

## QGP Phase Boundary and Plasma Lifetime From Thermal Properties of $\phi$ Mesons

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Heavy ion collisions at ultra-relativistic energies could produce a state of matter which is governed by partonic degrees of freedom, called Quark-Gluon Plasma (QGP). The nature of the deconfinement transition is of great importance. In this article, we have used  $\phi$  meson as a probe to study the QCD phase boundary and to estimate the QGP lifetime.

The transverse momentum ( $p_T$ ) spectra of  $\phi$ -mesons measured for AGS, SPS and RHIC energies have been analyzed in order to obtain the inverse slope parameter,  $T_{\text{eff}}$ . This effective temperature has contributions from both the (random) thermal and the collective motions in the transverse direction.  $T_{\text{eff}}$  of a hadron of mass ' $m$ ' can be related to the 'true' freeze-out temperature ( $T_{\text{th}}$ ) and average radial flow velocity ( $\langle v_r \rangle$ ) at the decoupling surface as:  $T_{\text{eff}} = T_{\text{th}} + \frac{1}{2}m \langle v_r \rangle^2$ . The  $\phi$  meson is expected to obtain most of its collective flow from the partonic phase, as it suffers from less hadronic rescattering. The  $\langle v_r \rangle$  values for central heavy ion collisions are obtained from the best fit of the blastwave formula to the  $p_T$  spectra. The large values of  $\langle v_r \rangle$  at RHIC energies indicate that QGP has undergone substantial radial flow. The similarity of the extracted  $T_{\text{th}}$  to the critical temperature  $T_c$  predicted by lattice calculations imply that the  $\phi$  meson freezes out near the phase boundary and could be used to extract the properties of QCD matter near the transition point.

A compilation of measured data for central heavy ion collisions for  $T_{\text{eff}}$ , for low  $p_T$  range

( $0 < p_T \leq 3.0$  GeV/c) and at mid-rapidity has been made. The results are depicted in Fig. 1 as a function of collision energy  $\sqrt{s_{\text{NN}}}$ .  $T_{\text{eff}}^\phi$  shows an increase with  $\sqrt{s_{\text{NN}}}$  (also with  $\epsilon_{\text{Bj}}$ , the Bjorken energy density) up to lower SPS energy. A further increase in collision energy leads to an increase of early temperature and pressure of the system. As a consequence, the  $p_T$  of the produced hadrons and hence the  $T_{\text{eff}}$  increase with collision energy. This is followed by the region of constant temperature which corresponds to higher SPS energies, where it is expected that the system undergoes a deconfinement transition with the creation of a co-existing phase of partons and hadrons, which is signaled by a plateau-like structure in the above spectrum. The resulting modification of Equation of State (EoS) suppresses the transverse expansion. This observation of a plateau-like structure is equivalent to a liquid-gas phase transition with the involvement of latent heat of the system, leading to a first order phase transition. Thus the experimental data indicate a first order phase transition, with a mixed phase stretching energy density between  $\sim 1$  and  $3.2$  GeV/fm<sup>3</sup> (figure to be shown), corresponding to higher SPS energy regime. At higher energies (corresponding to RHIC energies),  $T_{\text{eff}}^\phi$  further increases with collision energy. The EoS at the early stage becomes stiff, again leading to increase in early pressure and temperature. The thermal component,  $T_{\text{th}}^\phi$ , also shows a similar behavior with an increase in going toward the RHIC energies. This could be related to the onset of deconfinement corresponding to SPS energy regime.

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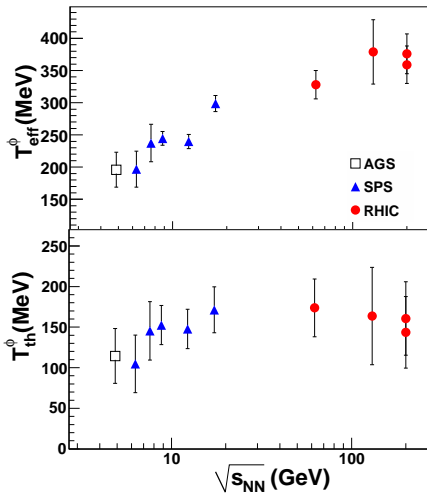


FIG. 1: Top panel:  $T_{\text{eff}}^{\phi}$  and the bottom panel:  $T_{\text{th}}^{\phi}$  as a function of  $\sqrt{s_{\text{NN}}}$ .

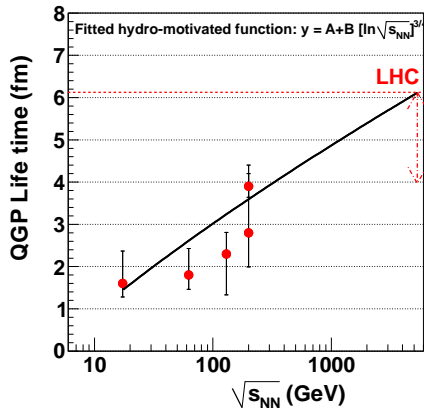


FIG. 2: QGP lifetime as a function of center-of-mass energy.

One of the important observables to look at a possible deconfinement phase transition is the effective thermodynamic degrees of freedom of the system,  $g$  which is given by:  $\frac{\epsilon}{T^4} = \frac{\pi^2}{30}g(T)$ , where  $\epsilon$  is the energy density of the

system. This value of  $g$  changes by a large factor while going from hadronic to a deconfined state of matter. For an ideal gas of massless, non-interacting constituents,  $g$  counts the number of bosonic degrees of freedom and fermionic degrees of freedom weighted by  $7/8$ . Thus for the deconfined QGP state with three flavors, it yields,  $g = 47.5$  and so  $\epsilon/T^4 \sim 16$ . IQCD calculations for  $\epsilon/T^4$  as a function of  $T$  show that energy density rises rapidly in a narrow temperature interval around  $T = 170$  MeV. If the transition occurs at  $T_c \sim 170$  MeV as expected, the temperature of  $\phi$  at the ‘true’ freeze out should saturate close to  $T_c$  which is observed in the present work.

To estimate the quantity,  $\epsilon/T^4$  at the phase boundary we need to know both the energy density and the temperature at the transition point. The value of  $T_{\text{th}}$  is determined from the  $p_T$  spectra of the  $\phi$ . For  $\epsilon$  we replace the formation time in the formula for  $\epsilon_{\text{Bj}}$  by some proper time  $\tau_c$  when the transition is over. Our aim is to estimate  $\tau_c$  which may be treated as the time that marks the end of the partonic phase. We compare the values of  $\epsilon/T^4$  obtained from the analysis of the data with the value obtained from lattice QCD calculations. The quantity  $\epsilon_{\text{Bj}}/T_{\text{th}}^4$  for RHIC is compared with the highest value of  $\epsilon/T^4$  obtained in IQCD calculations corresponding to value of  $g \sim 47.5$  for 3-flavor partonic phase and estimate the value of  $\tau_c$ . Similar procedure is followed for other colliding energies. The variation of  $\tau_c$ , the life time of plasma, with  $\sqrt{s_{\text{NN}}}$  is depicted in Fig. 2. We note that the  $\tau_c$  increases logarithmically with  $\sqrt{s_{\text{NN}}}$ . A value of  $\tau_c \sim 4 - 6$  fm/c is predicted at  $\sqrt{s_{\text{NN}}} = 5.5$  TeV for LHC [1]. The form of the fitted function is motivated by the fact that energy density scales logarithmically with collision energy and the time  $\sim \epsilon^{3/4}$ : a solution expected for asymptotically large  $\epsilon$  from the Bjorken scaling solution.

## References

- [1] R. Sahoo *et al.*, *Phys. Rev. C* (Under Review), arXiv: 1007.4335 [nucl-ex] and references therein.