

## On the missing 1n contribution in $^{12}\text{C}+^{93}\text{Nb}\rightarrow^{105}\text{Ag}^*$ reaction at below barrier energies

Sahila Chopra<sup>1</sup>, Manie Bansal<sup>1</sup>, Manoj K. Sharma<sup>2</sup>, and Raj K. Gupta<sup>1</sup>

<sup>1</sup>Physics Department, Panjab University, Chandigarh-160014, INDIA and

<sup>2</sup>School of Physics and Materials Science, Thapar University, Patiala 147004, INDIA

### Introduction

In a recent experiment [1], the excitation functions (EFs) of various evaporation residues were measured in the decay of compound nucleus  $^{105}\text{Ag}^*$  formed in  $^{12}\text{C}+^{93}\text{Nb}$  reaction at below barrier energies. The observed decay products are the heavy residues whose complementary fragments are the light particles (LPs, with  $\text{mass}\leq 4$ ) and intermediate mass fragments (IMFs,  $4\leq\text{mass}\leq 13$ ). The data is compared with statistical model code PACE2, which gives the complete fusion (CF) cross-section consisting of 2n, 3n, 4n and possibly  $^4\text{H}$ . The disagreement of PACE2 for heavier residues, the IMFs, is taken by the authors as signatures of incomplete fusion ICF and/or the direct reaction process, supported by their recoil range distribution measurements. The ICF consists of (complementary fragments) of  $^4,5,6\text{H}$ ,  $^4,5\text{He}$ ,  $^8\text{Li}$ ,  $^9,10,11\text{Be}$ ,  $^{12}\text{B}$ , and  $^{13}\text{C}$ . What is striking of the PACE2 predictions is that 1n emission is, in general, found missing (see, e.g., also Ref. [2] for  $^{64}\text{Ni}+^{100}\text{Mo}$  reaction) whereas the dynamical cluster-decay model [3] predict the contribution due to 1n to be the largest to total evaporation residue cross-section.

In this contribution, we attempt to determine the missing 1n content of evaporation residue cross-section in the decay of  $^{105}\text{Ag}^*$  formed in reaction  $^{12}\text{C}+^{93}\text{Nb}$ , using the dynamical-cluster decay model (DCM) [3, 4] at an illustrative lab. energy  $E_{\text{Lab}}=61.2$  MeV ( $E_{\text{c.m.}}=54.205$  MeV). It is important to stress that the study of EFs serves as a powerful tool to understand the reaction mechanism governing the collision dynamics, ranging from CF, ICF to more complex non-compound nucleus (nCN) pre-equilibrium or quasi-fission

(qf) processes. Our DCM calculations include deformation effects up to hexadecapole deformations ( $\beta_2, \beta_3, \beta_4$ ) with compact orientations of hot fusion process, for the case of coplaner nuclei only (azimuthal angle  $\Phi=0^\circ$ ).

### Dynamical cluster-decay model

The DCM of Gupta and collaborators [3, 4] is based on collective coordinates of mass (and charge) asymmetries  $\eta$  (and  $\eta_z$ ) [ $\eta = (A_1 - A_2)/(A_1 + A_2)$ ,  $\eta_z = (Z_1 - Z_2)/(Z_1 + Z_2)$ ], and relative separation R, with multipole deformations  $\beta_{\lambda i}$  ( $\lambda=2,3,4$ ;  $i=1,2$ ), and orientations  $\theta_i$ . In terms of these coordinates, we define the compound nucleus decay cross section for  $\ell$  partial waves as

$$\sigma = \frac{\pi}{k^2} \sum_{\ell=0}^{\ell_{\text{max}}} (2\ell + 1) P_0 P; \quad k = \sqrt{\frac{2\mu E_{\text{c.m.}}}{\hbar^2}}$$

where  $P_0$  is preformation probability referring to  $\eta$  motion and  $P$ , the penetrability, to R motion, both dependent on  $\ell$  (angular momentum) and T (temperature).  $\mu$  is reduced mass with m as the nucleon mass.  $\ell_{\text{max}}$  is maximum angular momentum, defined for light particle evaporation residue cross section  $\sigma_{ER} \rightarrow 0$ .  $P_0$  is the solution of stationary schrödinger equation in  $\eta$  and  $P$  is the WKB penetrability of preformed fragments. Note that  $P_0$  imparts the important nuclear structure information which is otherwise missing in alternative statistical models. The only parameter of the model is the T dependent neck-length parameter  $\Delta R(T)$ , defining the first turning point  $R_a = R_1(\alpha_1, T) + R_2(\alpha_2, T) + \Delta R(T)$  for the penetration  $P$ . DCM is also applied to nCN, qf process where  $P_0=1$  since the incoming channel remains unaffected.

**Table 1:** Experimental and calculated channel cross-sections at  $E_{c.m.}=54.205$  MeV for fitted  $\Delta R=1.363$  fm,  $\ell_{max}=76 \hbar$ .

No. of neutrons emitted	Cross-section (mb)	
	Experimental	Calculated
1n	—	75.4
2n	$7.6 \pm 1.4$	7.6
3n	$398.2 \pm 46.9$	0.156
4n	$203.4 \pm 26.4$	0.0022

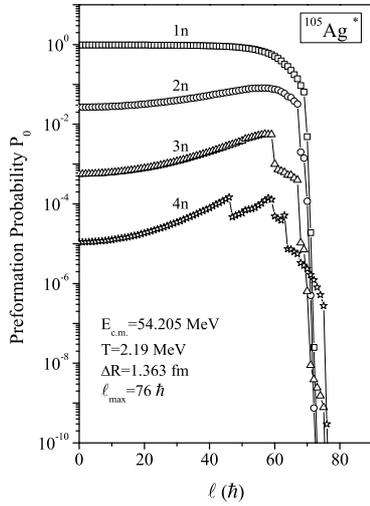


FIG. 1: Preformation probability  $P_0$  as a function of angular momentum  $\ell$  for  $xn$  decay of  $^{105}\text{Ag}^*$  formed in  $^{12}\text{C}+^{93}\text{Nb}$  reaction at  $E_{c.m.}=54.205$  MeV, using  $\Delta R=1.363$  fm and  $\ell_{max}=76 \hbar$ .

## Calculations and results

First of all, we look for the possible decay processes in  $^{105}\text{Ag}^*$  formed in  $^{12}\text{C}+^{93}\text{Nb}$  reaction. We are interested in evaporation residue channels  $xn$ ,  $x=2,3,4$  populated by CF. Also,  $^4\text{H}$  residue channel could be due to CF. Apparently, the 1n evaporation channel is not observed in this reaction.

Our calculated fragmentation potential  $V(\eta)$  on DCM, for best fitted  $\Delta R$  to the measured 2n cross-section, gives potential energy minima at 1n, 2n, 3n and  $^4\text{H}$ , i.e.  $^4\text{H}$  is energetically more favoured than 4n. Though  $^4\text{H}$  could also be taken as CF product, we have replaced it by the binding energy of 4n (and

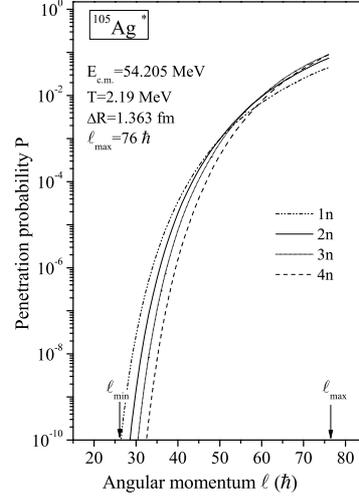


FIG. 2: Same as for Fig. 1, but for the penetrability  $P$ .

the complementary heavy fragment).

Figs. 1 and 2 show the calculated  $P_0$  and  $P$  as a function of  $\ell$  for  $xn$ ,  $x=1,2,3,4$  channels. Interesting enough, 1n channel competes with other three neutron channels. In fact 1n product is not only most strongly preformed, but has the largest penetrability. In other words, the  $\sigma_{1n}$  makes the largest contribution to evaporation residue cross-section  $\sigma_{ER} = \sum_x \sigma_{xn}$ . Table 1 suggests the use of different  $\Delta R$  for 3n and 4n emissions, which means different reaction times for different residue products.

Concluding, fitted  $\Delta R$  for 2n emission in DCM suggests that the experimentally missing 1n contribution is the largest to CF cross-section, and that the reaction times are likely to be different for different  $xn$  emissions.

## References

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