

S.V. Surayanarayana<sup>a</sup>, V. K. Mulik<sup>b</sup>, Sadhana Mukerji<sup>c</sup>, H. Naik<sup>d</sup>, P. M. Prajapati<sup>e</sup>, B. S. Sivashankar<sup>f</sup>,  
 A.K. Mohanty<sup>a</sup>, A. Goswami<sup>d</sup>, S. D. Dhole<sup>a</sup>, V.N. Bhoraskar<sup>a</sup>

<sup>d</sup>Nuclear Physics Division Bhabha Atomic Research Centre, <sup>b</sup>Department of Physics, University of Pune - 411 007, India,

<sup>c</sup>Reactor Physics Design Division <sup>d</sup>Radiochemistry Division, Bhabha Atomic Research Centre, Mumbai -400 085, India

<sup>e</sup>Physics Department, Faculty of Science, The M. S. University of Baroda, Vadodara-390 002, India

<sup>f</sup>Department of Statistics, Manipal University, Manipal - 576 104, India.

Presently, the reactors operating over the globe are light or heavy water reactors that use enriched or natural uranium as a fuel. However, recent innovations in reactor technology are aiming for nuclear power based on the fast reactor and advanced heavy water reactor (AHWR) [1,2] in order to meet the increased demand of power production. In the fast reactor, the fissile material  $^{239}\text{Pu}$  along with  $^{238}\text{U}$  is used as a fuel for power generation. Therefore isotope  $^{238}\text{U}$  is very important in this regard, as  $^{238}\text{U}$  is used as the breeding material to regenerate the fissile material  $^{239}\text{Pu}$ . The required initial stock of  $^{239}\text{Pu}$  is first generated in a thermal reactor from  $^{238}\text{U}(n, \gamma)^{239}\text{U}$  reaction followed by two successive  $\beta$  decays. Therefore, in the present work the  $^{238}\text{U}(n, \gamma)$  and  $^{238}\text{U}(n, 2n)$  reaction cross-sections have been studied. These reactions were determined at average neutron energy of  $15.5 \pm 0.7$  MeV using the neutron beam from  $^7\text{Li}(p, n)$  reaction and by activation technique followed by off-line  $\gamma$ -ray spectrometry.

The neutron beam for the present work was obtained from the  $^7\text{Li}(p, n)$  reaction, utilizing the proton beam of energy 18 MeV and current of 250 nA at 14UD BARC-TIFR Pelletron facility, Mumbai, India. The natural  $^{238}\text{U}$  metal foil of 571.5 mg/cm<sup>2</sup> thickness was wrapped with 0.025 mm thick aluminum foil and the sample was irradiated for 5 h with neutrons. After sufficient cooling, the  $\gamma$ -rays of fission/reaction products from the irradiated U sample was counted in an energy and efficiency calibrated 80 c.c. HPGe detector coupled to a PC-based 4K channel analyzer.

The half-life of  $^{239}\text{U}$  is 23.54 min., which decays 99.6% to  $^{239}\text{Np}$  within 3 h. In view of this, the  $^{238}\text{U}(n, \gamma)$  reaction cross-section was calculated from the observed photo-peak activity of  $^{239}\text{Np}$  and was identified through the  $\gamma$ -lines of 103.73, and 277.85 keV. On the other hand the  $^{238}\text{U}(n, 2n)$  reaction was calculated from the 208.0 keV  $\gamma$ -line of  $^{237}\text{U}$ . The fission products from  $^{238}\text{U}$  have varying half-lives. The net photo-peak area ( $A_{\text{net}}$ ) for  $\gamma$ -lines of  $^{239}\text{Np}$  and for different  $\gamma$ -lines of fission products (e.g.

743.3 keV of  $^{97}\text{Zr}$ ) were obtained from their total peak area after subtracting the linear background due to Compton effects. From the  $A_{\text{net}}$  of a particular fission products (e.g.  $^{97}\text{Zr}$ ), neutron flux ( $\phi$ ) was obtained using decay equation (1)

$$A_i = N\sigma\Phi a\epsilon Y (1 - e^{-\lambda t}) e^{-\lambda T} (1 - e^{-\lambda \Delta T}) / \lambda \quad (1)$$

where N is the number of atoms of the isotope of the element and  $\sigma$  is the fission cross-section of  $^{238}\text{U}$ . Y is the cumulative yield of  $^{97}\text{Zr}$ . 'e' is the detection efficiency, 'a' is the  $\gamma$ -ray abundance and  $\lambda$  is the decay constant of the product nuclide. 't', T and  $\Delta T$  are irradiation, cooling and counting time respectively. The incident proton energy in the present experiment was 18 MeV. The neutrons from  $^7\text{Li}(p, n)$  reaction are not mono-energetic, and their energy spectra were obtained from literature [3]. The average neutron energy corresponding to quasi mono energetic peak was obtained as  $15.5 \pm 0.7$  MeV after removing the tailing distribution of the neutron spectrum below 13 MeV. The neutron flux ( $\Phi$ ) of  $3.78 \times 10^9$  n cm<sup>-2</sup> h<sup>-1</sup> was used to calculate the  $^{238}\text{U}(n, \gamma)$  reaction cross-section, which is  $2.389 \pm 0.075$  mb. Similarly, the  $^{238}\text{U}(n, 2n)$  cross-section was calculated using the neutron flux of  $2.15 \times 10^9$  n cm<sup>-2</sup> h<sup>-1</sup>, which is  $1094.4 \pm 75.6$  mb.

For  $^{238}\text{U}(n, \gamma)$  reaction, the low energy neutrons below 13 MeV also contribute to the cross-section. This contribution from the tail was estimated using the ENDF/B-VII.0 [4] and JENDL-4.0 [5] by folding the cross-sections with neutron flux distributions of ref. [3]. The contributions from the above evaluations are 1.912 and 1.476 from ENDF/B-VII.0 and JENDL-4.0 respectively. The actual cross-section was obtained as  $0.623 \pm 0.075$  mb, which are given in Table 1. Similarly, the tailing contribution for  $^{238}\text{U}(n, 2n)$  reaction was calculated from 6.16 to 13 MeV, which are 565.36 mb from ENDF/B-VII.0 and 550.23 mb from JENDL-4.0. The actual cross-section was obtained as  $536.58 \pm 75.67$  mb, which is given in Table 1. The evaluated  $^{238}\text{U}(n, \gamma)$  and (n, 2n) reaction cross-sections from ENDF/B-VII.0 and JENDL-4.0

are also given in Table 1 for comparison. It can be seen from Table 1 that the measured  $^{238}\text{U}(n, \gamma)$  and  $(n,2n)$  reaction cross-sections are almost within the range of the evaluated data of ENDF/B-VII.0 and JENDL 4.0.

Table 1.  $^{238}\text{U}(n, \gamma)^{239}\text{U}$  and  $^{238}\text{U}(n, 2n)^{237}\text{U}$  reaction cross-sections at neutron energy of  $15.5 \pm 0.7$  MeV.

Neutron Flux ( $\text{n cm}^{-2}\text{h}^{-1}$ )	Cross-section (mb)		
	Expt.	ENDF/B-VII	JENDL-4.0
	$^{238}\text{U}(n, \gamma)^{239}\text{U}$		
$3.78 \times 10^9$	$0.623 \pm 0.075$	0.710-0.580	0.415
	$^{238}\text{U}(n, 2n)^{237}\text{U}$		
$2.15 \times 10^9$	$536.58 \pm 75.67$	508.3	483.4

To examine this, the  $^{238}\text{U}(n, \gamma)$  and  $(n,2n)$  reaction cross-sections from the present work and evaluated data from ENDF/B-VII.0 and JENDL-4.0 JEFF 3.1/A [6] and CENDL-3.1[7] along with data given in EXFOR [8] along with are plotted in Fig. 1 and Fig.2, respectively. It can be seen from Fig. 1 that the  $^{238}\text{U}(n, \gamma)$  reaction cross-section at neutron energy of 15.47 MeV is in agreement with the value of Mc Daniels et al [9] at 11.2-12.2 MeV but not with CENDL-3.1 [7] evaluation. The  $^{238}\text{U}(n,2n)$  reaction cross-sections from the present work are in agreement with evaluated data (Fig. 2).

Theoretically, the  $^{238}\text{U}(n, \gamma)$  reaction cross-sections at different neutron energy beyond 100 keV and  $^{238}\text{U}(n,2n)$  above threshold were also calculated using the nuclear model based computer code TALYS 1.2 [10] and are plotted in the Fig. 1 and Fig.2. It can be seen that the trend of the experimental and evaluated  $^{238}\text{U}(n, \gamma)$  and  $(n,2n)$  reaction cross-sections are well produced by the TALYS 1.2. However, theoretical  $^{238}\text{U}(n, \gamma)$  reaction cross-section from TALYS are slightly higher than the experimental and evaluated values for neutron energy from 100 keV to 2 MeV. Further, it can be seen from Fig. 1 that the experimental, evaluated and the theoretical  $^{238}\text{U}(n, \gamma)$  reaction cross-section decreases from 100 keV to 7 MeV and predict a dip in around 7.3-8.0 MeV. Beyond 8.0 MeV, it increases up to neutron energy of 14 MeV and then again decreases. The dip in the  $^{238}\text{U}(n, \gamma)$  reaction cross-section around neutron energy of 7.5-8.5 MeV indicates the opening of  $(n, 2n)$  reaction channel besides  $(n, nf)$  channel.

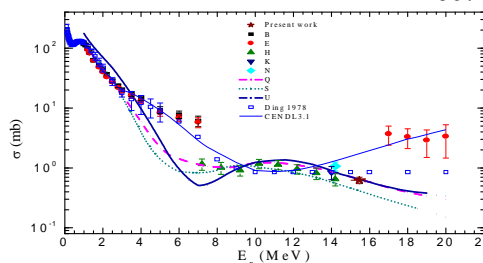


Fig. 1: Plot of  $^{238}\text{U}(n, \gamma)$  reaction cross-section as a function of neutron energy

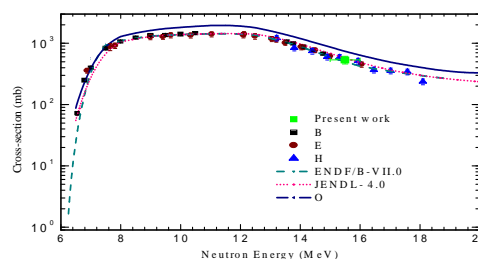


Fig. 2: Plot of  $^{238}\text{U}(n,2n)$  reaction cross-section as a function of neutron energy

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