

## Fusion excitation functions of $^{66}\text{As}^*$ formed in $^8\text{B} + ^{58}\text{Ni}$ reaction at near barrier energies.

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### Introduction

The study of proton-halo systems has become an interesting area of research in recent years. This is so because the knowledge of p-halo's is very scarce in comparison to neutron-halo systems. In view of having worthy information regarding nuclear structure and dynamics related to p-halo systems, the fusion of exotic p-halo  $^8\text{B}$  nucleus with  $^{58}\text{Ni}$  target has been studied. The experimental data for the fusion cross-sections of  $^8\text{B} + ^{58}\text{Ni} \rightarrow ^{66}\text{As}^* \rightarrow A_1 + A_2$  reaction has been reported at various energies around the Coulomb barrier [1]. The same is tested in framework of dynamical cluster decay model (DCM)[2] and the  $\ell$ -summed Wong model [3]. Since the projectile involved is halo, the extended radius of halo nuclei is assimilated in DCM through the neck-length parameter  $\Delta R$ . However, in Wong model calculations 26% increased radius is employed as suggested in experimental paper [1].

### The Model

#### Dynamical cluster decay model (DCM):

The DCM, carries distinct advantage over available statistical models as it treats the Evaporation Residue (ER), intermediate mass fragment (IMF) and fusion fission (ff) on equal footing [2]. The missing nuclear structure information of the Compound Nucleus in statistical models enters in the DCM via the preformation probability  $P_0$  of the fragments and is calculated by solving stationary Schrödinger equation in  $\eta$  co-ordinate. For  $\ell$ -partial waves, the compound nucleus decay cross-section is given by

$$\sigma = \frac{\pi}{k^2} \sum_{\ell=0}^{\ell_{max}} (2\ell + 1) P_0 P; \quad k = \sqrt{\frac{2\mu E_{c.m.}}{\hbar^2}} \quad (1)$$

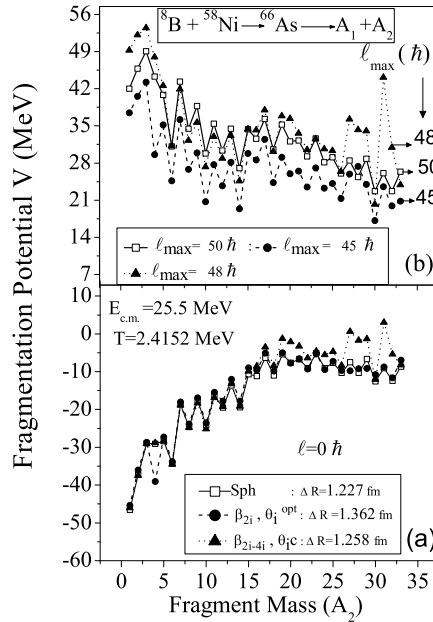


FIG. 1: Variation of fragmentation potential with fragment mass  $A_2$  at (a)  $\ell=0\hbar$  and (b)  $\ell=\ell_{max}$

where  $\mu = [A_1 A_2 / (A_1 + A_2)] m$  is the reduced mass,  $m$  is the nucleon mass, and  $\ell_{max}$  is the maximum angular momentum. The angular momentum  $\ell_{max}$  is fixed for the vanishing of the fusion barrier of incoming channel  $\eta_i$  or light particle cross-section  $\sigma_{LP} \rightarrow 0$ . The preformation probability  $P_0$  refers to  $\eta$  motion and the penetrability  $P$  to  $R$  motion, both depending on angular momentum  $\ell$  and temperature  $T$  of compound nucleus state. The DCM has been used to study the decay of  $^{66}\text{As}^*$  in terms of light particles (LPs), intermediate fragments (IMFs) and the fission fragments.

#### The $\ell$ -summed Wong model:

The extended  $\ell$ -summed Wong Model is a special case of DCM where cross-sections are cal-

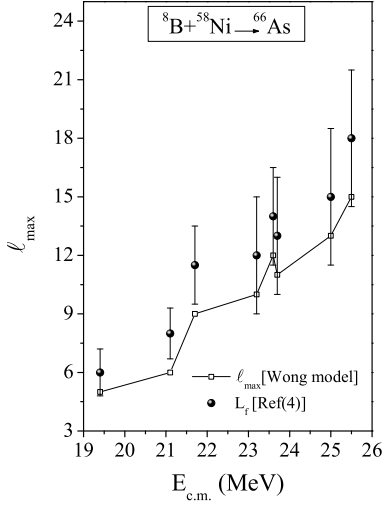


FIG. 2: Deduced  $l_{max}$  values using  $l$ -summed Wong model.

culated through Eq.(1), by taking  $P_0=1$  and penetrability  $P$  calculated in Hill-Wheeler approximation, which in DCM is calculated by the WKB integral for each decay channel. The penetrability  $P$ , in terms of the barrier height  $V_B$  and curvature  $\hbar\omega_\ell$ , is given as

$$P = \left[ 1 + \exp \left( \frac{2\pi(V_B - E_{c.m.})}{\hbar\omega_\ell} \right) \right]^{-1} \quad (2)$$

with  $\hbar\omega_\ell$  evaluated at the barrier position  $R = R_B$  corresponding to barrier height  $V_B$ .

### Calculations and Discussions

The fusion excitation function have been calculated for the reaction in which exotic proton-halo nucleus  ${}^8B$  is incident on a  ${}^{58}Ni$  target at several energies near the Coulomb barrier [1], using dynamical cluster-decay model (DCM) and  $l$ -summed Wong model. In framework of DCM, the fusion cross-section is the sum of decay processes involving ER, IMFs, and fission fragments, and the calculations have been done for spherical,  $\beta_2$  deformed and  $\beta_2$ - $\beta_4$  deformed choice of fragmentation. At  $E_{c.m.}=25.5$  MeV, the fusion cross-section of 864.67mb has been fitted within one parameter of DCM, the neck-length parameter  $\Delta R$ . The cross-section of LPs is 333.55mb, IMFs is 445.73 mb and that of fission fragments is 84.52 mb. Similar trend of decay

cross-sections is observed at all other energies. These results help us to conclude that IMFs are the major contributors towards the decay of  ${}^{66}As^*$  followed by LPs and only 10% contribution is through fission. As shown in Fig.1 (a) and (b), the fragmentation path is almost independent of deformation effects for  $\ell=0\hbar$ . However, at  $\ell=l_{max}$  the potential energy surfaces get influenced with the inclusion of deformation and orientation effects. The fragmentation is more prompt for  $\beta_2$  deformation. Besides this, DCM has also been tested for one of non-compound nucleus process i.e the incomplete fusion (ICF). In this case the break-up of  ${}^8B$  into  ${}^7Be$  is observed, which in turn interact with the target  ${}^{58}Ni$ . After applying appropriate energy correction the decay cross-section for  ${}^8B + {}^{58}Ni \rightarrow {}^{65}Ge^* \rightarrow A_1 + A_2$  are calculated. The results obtained (for  $\beta_2$  deformed case) are in agreement with the ones reported in [1]. Finally, the fusion cross-sections have also been calculated by extended-Wong model. The cross-sections have been fitted at all the energies except one lying below the barrier as it requires modification at below barrier region. The deduced  $l$ -values from Wong calculations are in good agreement with those of  $L_f$ -values extracted from the fusion data as is clear from Fig.2. The  $l$ -values in case of DCM are relatively of higher magnitude because sticking moment of inertia is used in these calculations to account for proximity interaction. In summary the fusion-fission cross-section compare nicely with experimental data. The role of deformation is seen in decay of  ${}^{66}As^*$  and the contribution of IMFs is found to dominate, followed by LPs and fission.

### References

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