

## Decay of $^{196}\text{Pt}^*$ formed in inverse kinematic $^{132}\text{Sn}+^{64}\text{Ni}$ reaction using the dynamical cluster-decay model

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### Introduction

Evaporation residue (ER) and fission cross-sections are measured in  $^{132}\text{Sn}+^{64}\text{Ni}$  reaction at the near and sub-barrier energies [1]. At sub-barrier energies only the ER cross-sections are measured because the fission decay of the compound nucleus (CN) is negligible. In this experiment, quasi-fission (qf) component could not be separated from fusion-fission because the reaction was in inverse kinematics and the beam intensity was several orders of magnitude lower than that of stable beams (the heaviest stable Sn-nucleus is  $^{124}\text{Sn}$ ). A coupled-channel calculation including neutron transfer and inelastic excitation of the projectile and target (CCFULL), and a density-constrained TDHF calculation reproduce the measured excitation functions (the sum of ER and fission cross-sections as a function of centre-of-mass energy  $E_{c.m.}$ ) very well at sub-barrier energies though there are deviations (under- and over-estimation) at the above and near Coulomb barrier energies, respectively.

In this paper, we have used the dynamical cluster-decay model (DCM) of Gupta and collaborators [2] to estimate for the first time the CN decay cross-section, as well as the possible non-compound qf content, in this reaction with inverse kinematics. In DCM, all decay products are calculated as emissions of preformed clusters through the interaction barriers, treating the ER and fusion-fission decays on equal footings, and then allowing the incoming target and projectile nuclei to keep their identity (pre-

formation probability  $P_0=1$ ) for the qf process.

### The Model

The DCM uses the collective co-ordinates of mass asymmetry  $\eta = (A_1 - A_2)/(A_1 + A_2)$  and relative separation  $R$ , which allows to define the CN decay cross-section, in terms of the partial waves, as

$$\sigma = \frac{\pi}{k^2} \sum_{l=0}^{l_{max}} (2l+1) P_0 P; \quad k = \sqrt{\frac{2\mu E_{c.m.}}{\hbar^2}}, \quad (1)$$

with  $\mu$  as the reduced mass and  $l_{max}$ , the maximum angular momentum, fixed for the vanishing of the fusion barrier of the incoming channel, or else for the light particle cross-section  $\sigma_{ER}(\ell) \rightarrow 0$ .  $P_0$ , the pre-formation probability, is the solution of stationary Schrödinger equation in mass asymmetry coordinate  $\eta$  and  $P$  is the WKB penetrability of preformed fragments or clusters in  $R$ -motion. The only parameter of the model is the temperature dependent neck length parameter  $\Delta R(T)$ , defining the first turning point  $R_a = R_1(\alpha_1, T) + R_2(\alpha_2, T) + \Delta R(T)$  for the penetration of preformed fragment. The decay fragments are considered to be deformed with quadrupole deformation  $\beta_2$  alone, having "optimal" orientations of hot configuration [3].

### Calculations and discussion

Fig. 1 shows the preformation probability  $P_0$  as a function of fragment mass for a fixed  $E_{c.m.}$  and  $\ell=0$  and  $\ell_{max}$  values. We notice that at  $\ell=0$  only the ER contribution is most prominent, whereas the same at  $\ell_{max}$  is prominent more for the fission fragments in the symmetric fission (SF) and near symmet-

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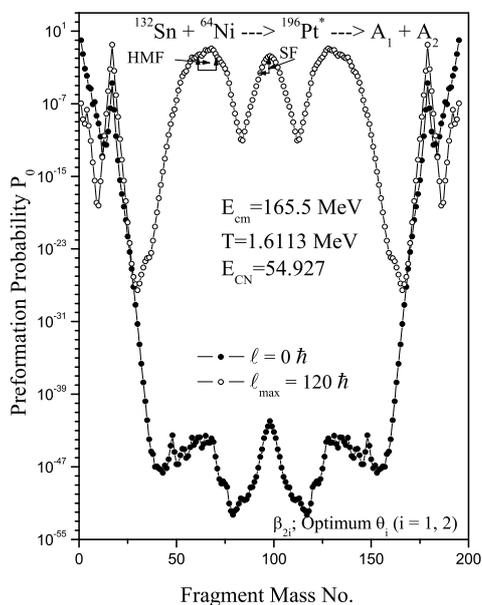


FIG. 1: Preformation probability as a function of fragment mass number for the decay of  $^{196}\text{Pt}^*$  at  $\ell_{min}$  and  $\ell_{max}$  values.

ric fission (nSF) regions. A small structure at intermediate mass fragmentation (IMF) starts appearing for  $\ell=0$ , but gets washed at higher  $\ell$ -values. The calculated cross-sections for the SF is very small compared to the nSF cross-section, i.e.,  $\sigma_{nSF} \sim 10^5 \times \sigma_{SF}$ . In other words, calculation of cross-sections for the fusion-fission of  $^{196}\text{Pt}^*$  favors asymmetric fission process, the SF contribution being almost negligible. The other component contributing to total cross-section is ER only. It may be noted that the asymmetric-fission peaks are placed around the target and projectile masses, and hence could contain a qf component, not measured in experiments.

Fig. 2 shows a comparison of our DCM calculated total fusion cross-section  $\sigma$  ( $=\sigma_{ER} + \sigma_{fiss} + \sigma_{qf}$ ) for the decay of  $^{196}\text{Pt}^*$  and the experimental data for the  $^{132}\text{Sn} + ^{64}\text{Ni}$  reaction [1]. The individual contributions of ER, fission, and qf (for the highest energies) are also shown, for both the data and DCM calculations. Apparently, the comparisons

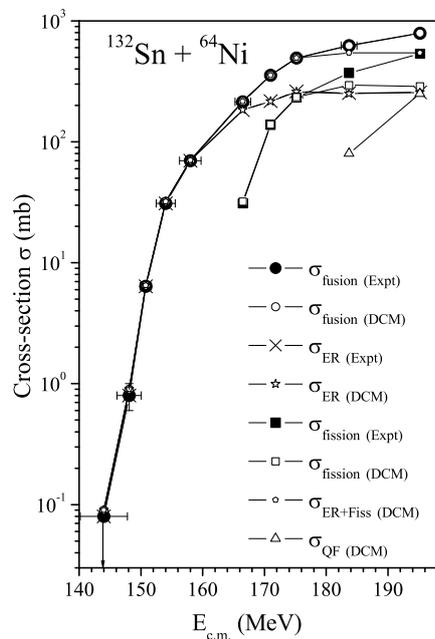


FIG. 2: The DCM calculated  $\sigma$ , and the individual  $\sigma_{ER}$ ,  $\sigma_{fiss}$ , and  $\sigma_{qf}$ , compared with experimental data on  $\sigma$  for the  $^{132}\text{Sn} + ^{64}\text{Ni}$  reaction.

between the experiments and DCM model are excellent for both the individual and total cross-sections, better than the CCFULL and density-constrained TDHF calculations (see Fig. 5 in first part of Ref. [1]), showing that a significant content of non-compound, qf process exists in this reaction, which need be measured.

## References

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