

Impact of spallation neutrons on criticality

Chitra Bhatia and V. Kumar*

HENP Laboratory(ADS Program), Physics Department, University of Rajasthan, Jaipur, India

* e-mail: ykv1951@gmail.com

Neutron spectrum in an accelerator driven sub critical system is different to the neutron spectra of thermal and fast reactors. The spallation neutron spectrum falls over a nuclear fuel in an ADS. Thus, for an arbitrary nuclear fuel, production and absorption cross sections of neutrons would vary with the incident neutron energy and the criticality coefficient, k_{eff} is expected to behave differently to a thermal or the fast reactor. For developing methods of estimation of k_{eff} in an ADS, let us assume that there is a cylindrical spallation target of dimension $d \times L = 20 \times 50 \text{ cm}^2$ enclosed in side a fuel blanket of thickness X . The energy spectra of spallation neutrons produced in collision of 1 GeV proton with a thick lead target is simulated using the CASCADE code [1] and the standard neutron spectra [2] of the thermal and a fast reactor have been used to estimate the spectrum average cross sections of different reactions occurring in different fuel elements such as ^{232}Th and ^{235}U of an arbitrarily assumed fuel system of an ADS. Spallation neutron spectrum spread up to several hundred MeV, may cause several reactions such as ^{232}Th ($n,3n$) ($E_{th} = 11.61 \text{ MeV}$), ^{232}Th ($n,3np$) ($E_{th} = 18.82 \text{ MeV}$) and ^{232}Th ($n,8n$) ($E_{th} = 43.55 \text{ MeV}$) etc. that were not possible in a critical reactor. The main problem arising because of spallation neutron spectrum is that in the energy range $E_n > 20 \text{ MeV}$ the experimental data is scarce and in this situation one has to depend on the calculations from the models or a deterministic code and a combination of a deterministic code and a Monte Carlo simulation code.

Thus, one can estimate k_{eff} from the knowledge of cross sections of all the production channels such as, i) single neutron type - (n,n'), (n,np), (n,nd), (n,nt), ($n,n\alpha$), (n,nh) and ($n,2np$) ii) multiple neutron type - (n, xn) where $x = 2, 3, \dots$, (n, f), ($n,2nh$),

($n,2nd$), ($n,2nt$), ($n,2npd$), ($n,2n2p$), ($n,3np$), ($n,3nd$), ($n,3nh$), ($n,3nt$) and ($n,3n\alpha$) and iii) neutron removal type - (n,γ), (n, p), (n, d), (n, t) and (n,α). The (n, f) reaction contributes in both production and the removal channels. Considering all the aforesaid reaction channels we have calculated [3] the k_{eff} for the three neutron spectra by using the following formulas and further details of the procedure will be published elsewhere. Let us assume I_0 is the incident neutron intensity falling on a material then the intensity of surviving neutrons after passing through x distance, $I_x = I_0 \exp(-x\Sigma_t)$, here, Σ_t is the total macroscopic cross section of a neutron in the given material, Considering $I_0 = 1$ for a single neutron intensity, then the removal term, R of the neutron by way of all removal processes, (n, γ), (n, p) and ($n,2n$) etc. may be written as,

$$R = (1 - \exp(-x\Sigma_t)) [(\Sigma(n, \gamma) + \Sigma(n,p) + \Sigma(n,d) + \dots + \Sigma(n,np) + \Sigma(n,2n) + \dots + \Sigma(n,9n) + \dots + \Sigma(n,f) + \Sigma(n,2np) + \dots + \Sigma(n,3n\alpha) \dots] / \Sigma_{tot} \quad (1)$$

and the remaining fraction, $1 - R = L$ may be assumed to leak out from the fuel system.

Similarly, a production term, P may be written as follows,

$$P = (1 - \exp(-x\Sigma_t)) [2\Sigma(n,2n) + 3\Sigma(n,3n) + 4\Sigma(n,4n) + \dots + 9\Sigma(n,9n) + \dots + \langle v \rangle \Sigma(n,f) + 2\Sigma(n,2np) + \dots + 3\Sigma(n,3n\alpha) + \dots] / \Sigma_{tot} \quad (2)$$

Here $\langle v \rangle$ = fission neutrons of a fuel element.

In the estimation of P , R and L contribution of the (n, n') channel is not included because of the obvious reasons. Thus,

$$k_{eff} = P / (R+L) \quad (3)$$

The elementary cross sections of different reactions up to 250 MeV neutron energy are calculated using the TALYS-1.0 code and for the energy 250 MeV the values of the cross sections are considered constant at the last value corresponding to 250MeV. This is a fair approximation because n-flux at $E_n > 250$ MeV is very small.

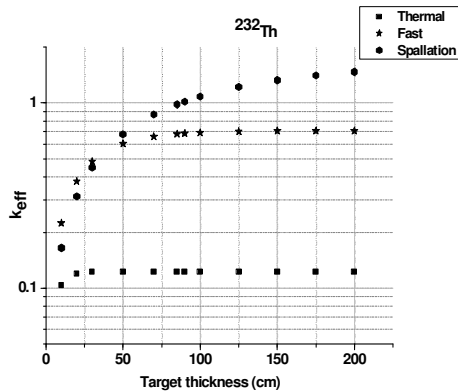


Fig.1a) Variation of k_{eff} with the target thickness for ^{232}Th target.

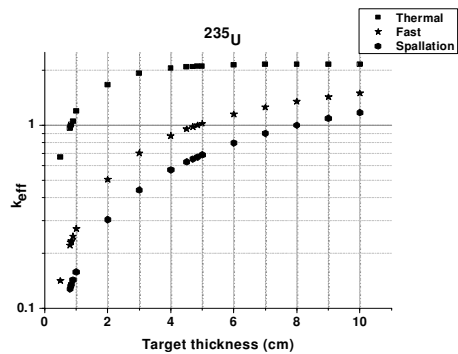


Fig.1b) Variation of k_{eff} with the target thickness for ^{235}U target.

From the data given in figs 1a) and b) we can infer that for the spallation spectrum k_{eff} is dominant in case of ^{232}Th while it is below the values corresponding to the thermal & fast spectra in case of ^{235}U . This shows that (n,xn) reaction play dominant role in case of fertile Th-fuel because of the presence of high energy

neutrons compared to the fissile ^{235}U fuel. This kind of behavior was pointed out earlier in ref. [4] in a detailed study of (n,xn) reactions in fertile and fissile fuels.

We have compared results of our calculations with that estimated by Cullen et al. [5] and a good agreement is seen in the two calculations.

Acknowledgement: We are thankful to the ILTP(DST) for the financial support.

References

- [1] V.Kumar, et.al Pramana-jn. of Physics 68(2)(2007) 315.
- [2] Accelerator Driven Systems (ADS) and Fast Reactors (FR) in Advanced Nuclear Fuel Cycles, A Comparative Study, www.nea.fr/html/ndd/reports/2002/nea3109.html
- [3] V.Kumar, Neutron Physics for ADS, a talk delivered at BHU, Sept. 19-20, 2009.
- [4] V.Kumar, Chitra Bhatia and H.Kumawat 'CASCADE Data for the A.D.S. Materials for its Benchmarking', published in ARIA'08 proceedings, PSI Proceedings 09-01, January 2009, ISSN 1019-0643, Pg no. 30.
- [5] D.E.Cullen et al. UCRL-TR-23310 (2007).