

Description and Prediction of Long -Lived Isomers in Neptunium (Z=93) Isotopes

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Long-lived isomer pairs have been identified/indicated [1] for well over 50 years in the odd-odd $^{236-244}\text{Np}$ with the exception of ^{238}Np isotope. However, firm two-quasiparticle (2qp) configuration assignments, or even the relative energy placement, for any pair of these isomers remains undefined to-date. Using the rotor-particle model with the inclusion of residual n-p interaction contribution, Sood *et al.* [2] had developed a formalism, which has been extensively employed for credible description/prediction of 2qp structures of deformed odd-odd nuclei, both of the actinide [2,3] and the rare earth [4] regions. In a simplified semi-empirical version of this approach [4,5], the energy of a specific 2qp(Ω_p, Ω_n) state of an odd-odd deformed nucleus is written as follows:

$$E(K; \Omega_p, \Omega_n) = E_0 + E_p(\Omega_p) + E_n(\Omega_n) + E_{rot} + \langle V_{pn} \rangle,$$

$$E_{rot} \approx -\frac{\hbar^2}{2I} (2\Omega_{<}) \delta_{K,K-};$$

$$\langle V_{pn} \rangle = -\left(\frac{1}{2} - \delta_{\Sigma,0}\right) E_{GM} + (-)^J E_N \delta_{K,0} \dots (1)$$

wherein, E_p/E_n are the experimental energies of neighboring odd-A isotope/isotone, and the GM doublet splitting energy E_{GM} and Newby parameter E_N for $K=0$ bands are the values of these quantities for respective band as observed in any odd-odd neighbor. Following the 3-step procedure adopted in our recent publications [5], we have carried out detailed analysis for $^{240,242,244}\text{Np}$ isotopes as summarized below.

As our first step, we map the 1qp configuration space by plotting, in Fig.1, the experimental E_x in the core ^{239}Np nucleus for proton (on right) and in ^{243}Pu for neutron orbitals (on left) of interest in n-rich odd-odd Np isotopes. The circled numbers denote large energy gaps, which give rise to isomer pairs in Np isotopes having A values shown in square boxes on connecting lines in Fig. 1.

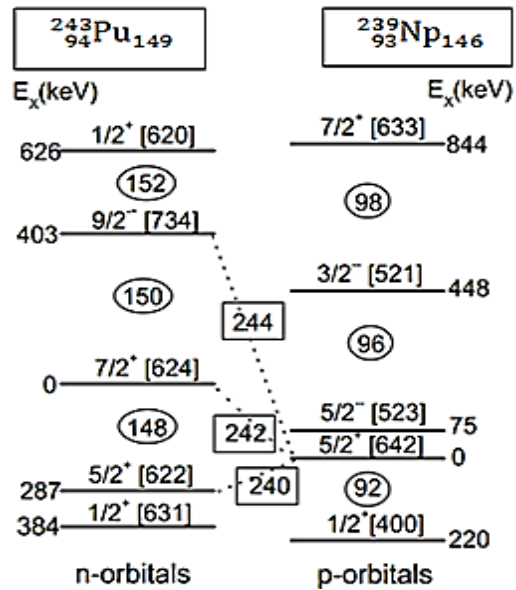


Fig. 1: Schematic plot of p/n orbitals of interest for n-rich odd-odd $_{93}\text{Np}$ level structures. The orbital excitation energies E_x are experimental [1] values for ^{239}Np (on the right) and ^{243}Pu (on the left). Circled numbers indicate significant shell gaps. Numbers in square boxes denote A for the respective odd-odd Np isotope.

In the second step, we use the data in Fig.1 to enumerate the $K^\pm(\Omega_p \pm \Omega_n)$ doublets expected around the Fermi surface for each isotope. Table 1 lists the lowest 2qp structures, denoted by spins-triplet K_T (on left) and spins-singlet K_S (on right) following GM rule in each case. The entries herein are limited to summed $(E_p + E_n) < 150$ keV in each case, since we are presently interested only in isomers close to Fermi surface. As evident herein, the lowest two bands in each case have $\Delta K=5$, thus resulting in long-lived isomers having no significant IT/ γ connection and hence no precise energy values.

Table 1: Low lying ($E_x < 200$ keV) pairs of GM doublet bands expected in $^{240,242,244}\text{Np}$ nuclei. The entries in each box list respectively K^π for the spins parallel (K_T) and spins antiparallel (K_S) band from the indicated 2qp configuration.

n↓ \ p→	p ₀		p ₁	
	5/2 ⁺ [642↑]		5/2 ⁻ [523↓]	
147 5/2 ⁺ [622↑]	5 ⁺	0 ⁺	0 ⁻	5 ⁻
149 7/2 ⁺ [624↓]	1 ⁺	6 ⁺	6 ⁻	1 ⁻
151 9/2 ⁻ [734↑]	7 ⁻	2 ⁻	2 ⁺	7 ⁺

As an illustration, we briefly discuss here our results for ^{240}Np . Out of the 4 bands for ^{240}Np , listed in row 1 of Table 1, the 5⁺ band appears [6] as ^{240}Np g.s., while the other 3 bands are observed in the isotonic neighbor ^{242}Am [1]. Using these data in eq (1), the evaluated energies of these 4 bandheads and also J=1 level of K=0 band are shown in Fig.2. Our analysis confirms $\{\pi 5/2[642] + \nu 5/2[622]\}$ configuration, which is listed as ‘probable’ in latest NDS [6], for the 62m ^{240}Np high-spin isomer (HSI). For the 7.22m low-spin isomer [LSI], NDS [6] lists $1^+0\{\pi 7/2[633]-\nu 7/2[624]\}$ configuration. Hseuh *et al.* [7] had proposed this configuration, albeit with many ifs and buts, by considering just one β -branch from 7.22m ^{240}Np decay to 597 keV $J^\pi K=1^-0$ level in ^{240}Pu and a 1976 model (QPM) based 2qp decomposition of this octupole state. However, they did point out that E_x of suggested p-orbital is too high (>500 keV) to constitute a particle-hole structure for this vibrational level. Even the β -branches $^{240}\text{U}(0^+) \rightarrow ^{240}\text{Np}(1^+)$ and $^{240}\text{Np}(1^+) \rightarrow ^{240}\text{Pu}(0^+)$ were not considered to cross check this assignment.

As seen in Fig.1, the $\pi 7/2[633]$ 1qp state is experimentally [1] indicated at 844 keV in ^{239}Np and hence cannot form a constituent of LSI close to ^{240}Np g.s. Further with $Q_\beta(^{240}\text{U}) \sim 400$ keV, such a high-lying state cannot be populated in this β -decay. These, and other equally strong, arguments clearly establish

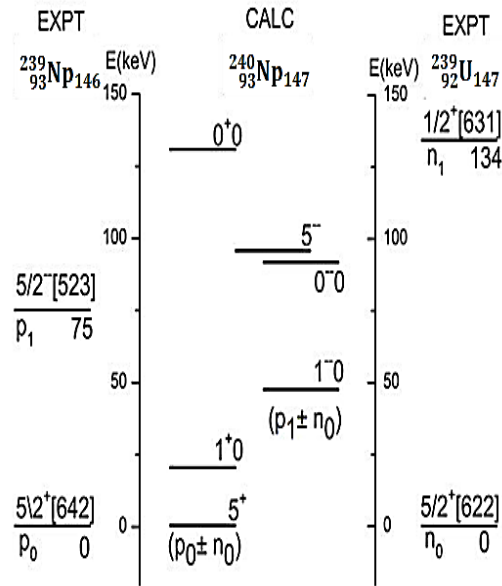


Fig. 2: Plot of experimental 1qp bandhead energies in isotopic ^{239}Np (on the left) and in isotonic ^{239}U (on the right), and model calculated 2qp level energies in ^{240}Np (in the middle).

$J^\pi K=1^+0\{\pi 5/2[642]-\nu 5/2[622]\}$ --- (2) as the only physically acceptable configuration for the 7.22m ^{240}Np isomer.

Similar analysis and critical examination of isomer pairs in other odd-odd Np isomers form a part of this study, and will be reported.

References

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