

MICROSCOPIC STUDY ON STRUCTURE OF YRAST BANDS IN EVEN-EVEN ⁹⁰⁻¹⁰⁰Ru ISOTOPES

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Introduction

The research in nuclear structure physics reached a new height with the advancement in the development of accelerators and efficient γ -ray detectors. This made it possible to carry out nuclear reactions where in nuclei could be produced in highly excited states, which in turn decayed to their ground state by emitting a cascade of γ -rays. Fusion-evaporation reactions provided the standard mechanism to populate states with high angular momentum.

In recent years, Projected Shell Model has become quite successful in explaining a broad range of properties of deformed nuclei [1]. The advantage in this method is that the numerical requirements are minimal and therefore, it is possible to perform a systematic study for a group of nuclei in a reasonable time frame. A systematic study of the neutron-rich nuclei in the mass region $A=100$ using PSM has been carried out recently and the agreement between PSM results and experimental data has been found to be quite good [2,3].

The Hamiltonian which has been used throughout the present work is described as follows.

$$\hat{H} = H_0 - \frac{1}{2} \chi \sum_{\mu} \hat{Q}_{\mu}^{\dagger} \hat{Q}_{\mu} - G_M \hat{P}^{\dagger} \hat{P} - G_Q \sum_{\mu} \hat{P}_{\mu}^{\dagger} \hat{P}_{\mu}$$

where H_0 is the spherical single particle Hamiltonian. The second term is the quadrupole-quadrupole interaction and the last two terms are the monopole and quadrupole pairing interactions, respectively. The strength of the quadrupole-quadrupole term is obtained via self-consistent conditions with a given deformation parameter ϵ_2 .

A microscopic description of yrast states, back-bending phenomena, B(E2) transition

probabilities and g-factors has been carried out in Ru isotopes with $A=90-100$ studied in the framework of PSM. The prolate structure of the bands has been investigated. It is observed that the proton sub-shells $2p_{1/2}$, $2p_{3/2}$, $1f_{5/2}$ and $1f_{7/2}$ are polarized. The maximum polarization is suffered by $1f_{5/2}$ and $2p_{1/2}$ proton sub-shells. The variation in the occupation numbers of $2p_{1/2}$ proton sub-shells from isotope to isotope is very very small. In fact, for most of the isotopes, the $2p_{1/2}$ occupation is around 1.2 and $1f_{5/2}$ is around 5. This type of variation for $2p_{1/2}$, $1f_{5/2}$ sub shells can be linked to producing a small amount of background collectivity in all Ru isotopes. It is therefore, very unlikely that these core polarization effects of $2p_{1/2}$ and $1f_{5/2}$ proton sub-shells could be responsible for the observed low lying systematics of Ru isotopes. The occupation probability of $2d_{5/2}$ neutron orbit is found to increase from 0.8 in ⁹⁰Ru to 2.21 in ¹⁰⁰Ru. This means that $k=1/2, 3/2$ components of $2d_{5/2}$ neutron orbit are getting occupied. Since these orbits are down sloping, their occupation will produce increase of deformation. Thus, this factor could cause increase of deformation as one proceeds from ⁹⁰Ru to ¹⁰⁰Ru. The observed deformation systematics in ⁹⁰⁻¹⁰⁰Ru isotopes depends on the increase of occupation probability of $(1h_{11/2})_v$ orbit.

References:

- [1] K. Hara and Y. Sun, Int. J. Mod. Phys. **E4**, 637 (1995).
- [2] Arvind Bhat, Arun Bharti and S.K. Khosa IJMP E, Vol 21, No.3(2012) 1250030
- [3] Arvind Bhat, Arun Bharti and S.K. Khosa Eur.Phys.J.A (2012) 48:39

Table1.1: the BCS sub shell occupation numbers of protons in the ground states of $^{90-112}\text{Ru}$

Isotope	$2s_{1/2}$	$1d_{3/2}$	$1d_{5/2}$	$2p_{1/2}$	$2p_{3/2}$	$1f_{5/2}$	$1f_{7/2}$	$3s_{1/2}$	$2d_{3/2}$	$2d_{5/2}$	$1g_{7/2}$	$1g_{9/2}$
^{90}Ru	2	3.99	6	1.2	3.75	4.96	7.94	0.04	0.05	0.52	0.18	5.35
^{92}Ru	1.97	3.99	6	1.28	3.78	5.08	7.94	0.03	0.04	0.42	0.14	5.3
^{96}Ru	2	3.99	6	1.28	3.79	5.1	7.94	0.03	0.04	0.42	0.13	5.29
^{98}Ru	2	3.99	6	1.25	3.78	5.06	7.94	0.03	0.04	0.46	0.14	5.96
^{100}Ru	2	3.99	6	1.27	3.79	5.11	7.94	0.03	0.029	0.42	0.13	5.29

Table1.2: The BCS sub shell occupation numbers of neutrons in the ground states of $^{90-112}\text{Ru}$

Isotope	$2p_{1/2}$	$2p_{3/2}$	$1f_{5/2}$	$1f_{7/2}$	$3s_{1/2}$	$2d_{3/2}$	$2d_{5/2}$	$1g_{7/2}$	$1g_{9/2}$	$3p_{1/2}$	$3p_{3/2}$	$2f_{5/2}$	$2f_{7/2}$	$1h_{9/2}$	$1h_{11/2}$
^{90}Ru	1.59	3.8	5.2	7.9	0.12	0.17	0.8	0.4	5.66	0.005	0.013	0.016	0.051	0.04	0.24
^{92}Ru	1.84	3.9	5.66	7.95	0.12	0.17	0.8	0.42	6.79	0.004	0.009	0.013	0.044	0.035	0.47
^{96}Ru	1.96	3.95	5.87	7.96	0.3	0.42	1.41	1.09	8.47	0.004	0.014	0.018	0.071	0.041	0.83
^{98}Ru	1.96	3.96	5.89	7.96	0.37	0.57	1.78	1.65	8.89	0.005	0.015	0.019	0.107	0.052	0.74
^{100}Ru	1.97	3.97	5.92	7.96	0.4	0.65	2.21	2.22	9.33	0.006	0.016	0.021	0.147	0.064	1.1