

Prediction of energies of yrast band in some even-even nuclei

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Introduction:

The deformation parameter β and γ of the collective model of Bohr and Mottelson [1] are basic descriptors of the nuclear equilibrium shape and structure. The researchers found that the values of γ obtained from energies ($= \gamma_e$) are nearly equal to the value of γ derived from transition rate ($= \gamma_b$) in even Xe, Ba and Ce nuclei (A~120-140) and Hf, W, Os, Pt and Hg nuclei (A~160-200) using rigid triaxial rotor model of Davydov – Filippov [2-3]. In the present study, the relatively light mass nuclei (Mo, Ru and Pd) have been taken. As far as γ is concerned, it is known that the Ru chains of nuclei is intermediate between the two having opposite trends for parameter γ , decreasing for Mo & increasing for Pd, and has an irregular behaviour in itself with the increase of neutron number [4].

Procedure for evaluating asymmetric parameter (γ) and moment of inertia (I_0):

The Hamiltonian in Asymmetric Rotor Model of Davydov – Filippov [5] is written as

$$H = \frac{\hbar^2}{2} \sum \frac{I_i^2}{J_i} \quad (1)$$

Where, I_i is the projection of the angular momentum on the intrinsic axes. The moment of inertia of the model are given by the hydrodynamic relation

$$J_i = \frac{4}{3} J_0 \sin^2(\gamma - \frac{2\pi}{3}) \quad (2)$$

$$\text{Where, } J_0 = 4B\beta^2 \quad (3)$$

Simple analytical expressions for the energy of two levels of the required symmetry are, for $I = 2$, defined by the expressions

$$E_{2^+ 1,2}^+ = \frac{9[1-(-1)^{\sigma_{1,2}}] \sqrt{1-\frac{8}{9}X}}{X} \quad (4)$$

$$\text{In the units of } \hbar^2/J_0 \quad (5)$$

where $\sigma_{1,2} = 0, 1$
and $X = \sin^2 3\gamma$

(a) Determination of asymmetric parameter (γ):

The energy ratio $E_{2^+ 2}^+/E_{2^+ 1}^+$ given by equation (6) depends only on γ and this is evaluated by feeding experimental energies $E_{2^+ 2}^+$ & $E_{2^+ 1}^+$ of the two 2^+ states.

$$\frac{E_{2^+ 2}^+}{E_{2^+ 1}^+} = \frac{1 + \sqrt{1 - \frac{8}{9}X}}{1 - \sqrt{1 - \frac{8}{9}X}} \quad (6)$$

We denote the parameter obtained from these energies by γ_e . In another way the branching ratio $B(E2; 2_2^+ \rightarrow 2_1^+)/B(E2; 2_2^+ \rightarrow 0_1^+)$ which is generally known experimental quantity and is a function of γ alone is used to determine this parameter as –

$$\frac{B(E2; 2_2^+ \rightarrow 2_1^+)}{B(E2; 2_2^+ \rightarrow 0_1^+)} = \frac{\frac{10}{7} \left[\frac{X}{9-8X} \right]}{\frac{1}{2} \left(1 - \frac{3-2X}{\sqrt{9-8X}} \right)} \quad (7)$$

The parameter so obtained is denoted by γ_b .

(b) Determination of Moment of Inertia (I_0):

The moment of inertia is extracted from equation (4) using the energy of first excited 2^+ state and is denoted by I_0

$$I_0 = \frac{J_0}{\hbar^2} = \frac{9 - \sqrt{81 - 72X}}{XE_1^+} \quad (8)$$

Discussion:

The asymmetric parameters γ_e and γ_b are evaluated from experimental level energies and E2 transition rates and are tabulated (Table I). There is no general trend in values of γ with neutron number but, one thing is clear that the values of γ_b and γ_e are almost equal in all of the nuclei of Mo, Ru and Pd. The γ_b values are plotted against its valence nucleon product $N_p N_n$, an external parameter in fig. 1. The smooth behavior of γ with $N_p N_n$ is surprising. Although, this has been shown before in other regions i.e A~120-140 and A~160-200 [2, 3] but, in A~90-120 region where variation of γ in Ru, Mo and Pd

are opposite with $N_p N_n$ when considered individually (Table I).

The asymmetric parameter γ , for those nuclei whose energy of 2_2^+ state are not known, have been predicted from fig.1. This value of γ is used to calculate the energies of 4^+ , 6^+ and 8^+ state of yrast band using ARM for some of Ru, Pd and Mo nuclei and are listed in Table II. The systematic study of γ shape variables and moment of inertia I_0 in Mo, Ru and Pd nuclei concludes that the γ values evaluated separately from energies and transition rates are found very close to each – other and this gives confidence in the values obtained.

Table I

The values of moment of inertia I_0 and asymmetric parameter γ for even – even Mo, Ru and Pd Nuclei alongwith $N_p N_n$

Nucleus	R_e	R_b	γ_e	γ_b	I_0
^{92}Mo	2.05	-	28.0	-	3.93
^{94}Mo	2.14	-	27.0	-	6.73
^{96}Mo	2.09	-	28.0	-	7.63
^{98}Mo	2.01	-	29.0	-	6.86
^{100}Mo	1.90	30.25	30.0	26.0	11.21
^{102}Mo	2.86	-	23.0	-	18.28
^{104}Mo	4.22	5.80	19.0	20.0	25.84
^{106}Mo	4.14	4.20	19.0	17.5	28.90
^{98}Ru	2.17	50.00	27.0	26.8	8.98
^{100}Ru	2.52	24.40	24.0	25.5	10.28
^{102}Ru	2.32	26.30	25.5	25.7	12.02
^{104}Ru	2.49	18.52	24.5	24.5	15.65
^{106}Ru	2.93	-	22.5	-	19.84
^{108}Ru	2.93	10.20	22.5	22.3	22.12
^{110}Ru	2.54	14.20	24.0	23.8	23.04
^{112}Ru	2.21	25.30	26.0	25.5	23.98
^{102}Pd	2.76	8.30	23.0	21.5	9.96
^{104}Pd	2.41	20.00	25.0	25.0	10.18
^{106}Pd	2.20	37.40	26.5	26.3	11.35
^{108}Pd	2.15	71.40	27.0	27.3	13.51
^{110}Pd	2.18	71.43	26.7	27.3	15.58
^{112}Pd	2.11	-	28.0	-	17.00
^{114}Pd	2.09	-	28.0	-	17.82
^{116}Pd	2.16	-	27.0	-	17.81

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Table – II

Predicted values of energies (in KeV) of 4^+ , 6^+ and 8^+ states of yrast band for some Mo, Ru and Pd Nuclei

Nucl	γ	$E2_1^+$	$E4_1^+$	$E6_1^+$	$E8_1^+$
^{84}Mo	24	443.8	1284.8	2393.4	3506.9
^{86}Mo	25	566.6	1606.8	2940.0	4276.1
^{88}Mo	26	740.5	2058.6	3720.2	5435.2
^{90}Mo	30	947.9	2527.7	4860.8	6918.7
^{88}Ru	26	616.0	1713.3	3094.8	4521.4
^{90}Ru	27	738.1	2017.7	3642.7	5317.3
^{92}Ru	28	864.6	2330.8	4229.3	6175
^{94}Ru	30	1430.	3813.3	7330.0	10423.3
^{114}R	23	127.0	375.3	710.0	763.6
^{94}Pd	28	814.2	2195.9	3984.5	5817.6
^{96}Pd	30	1415.	3773.3	7256.1	10328
^{98}Pd	28	863.1	2328.1	4224.3	6167.8
^{100}P	27	665.5	1818.1	3282.4	4792.6
^{118}P	26	378.4	1051.3	1899.1	26

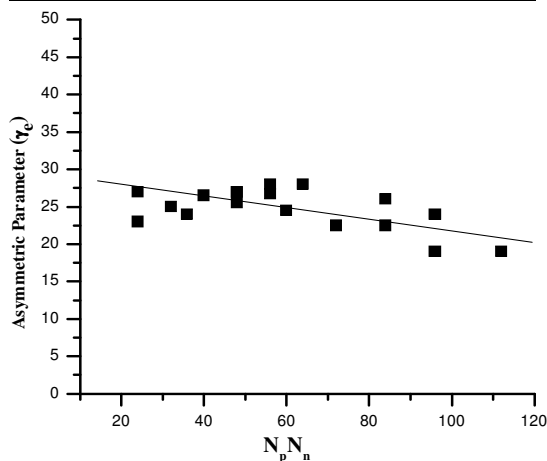


Fig. 1 The variation γ values plotted against its valence nucleon product $N_p N_n$ for Mo, Ru and Pd nuclei

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