

Configuration assignment to the ^{156}Pm isomers

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Level structures of transitional nuclei ($A \approx 150-160$), particularly of neutron-rich species, are of a great interest, but are rather poorly defined. In this context, the latest Nuclear Data Sheets (NDS) evaluator for $A=156$ mass chain [1,2] remarked that assigning of a configuration to $Z = 61$ isotope ^{156}Pm ground state (gs) by Hellstrom et al. [3], and relating it to the 26.7-s activity present problems. In the present report, we critically examine this situation, taken together with more recent related experimental results [4] and Rotor Particle Model calculations [5].

Characterisation of the 26.7-s ^{156}Pm was undertaken by Hellstrom et al.[3] based on its decay to levels in ^{156}Sm ; they concluded that ‘the angular momentum of ^{156}Pm gs is at least 4 units’, and that ‘a consistent picture of ^{156}Pm decay is only found if gs is assumed to have $I^\pi K = 4^- 4^-$ ’ with most likely configuration:

$$4^- \{p_1: 5/2 [413] + n_0: 3/2 [521]\} \quad (\text{A})$$

wherein we have added subscripts on p/n to denote their excitational order as experimentally observed [2] in the neighbouring odd-mass isotope/isotone.

The NDS evaluator Reich [1] argued that this two-quasiparticle (2qp) configuration for the 4^- level, being a higher-lying singlet member of this GM doublet cannot be accepted as ^{156}Pm gs, since its GM triplet partner with $I^\pi K=1^- 1^-$ has to lie lower. He further opined that if the 4^- is indeed the gs and also has the expected $p_0:5/2[532]$ orbital as its constituent, the n-orbital has to have $K^\pi=3/2^+$ with the most likely possible configuration:

$$4^- \{p_0: 5/2 [532] + n_3: 3/2 [651]\} \quad (\text{B})$$

More recently, Shibata et al.[4], in their study of ^{156}Nd decay, identified an excited state isomer ($t < 5s$) in ^{156}Pm which de-excites with an

150.3 keV M3 transition to 26.7-s gs. This experimental observation of a higher-lying low-spin isomer conclusively rules out the (A) configuration assignment of Hellstrom et al.[1]. Shibata et al.[4], without any further analysis, just referred to the previously suggested [1] assignment (B) for 4^- ^{156}Pm gs and invoked GM rule to propose that the 150 keV isomer has $I^\pi=1^-$, being the singlet state of the GM doublet with previously suggested configuration (B).

However, whereas NDS evaluator [1], adopted the $p_0:5/2[532]$ orbital, as a constituent of ^{156}Pm gs, he somehow missed out checking the proximity of the suggested n-orbital $3/2[651]$ to the gs in adjacent $N=95$ isotones. A scan of the latest NDS evaluated database [2] yields the following values for the observed excitation energies of the $3/2[651]$ n-orbital in the respective $N=95$ odd-A isotones:

$$\begin{aligned} \text{Experimental } E_x : 3/2[651] \\ ^{157}\text{Sm} (571 \text{ keV}); ^{159}\text{Gd} (602 \text{ keV}); \\ ^{161}\text{Dy} (678 \text{ keV}); ^{163}\text{Er} (619 \text{ keV}). \dots\dots (1) \end{aligned}$$

Clearly any levels arising from coupling of this n-orbital (at $E_x > 570$ keV) with any p-orbital will lie above ~ 500 keV in the $N=95$ odd-odd nucleus; in particular, such a 2qp configuration, e.g., that in (B), cannot correspond to the gs of the $N=95$ ^{156}Pm nucleus.

We had earlier reported [5] the results of our calculation of the bandhead energies of various 2qp configurations using Quasiparticle Rotor Model (QPRM) with experimental inputs for the 1qp energies from neighbouring odd-A isotopes/isotones and for the GM splitting energies and Newby odd-even shift parameters from adjacent odd-odd nuclei. Therein it was concluded that the two ^{156}Pm isomers constitute the GM doublet having the 2qp configuration:

$$4^+_{\text{gs}} \{p_0: 5/2 [532] \pm n_0: 3/2 [521]\} 1^+_{150} \quad (\text{C})$$

with $I^\pi K=4^+4$ as ^{156}Pm gs and the 150.3 keV isomer having $I^\pi K=1^+1$.

We now proceed to critically examine the evidence from ^{156}Pm decay[1,3], i.e., log ft values to specific levels in ^{156}Sm , which was said to rule out positive parity and to favour negative parity for ^{156}Pm gs. Essentially, this includes 3 sets of log ft values which are discussed one by one in the following:

(I). Significant β -feeding of two levels around 2.5 MeV, which have almost identical γ -decays to 2^+ , 2^- , 4^+ and 4^- levels in ^{156}Sm , indicating $I=3$ assignment. Log ft values of 5.51 and 5.82 are indicative of allowed β branches, and hence same parity for both these levels as that of the parent. These data do not establish, or favour, negative parity for the parent state; they are equally consistent with its positive parity.

(II). β -feeding of 6^+ rotational level of the ground ($K^\pi = 0^+$) band with log ft = 9.0; this has been interpreted as indicative of 1^u transition, and hence favouring negative parity of the parent state. This argument overlooks the significant fact that this β -branch involves $\Delta K=4$, and hence is highly K-forbidden, in addition to ΔI effect. Presumably the reported log ft needs to be rechecked.

(III). The supposedly decisive evidence for negative parity for parent state comes from the feeding of 1515 keV 5^- level with log ft = 5.95 which is hence 'termed' as an allowed transition with $\Delta\pi=0$ and $\Delta I=1$. However, this interpretation is not as conclusive as presumed. Using the comprehensive global database for evaluated log ft values by Singh et al.[6], we find 27 cases of $1f$ ($\Delta\pi=\text{yes}$) transitions with log ft ≤ 6.1 just in the surrounding mass range $146 \leq A \leq 167$. Even more significantly, β -decay of our core nucleus ^{155}Pm ($I^\pi_{\text{gs}}=5/2^-$: our p_0 orbital) has [2] a β -branch with log ft =5.9 to 1362.1 keV ^{155}Sm level with an unambiguous $3/2^+$ assignment based on n-capture and subsequent γ -decays. By analogy, our β -branch with log ft =5.95 to 1515 keV ^{156}Sm 5^- level can as well be a $\Delta\pi=\text{yes}$ transition, and hence consistent with 4^+ assignment for the parent state.

Summarising, we conclude that the configuration (A) with 4^- as ^{156}Pm gs is not acceptable, since it requires a lower-lying low-spin GM triplet partner, whereas a low spin isomer has since been experimentally identified at 150.3 keV above the 24.6-s gs. Configuration (B) is also not acceptable, since its constituent n-orbital $3/2[651]$ is observed at excitation energy of more than 570 keV in every one of the $N=95$ odd-A isotones; consequently this $2q$ configuration will appear certainly at more than 500 keV above the corresponding lowest state in the $N=95$ odd-odd ^{156}Pm . Configuration (C) has been shown to be consistent with all the known experimental inputs, including the observed log ft values for β -feedings to ^{156}Sm levels which were earlier interpreted to exclusively suggest a negative parity parent state. The fact that the two isomers are connected through an M3 ($\Delta I=3$, $\Delta\pi=\text{no}$) transition, directly yields $I^\pi=1^+$ for the higher isomer. This spin-parity matches the

$$1^+1 \{p_0: 5/2 [532] - n_0: 3/2 [521]\}$$

assignment for the GM singlet partner of 4^+ gs; their observed separation of 150 keV is also similar to the experimental GM splitting of this configuration in isobaric ^{156}Eu and other odd-odd neighbouring nuclei.

References

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