Prediction for the ground state energy of ${}_{\Lambda\Lambda}^{14}\text{C}$ in the cluster model

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### Introduction

In the past, Shoeb and co-workers [1, 2], in the variational Monte Carlo (VMC) framework, have successfully explained the energies of the ground and excited states of  ${}^9_{\Lambda}\text{Be}$  and  ${}^{10}_{\Lambda\Lambda}\text{Be}$  in the  $\alpha$  cluster model. The success is attributed to the proposed phenomenological dispersive [1] three-body  $\Lambda\alpha\alpha$  cluster force inspired from the microscopic analysis of Bodmer *et al.* [3]. Recently, we have analyzed the ground state energies of  ${}^{12}\text{C}$  and  ${}^{13}\text{C}$  in the  $\alpha\alpha\alpha$  and  $\Lambda\alpha\alpha\alpha$  models, respectively, alongwith the other  $\alpha$  cluster strange and multistrange hypernuclei [4]. The phenomenological dispersive [1] three-body  $\Lambda\alpha\alpha$  and  $\alpha\alpha\alpha$  forces along with the two-body  $\alpha\alpha$  and  $\Lambda\alpha$  potentials explained the ground state energies. The VMC calculation of the energy [5] of the excited  $2^+$  state of  ${}^{12}\text{C}$  in the  $\alpha$  cluster model is found to be consistent with the observed value. Looking at the success of VMC calculation in explaining the ground and excited state properties of  ${}^{12}\text{C}$  and ground state of  ${}^{13}_{\Lambda}\text{C}$ , we have extended it to predict the energy of  ${}^{14}_{\Lambda\Lambda}\text{C}$  in view of the ongoing programme to produce  $S = -2$  hypernuclei at J-PARC.

### Hamiltonian, wavefunction and energy calculation

The Hamiltonian for the five-body system

${}^{14}_{\Lambda\Lambda}\text{C}$  in the  $\Lambda\Lambda\alpha\alpha\alpha$  model takes the form:

$$\begin{aligned}
 H_{\Lambda\Lambda} = & \sum_{i=1}^2 K_{\Lambda}(i) + \sum_{i=3}^5 K_{\alpha}(i) + V_{\Lambda\Lambda}(r_{12}) \\
 & + \sum_{i=1}^2 \sum_{j=3}^5 V_{\Lambda\alpha}(r_{ij}) \\
 & + \sum_{i<j, i=3,4, j=4,5} V_{\alpha\alpha}(r_{ij}) \\
 & + \sum_{i=1}^2 \sum_{j<k, j=3}^5 V_{\Lambda\alpha\alpha}(r_{ij}, r_{ik}) \\
 & + V_{\alpha\alpha\alpha}(r_{34}, r_{35}, r_{45}), \quad (1)
 \end{aligned}$$

where indices (1,2), (3,4,5) label  $\Lambda$ s, and  $\alpha$  particles, respectively,  $K_Y$  is the kinetic energy operator for the particle  $Y (= \Lambda, \alpha)$ ,  $V_{hh}$  denotes the potential for the pair of particles  $hh$  ( $= \Lambda\Lambda, \Lambda\alpha, \alpha\alpha$ ),  $r_{kl}$  is the inter-particle separation for the pair having indices  $k$  and  $l$ . Here we have considered that all the  $hh$  pairs in a given system interact in the relative angular momentum  $s$ -state and  $V_{\Lambda\alpha\alpha}$  is the phenomenological dispersive three-body  $\Lambda\alpha\alpha$  potential and has a simple form

$$V_{\Lambda\alpha\alpha} = W_0 f(r_{\Lambda\alpha_1}) f(r_{\Lambda\alpha_2}), \quad (2)$$

where the strength  $W_0 > 0$  gives repulsion. The radial factor  $f(r)$  is of Yukawa form:  $f(r) = \exp(-ar)/ar$ , where  $a$  is the range parameter. The  $V_{\Lambda\alpha\alpha}$  ( $W_0 = 17.0$  MeV and  $a = 0.5 \text{ fm}^{-1}$  for the case of two-body  $\Lambda\alpha$  potential) is constrained to fit the  $B_{\Lambda}$ ,  $\Lambda$ -binding energy [1] of the  ${}^9_{\Lambda}\text{Be}$  in the three-body  $\Lambda\alpha\alpha$  model. The inclusion of the dispersive  $\Lambda\alpha\alpha$  force [1, 2] gives a good account of the ground and excited states of  ${}^9_{\Lambda}\text{Be}$  and  ${}^{10}_{\Lambda\Lambda}\text{Be}$ , as stated in the preceding section. The expressions for  $V_{\alpha\alpha\alpha}$ , three-body  $\alpha\alpha\alpha$  Gaussian shape potential and for the two-body  $V_{\alpha\alpha}$  potential in the

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relative angular momentum  $s$ -state have been given in our other paper [5].

The phenomenological  $s$ -state  $\Lambda\alpha$  potential of Isle form, the  $\Lambda\Lambda$  potential in the singlet state of the three-range Gaussian form adjusted to fit the ground state binding energy (7.26 MeV) of  ${}^6_{\Lambda\Lambda}\text{He}$  and  $\alpha\alpha$  potential of Ali and Bodmer [6] are taken from our earlier work [1, 2].

The trial wavefunction for the system  ${}^{14}_{\Lambda\Lambda}\text{C}$  is the product of two-body correlation functions  $f_{hh}$  for the pair of particles,  $hh$  and the appropriate spin function  $\zeta_0^0$  and has the form:

$$\Psi_{\Lambda\Lambda} = \left[ \prod_{i=1}^2 \prod_{j=3}^5 f_{\Lambda\alpha}(r_{ij}) \right] f_{\Lambda\Lambda}(r_{12}) \times \left[ \prod_{i < j, i=3,4, j=4,5} f_{\alpha\alpha}(r_{ij}) \right] \zeta_0^0. \quad (3)$$

The correlation functions  $f_{hh}(r)$  have the asymptotic behavior  $f_{hh}(r) \sim r^{-\nu_{hh}} \exp(-\kappa_{hh}r)$  and are obtained from the solution of the Schrödinger-type equation.

The total energy  $E_B$  for the system  ${}^{14}_{\Lambda\Lambda}\text{C}$  in the cluster model for the trial wavefunction Eq. (3) is evaluated using the following relation:

$$E_B({}^{14}_{\Lambda\Lambda}\text{C}) = \frac{\langle \Psi_{\Lambda\Lambda} | H_{\Lambda\Lambda} | \Psi_{\Lambda\Lambda} \rangle}{\langle \Psi_{\Lambda\Lambda} | \Psi_{\Lambda\Lambda} \rangle}. \quad (4)$$

The VMC estimates of the energy were made for 100 000 points. A distance of 20 fm was chosen for the integration and the energy was optimized with respect to variational parameters using standard code.

## Results and Discussion

The ground state energy of  ${}^{12}\text{C}$  which we have calculated earlier [4] in the VMC framework in the  $\alpha$  cluster model is 7.17 MeV, quite consistent with the experimental value 7.26

MeV. The results of variational calculation for  ${}^{14}_{\Lambda\Lambda}\text{C}$  are as under;

Energy  $-E_B = 31.29 \pm 0.0002$  MeV with break up: total kinetic energy = 30.58 MeV, total potential energy = -61.74 MeV, three-body dispersive energy  $\langle V_{\Lambda\alpha\alpha} \rangle = 8.72$  MeV, three-body cluster energy  $\langle V_{\alpha\alpha\alpha} \rangle = -8.85$  MeV. The root mean square radius between  $\alpha\alpha$ ,  $\mathbf{R}_{\alpha\alpha} = 3.48$  fm; between  $\Lambda\Lambda$ ,  $\mathbf{R}_{\Lambda\Lambda} = 3.22$  fm; between c.m. of 3-alpha and  $\Lambda$ ,  $\mathbf{R}_{(3\alpha)\Lambda} = 2.36$  fm; between  $\Lambda$  and  $\alpha$ ,  $\mathbf{R}_{\Lambda\alpha} = 3.09$  fm; between c.m. of 3-alpha and  $\alpha$ ,  $\mathbf{R}_{(3\alpha)\alpha} = 2.01$  fm; between c.m. of 3-alpha and  $\Lambda\Lambda$ ,  $\mathbf{R}_{(3\alpha)(\Lambda\Lambda)} = 1.72$  fm.

The predicted  $B_{\Lambda\Lambda}$  ( $= E_B({}^{12}\text{C}) - E_B({}^{14}_{\Lambda\Lambda}\text{C})$ ) for  ${}^{14}_{\Lambda\Lambda}\text{C}$  turns out to be 24.12 MeV.

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