

## Search the nature of multiphonon 2 $\gamma$ -band of $^{158}\text{Dy}$

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### Introduction

Multiphonon 2 $\gamma$ -band in deformed nuclei has been studied experimentally and theoretically during the recent decades. In deformed nuclei, the multiphonon excitations can stipulate worthy information of nuclear structure. Zamfir and Casten [1] have reviewed a number of signatures of  $\gamma$ -softness versus  $\gamma$ -rigidity in nuclei, giving attention on the staggering properties of  $\gamma$ -band energies. Liao Ji-Zhi [2] studied the behavior of staggering indices  $S(I)$  in both axial and triaxial rotor model, and concluded that no zigzag behavior appeared in  $S(I)$  versus spin  $I$  plots of the axial nucleus. Brentano *et al.* [4] proposed the soft rotor energy formula (SRF) for the ground band and Bihari *et al.* [5] used this SRF to calculate the perturbed energies of  $\gamma$ -band produced by rotation of rigid asymmetric atomic nucleus.

The purpose of the present study is to search the nature of multiphonon 2 $\gamma$ -band of  $^{158}\text{Dy}$  nucleus. In  $^{158}\text{Dy}$ , one phonon  $\gamma$  and two phonon 2 $\gamma$ -band have been observed (see Ref.6).

### Method and Calculations

The SRF formula [4] has been used

$$E = \frac{I(I+1)}{J_0(1+\alpha I)}$$

where  $J_0$  is the moment of inertia parameter and  $\alpha$  is the variable of moment of inertia parameter. The values of  $J_0$  and  $\alpha$  are calculated in the same way as done earlier by Bihari *et al.* [5] for  $\gamma$ -band. In the present paper, the energy values are calculated by fitting the odd and even spin energy sequences, using  $2_3^+$  and  $4_3^+$  energies and  $3_2^+$  and  $5_2^+$  energies for 2 $\gamma$ -band. Staggering indices [7] is defined as

$$S(I, I-1, I-2) = \frac{(E_1 - E_{I-1}) - (E_{I-1} - E_{I-2})}{E_1^2} \quad (1),$$

which shows alternating behavior with spin  $I$ .

### Results and Discussion

The energy levels for ground,  $\gamma$  and 2 $\gamma$ -bands in  $^{158}\text{Dy}$  nucleus are plotted in Fig.1.

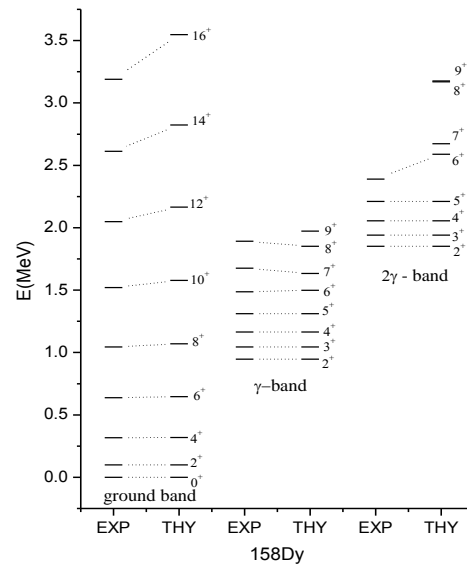


Fig.1: Comparison between experimental and calculated energy values of the ground,  $\gamma$  and 2 $\gamma$ -band using SRF formula.

For the ground band, calculated energy values show good agreement with the experimental energy up to  $I=12^+$ , where as in case of  $\gamma$ -band, the calculated energies show very good agreement with the experimental energies. In case of 2 $\gamma$ -band, the calculated energy values match excellently with the experimental energy values up to  $I=5^+$ , while a small deviation in the calculated values of energy possessing angular momentum  $I > 5^+$  have been observed.

The calculated values of  $J_0$  and  $\alpha$  for ground (g), even and odd spin sequences for  $\gamma$

and  $2\gamma$ -bands are listed in Table 1. For  $^{158}\text{Dy}$  nucleus, the sign of  $J_0$  and  $\alpha$  are positive for ground band energies and negative for  $\gamma$  and  $2\gamma$ -band energies.

Band	$\alpha_g$	$\alpha_{even}$	$\alpha_{odd}$	$J_0$ (g)	$J_0$ (even)	$J_0$ (odd)
g	58.08			0.02		
$\gamma$		-1.20	-1.01		-4.50	-5.59
$2\gamma$		-0.99	-0.75		-3.25	-4.89

Table 1: Fitted parameters  $J_0$  and  $\alpha$  used in present work.

Davydov-Filippov energy gap relations [3]  $\Delta E1 [= E3_1^+ - (E2_1^+ + E2_2^+)] = 0$  and  $\Delta E2 [= E3_1^+ - (2E2_1^+ + E4_1^+)]$  has been used to distinguish the rigid rotor from  $\gamma$ -soft rotor. For  $\gamma$ -rigid rotor,  $\Delta E1 \ll \Delta E2$  has been used [5] and found valid in the present study for  $\gamma$ -band. As studied earlier by Bihari *et al.* [8], it was also found valid.

For  $2\gamma$ -band, in order to distinguish the  $\gamma$ -rigid rotor, the new parameters  $\Delta E1 [= E3_2^+ - E22^{++} + E23^+]$  and  $\Delta E2 [= E3_2^+ - (2E2_2^+ + E4_2^+)]$  have been introduced. The parameter,  $\Delta E1 \ll \Delta E2$  as used to distinguish the  $\gamma$ -rigid and  $\gamma$ -soft nuclei has not been found valid for  $2\gamma$ -band as shown in Table 2.

For $\gamma$ -band $^{158}\text{Dy}$	$\Delta E1 [= E3_1^+ - E21^{++} + E22^+]$	$\Delta E2 [= E3_1^+ - 2E21^{++} + E41^+]$
	0.0	0.530
For $2\gamma$ -band $^{158}\text{Dy}$	$\Delta E1 [= E3_2^+ - E22^{++} + E23^+]$	$\Delta E2 [= E3_2^+ - 2E22^{++} + E42^+]$
	-0.858	-2.007

Table 2: Experimental values of energy differences  $\Delta E1$  and  $\Delta E2$  for  $\gamma$  and  $2\gamma$ -bands.

From the above evidences, it has been observed that this nucleus doesn't obey Davydov-Filippov energy gap relation for  $2\gamma$ -band. Hence, it doesn't show triaxial region in multiphonon  $2\gamma$ -band.

The variation of staggering indices  $S(I)$  with spin ( $I$ ) for  $^{158}\text{Dy}$  nuclei have been observed (see Fig.2), the spacing between odd-even spin levels in the present work show agreement with the experimental values for  $S(4)$ ,  $S(5)$  and the calculated values for  $S(6)$ ,  $S(7)$ ,  $S(8)$  and  $S(9)$  does not match with experimental values. For  $2\gamma$ -band, the experimental values for  $S(I)$  does not changes sign rather have positive sign with small values.

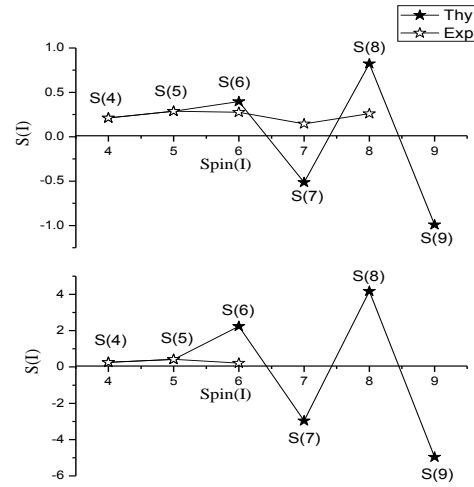


Fig.2: Staggering indices  $S(I)$  is plotted versus spin ( $I$ ) in the present work and experiment for  $^{158}\text{Dy}$  nuclei using Eq.(1) for  $\gamma$  and  $2\gamma$ -bands.

For axial rotor, all the  $S(I)$  are positive and small in magnitude and show no zigzag behavior, but increases slowly with increasing spin  $I$  [2]. In  $^{158}\text{Dy}$  nucleus, experimental values of  $S(I)$  have small positive values for all spins of  $2\gamma$ -band. Here, it appears that at very low spin  $^{158}\text{Dy}$  nucleus is axial but at higher spin ( $I^+ = 7^+$ ), it may become unpredictable. Hence, we need more experimental energy values in multiphonon  $2\gamma$ -band to study the nature of this nucleus.

## References

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