

## Synthesis of Z=108 <sup>269–271</sup>Hs nuclei: Entrance channel effects and role of double-magicity of <sup>48</sup>Ca beam

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### Introduction

The macro-microscopic model, as well as self-consistent mean-field calculations that also consider deformed nuclear shapes, predicted deformed shell closures at Z=108 and N=162 [1], creating a region of enhanced stability halfway between the heaviest nuclides found on earth and the predicted island of stability around Z=114, 120 or 126, N=172 or 184. These predictions are now confirmed [2].

Z=108 Hassium nucleus was first studied in the reaction <sup>248</sup>Cm(<sup>26</sup>Mg,4n/5n)<sup>270,269</sup>Hs [3] at GSI (Darmstadt) through 4n and 5n decay channels, which was later [4] extended also to 3n exit channel. A similar complete measurement of evaporation residues was later reported [5] for the fusion reaction <sup>238</sup>U(<sup>36</sup>S,5-3n)<sup>269–271</sup>Hs. More recently, <sup>226</sup>Ra+<sup>48</sup>Ca was used to synthesize the doubly-deformed magic nucleus <sup>270</sup>Hs [2]. The suitability of using <sup>48</sup>Ca projectile for the synthesis of superheavy elements (SHEs) has been shown repeatedly in the production of Z=112-118 nuclei.

Here, we concentrate on the entrance channels <sup>226</sup>Ra+<sup>48</sup>Ca and <sup>248</sup>Cm+<sup>26</sup>Mg [2–4], for the synthesis of Hs nucleus by individual light-particle (n=3,4,5) decay channels, on the basis of the Dynamical Cluster-decay Model (DCM) of Gupta and collaborators (see, e.g., [6, 7] and earlier references therein). In DCM, the deformation effects are included upto  $\beta_2$ , with the ‘hot’ optimum orientations. The DCM is a one parameter model, and, in this study, we have tried to analyze the significance of using <sup>48</sup>Ca as a projectile for synthesizing SHEs.

### The Model

In DCM, the compound nucleus (CN) decay cross-section in terms of partial waves is

$$\sigma = \frac{\pi}{k^2} \sum_{\ell=0}^{\ell_{max}} (2\ell+1)P_0P; \quad k = \sqrt{\frac{2\mu E_{cm}}{\hbar^2}} \quad (1)$$

where,  $\mu = [A_L A_H / (A_L + A_H)]m$  is the reduced mass and  $E_{c.m.}$ , the center of mass energy.  $\ell_{max}$ , the maximum angular momentum, is fixed for the light particles (LPs) cross-section  $\sigma_{LPs} \rightarrow 0$ . The preformation probability  $P_0$  is the solution of stationary Schrödinger equation in  $\eta = (A_H - A_L)/(A_H + A_L)$ , such that  $P_0(A_i) \propto |\psi(\eta(A_i))|^2$ . It contains the structure information of compound nucleus via the fragmentation potential,

$$V(R, \eta, T) = \sum_{i=1}^2 [V_{LDM}(A_i, Z_i, T)] + \sum_{i=1}^2 [\delta U_i] \exp(-\frac{T^2}{T_0^2}) + V_C(R, Z_i, \beta_{\lambda_i}, \theta_i, T) + V_P(R, A_i, \beta_{\lambda_i}, \theta_i, T) + V_l(R, A_i, \beta_{\lambda_i}, \theta_i, T) \quad (2)$$

used in stationary Schrödinger equation. Here,  $V_{LDM}$  is the T-dependent liquid drop energy and  $\delta U$ , the ‘‘empirical’’ shell corrections. The T-,  $\beta_{\lambda_i}$ - and  $\theta_i$ -dependent proximity  $V_P$ , Coulomb  $V_C$  and angular momentum-dependent potential  $V_l$  are [6]

$$V_P(s_0(T)) = 4\pi \bar{R}(T) \gamma b(T) \Phi(s_0(T)) \quad (3)$$

$$V_C = \frac{Z_1 Z_2 e^2}{R} + 3Z_1 Z_2 e^2 \sum_{\lambda, i=1,2} \frac{R_i^\lambda(\alpha_i, T)}{(2\lambda+1)R^{\lambda+1}} \times Y_\lambda^{(0)}(\theta_i) \left[ \beta_{\lambda_i} + \frac{4}{7} \beta_{\lambda_i}^2 Y_\lambda^{(0)}(\theta_i) \right], \quad (4)$$

$$V_l = \frac{\hbar^2 \ell(\ell+1)}{2I_s}. \quad (5)$$

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**Table 1:** Comparison of experimental and calculated  $\sigma_{xn}$ ,  $x=3-5$ , with fitted  $\Delta R$  at  $E^* \approx 40$  MeV.

Target-projectile combination used	No. of neutrons emitted	Cross-section (pb)		
		Experimental	Calculated	$\Delta R$ (fm)
$^{248}\text{Cm}+^{26}\text{Mg}$	3n	$1.38^{+1.75}_{-0.85}$	1.41	1.374
	4n	$3.0^{+2}_{-1.5}$	3.07	1.642
	5n	$0.72^{+1.54}_{-0.39}$	0.729	1.836
$^{226}\text{Ra}+^{48}\text{Ca}$	4n	$16^{+13}_{-7}$	16.4	1.618

with the shortest distance  $s_0$  giving compact  $\theta_{ci}$ .  $P$  in Eq. (2) is the WKB integral, with first turning point  $R_a(\eta, T) = R_1 + R_2 + \Delta R(T)$ ;  $R_i$  are radius vectors of the two nuclei and  $\Delta R(T)$ , a parameter that assimilates the neck formation effects, constant for all the fragments at a given excitation energy.

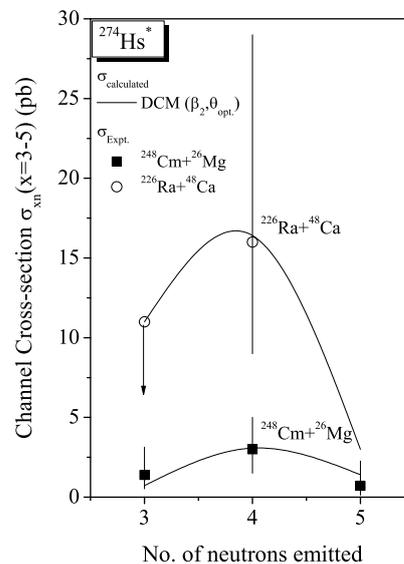
### Calculations and Results

We consider  $^{248}\text{Cm}+^{26}\text{Mg}$  and  $^{226}\text{Ra}+^{48}\text{Ca}$  hot-fusion reactions at  $E^*=40$  MeV, where  $3n$ ,  $4n$  and  $5n$  emission data are from [2–4]. Note, the two target-projectile (t,p) combinations are, respectively, the deformed+deformed and deformed+spherical nuclei, forming the same CN  $^{274}\text{Hs}^*$ .

Table 1 and Fig. 1 show a comparison between the experimental and DCM-calculated  $\sigma_{xn}$ ,  $x=3-5n$ , from  $^{274}\text{Hs}^*$  at a fixed  $E^*=40$  MeV for both the entrance channels. Clearly, the experimental data are reproduced within one parameter  $\Delta R$  fitting. Knowing that asymmetric t-p combinations are more suitable for the synthesis of new/ SHEs (lower interaction barrier and larger interaction radius [8]),  $^{248}\text{Cm}+^{26}\text{Mg}$  should give a higher cross-section at any given  $E^*$ , compared to that for  $^{226}\text{Ra}+^{48}\text{Ca}$ . However, Fig. 1 shows that, contrary to our expectations,  $^{248}\text{Cm}+^{26}\text{Mg}$  has lower  $\sigma_{xn}$ , than  $^{226}\text{Ra}+^{48}\text{Ca}$ . Apparently, this is because of the double magicity of spherical  $^{48}\text{Ca}$  nucleus. Also, out of three isotopes produced,  $^{270}\text{Hs}$  has the highest cross-section for the apparent reason of being a double deformed-magic nucleus.

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**FIG. 1:** The experimental and DCM calculated  $xn$  cross-sections for two different entrance channels.

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