

Establishing the Island of stability for superheavy elements: A new approach based on reaction data

Niyti and Raj K. Gupta

Physics Department, Panjab University, Chandigarh-160014, INDIA

Introduction

A fascinating challenge in the study of superheavy nuclei is the quest for island of stability where the magic numbers next to $Z=82$ and $N=126$ may be located. The earliest calculation of 1960's established that the center of island of stability for superheavy elements (SHEs) lies at $Z=114$, $N=184$. Recently, the microscopic mean field models, applied in the region of SHEs, predict $Z=120$ and $N=172$ or 184 as the next magic numbers [1, 2]. Also, $Z=126$, $N=184$ have been used/predicted as magic numbers [3]. The goal of present work is to identify which one of these three Z ($=114, 120$ or 126) is a better shell closure with $N=184$. To this end we use the hot fusion reaction $^{48}\text{Ca} + ^{238}\text{U} \rightarrow ^{286}112^*$, with measured evaporation residue (ER), fusion-fission (ff) and quasi-fission (qf) cross-sections, as a tool and analyse it on the basis of the dynamical cluster-decay model (DCM) of Gupta and collaborators (see, e.g., [4] and earlier references therein) where the effects of deformation up to hexadecapole deformation β_4 and "compact" orientations θ_c ($\theta_c=72^\circ$ for ^{238}U) are included. The DCM gives a good description of the ER (the light-particle emission), ff and qf (equivalently, capture) decay channels, and their excitation functions within a single parameter description, the neck length parameter ΔR .

The Methodology

The compound nucleus (CN) decay cross-section in DCM, in terms of partial waves, is

$$\sigma = \frac{\pi}{k^2} \sum_{\ell=\ell_{min}}^{\ell_{max}} (2\ell+1)P_0P; \quad k = \sqrt{\frac{2\mu E_{cm}}{\hbar^2}} \quad (1)$$

with reduced mass $\mu=[A_1A_2/(A_1+A_2)]m$, and center of mass energy E_{cm} . P is the WKB penetrability, with first turning point

$R_a = R_1(\alpha_1, T) + R_2(\alpha_2, T) + \Delta R(T)$. P_0 is the solution of stationary Schrödinger equation in mass asymmetry coordinate $\eta = (A_1 - A_2)/(A_1 + A_2)$, i.e., $P_0(A_i) \propto |\psi(\eta(A_i))|^2$, $i=1,2$, with mass fragmentation potential

$$V_R(\eta, T) = \sum_{i=1}^2 [V_{LDM}(A_i, Z_i, T)] + \sum_{i=1}^2 [\delta U_i] \exp(-\frac{T^2}{T_0^2}) + E_c(T) + V_P(T) + V_t(T). \quad (2)$$

Here, $V_{LDM}(T)$ is T-dependent liquid drop energy [5], with its constants at $T=0$ re-fitted to give the experimental binding energies B , defined within the Strutinsky renormalization procedure as $B = V_{LDM}(T=0) + \delta U$, with the shell corrections δU calculated in the "empirical method" of Myers and Swiatecki [3]. The magic numbers for superheavy region in the "empirical method" [3] are taken as $Z=126$, $N=184$. Evidently, the constants of $V_{LDM}(T=0)$ need to be refitted to give the experimental B , for the magicity at $Z=126$ changed to that at $Z=120$ or 114 , respectively. Then, P and P_0 are obtained for the three sets of magic numbers, and cross-sections calculated for the three decay processes of ER, ff and qf. Then, the excitation functions are fitted within a single parameter ΔR .

Calculations and Results

First of all, in Fig. 1, only the evaporation residue σ_{ER} is fitted for any one set of magic numbers and, using the parameter ΔR so obtained, the corresponding σ_{ER} 's are calculated for other two sets of magic numbers. Clearly, the σ_{ER} always remains the largest for magic set $Z=126$, $N=184$, independent of E^* , and the lowest for $Z=114$, $N=184$. However, in Fig. 2, when the fitting procedure is carried out simultaneously for all the three processes

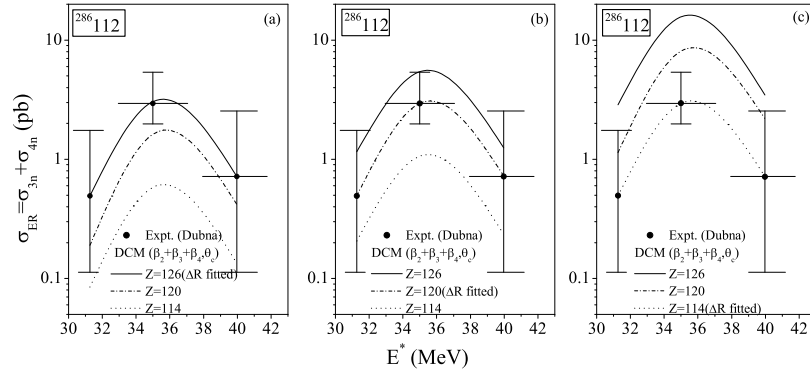


FIG. 1: ER as a function of CN E^* for $^{48}\text{Ca}+^{238}\text{U}\rightarrow^{286}112^*$ reaction, calculated for the parameter of DCM fitted to data, respectively, for $Z=126$, 120 or 114 , $N=184$ in the panels (a),(b) and (c).

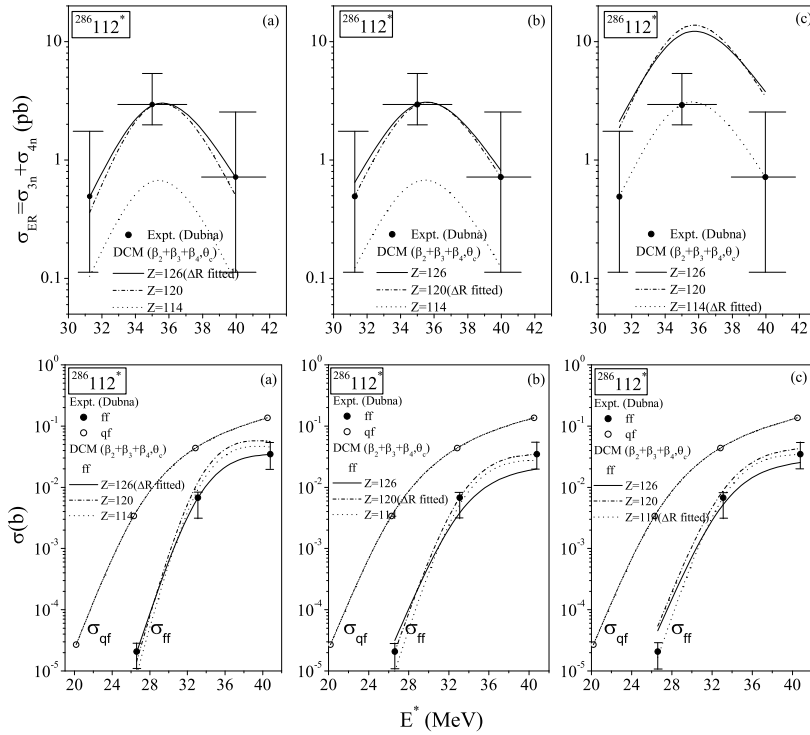


FIG. 2: Same as for Fig. 1, but now for ER, ff and qf. The qf is independent of magic shells.

of ER, ff and qf, the cross-sections are the largest and nearly indistinguishable for $Z=120$ and 126 , $N=184$. This preliminary result suggests that $Z=120$ or 126 with $N=184$ are the equally strong magic shells (largest shell corrections), and $Z=114$, $N=184$ are the weakest magic shells.

References

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