

## Nuclear structure of multiphonon $\gamma\gamma$ -band in neutron rich $^{112}\text{Ru}$ nucleus

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### Introduction

The neutron rich  $^{112}\text{Ru}$  nucleus ( $Z=44$  and  $N=68$ ) lie within the  $A=100$  deformed mass region. The experimental study of multiphonon bands was very difficult but now-a-days there are many methods to measure them. Recently, Xing-Lai et al.,[1] investigate the neutron-rich  $^{112}\text{Ru}$  nucleus with the gamma-sphere detector array by observing prompt  $\gamma$ -rays of spontaneous fission of  $^{252}\text{Cf}$ . They proposed two-side bands, one of them was predicted as two-phonon  $\gamma$ -vibrational band and the another one as two-quasiparticle band. The energy ratio  $R_{4/2} = \frac{E(4_1^+)}{E(2_1^+)}$  of the neutron rich  $^{112}\text{Ru}$  nucleus is 2.72, this shows that it lie in deformed region. Using cracked shell model, they predicted that in the ground state, the  $^{112}\text{Ru}$  nucleus has oblate shape deformation and posses triaxial deformation with increasing rotational frequency. Bohr and Mottelson [2] stated that at  $\gamma \geq 24^\circ$ , the nuclei believed to take any shape, including triaxial. In Ref.[3], a number of  $\gamma$ -softness and  $\gamma$ -rigidity signatures in various nuclei has been reviewed and gave most of its attention to the staggering properties of  $\gamma$ -band energies.

In the present work, we study the nature of multiphonon  $\gamma\gamma$ -band in neutron rich  $^{112}\text{Ru}$  nucleus and also calculate the energy value of one-phonon  $\gamma$ -band ( $K=2$ ) and two-phonon  $\gamma\gamma$ -band ( $K=4$ ) by using Modified Soft Rotor Formula (proposed by Gupta et al.,[4]).

### Method and Calculations

Brentano et al.,[5] proposed the soft rotor energy formula (SRF) for the ground band and later, Bihari et al.,[6] used this SRF to calculate the energy of one phonon  $\gamma$ -band. They

received both the positive and negative values of moment of inertia (MoI)  $\theta_0$  and also for softness parameter  $\sigma$  in Ru isotopes and also in many other nuclei. Recently, Gupta et al.,[4] illustrated that it is difficult to justify the negative values of MoI and also the large values of  $\sigma$ . As the softness parameter is only a perturbation correction of MoI [5], so  $\sigma$  is expected to be less than one and should be positive. Gupta et al.,[4] resolved the anomaly of negative MoI and the negative softness parameter  $\sigma$  and also calculate the energy of one phonon  $\gamma$ -band of deformed and shape transition nuclei.

The Modified Soft Rotor Formula (MSRF) is given as:

$$E(I) = EK + \frac{I(I+1)}{2\theta_0(1+\sigma I)}$$

where  $\theta_0$  is the MoI parameter and  $\sigma$  is the variable of MoI parameter. For the detail explanation for the calculation of the energy values see Ref.[4]. Here MSRF is true for the corresponding value of  $\gamma\gamma$ -band MoI =  $\frac{3}{E(5\gamma\gamma)-E(4\gamma\gamma)}$ , for  $\gamma$ -band MoI =  $\frac{3}{E(3\gamma)-E(2\gamma)}$  and ground MoI  $\frac{3}{E(2_1^+)}$ .

The staggering indices [3] given as:

$$S(I, I-1, I-2) = \frac{(E_I - E_{I-1}) - (E_{I-1} - E_{I-2})}{E(2_1^+)}$$

shows alternative behaviour with spin I. In case of  $\gamma$ -rigid triaxial, the clustering of the  $\gamma$ -band energy levels is predicted, which resulting in an oscillating behaviour of S(I) such that it is negative for odd-spin and positive for even spin levels.

### Results and Discussions

The energy levels for ground,  $\gamma$  and  $\gamma\gamma$ -bands in the neutron rich  $^{112}\text{Ru}$  nucleus are

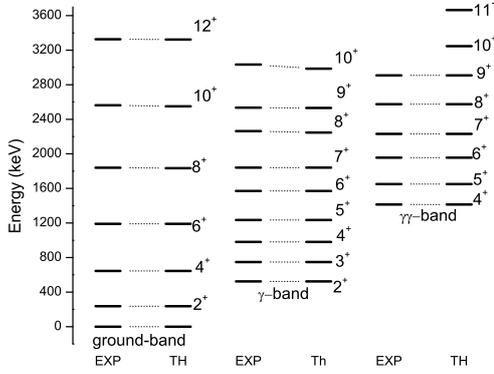


FIG. 1: Comparison between experimental and calculated values of the ground,  $\gamma$  and  $\gamma\gamma$ -bands energy using MSRF formula.

$^{112}\text{Ru}$ nucleus	Theoretical MoI	Expt. MoI	$\sigma$
Ground band	$\frac{3}{E(2_1^+)} = 0.012$	0.012	0.075
$\gamma$ -band	$\frac{3}{(E(3\gamma) - E(2\gamma))} = 0.013$	0.0102	0.099
$\gamma\gamma$ -band	$\frac{3}{(E(5\gamma\gamma) - E(4\gamma\gamma))} = 0.013$	0.0106	0.11

TABLE I: The softness parameter  $\sigma$ , MoI  $\theta_{\text{ground-band}}$ ,  $\theta_{\gamma\text{-band}}$  and  $\theta_{\gamma\gamma\text{-band}}$  ( $\text{keV}^{-1}$ ) from MSRF. The rotor model  $\theta_{\text{grnd}} = \frac{3}{E(2_1^+)}$  for ground band,  $3/(E(3\gamma) - E(2\gamma))$  for  $\gamma$ -band and  $3/(E(5\gamma\gamma) - E(4\gamma\gamma))$  for  $\gamma\gamma$ -band are listed for comparison.

plotted in Fig.1. The calculated energy values match excellently with the experimental energy for all spin values for ground,  $\gamma$  and  $\gamma\gamma$ -bands. The experimental data is taken from Ref.[1].

The calculated values of  $\theta_0$  and  $\sigma$  for ground,  $\gamma$  and  $\gamma\gamma$ -bands are listed in Table 1. In the neutron rich  $^{112}\text{Ru}$  nucleus, the calculated  $\theta_0$  for  $\gamma$  and  $\gamma\gamma$ -bands are almost equal to the calculated MoI for ground band, which is close to the corresponding rotor model values.

For axial rotor, all the staggering indices  $S(I)$

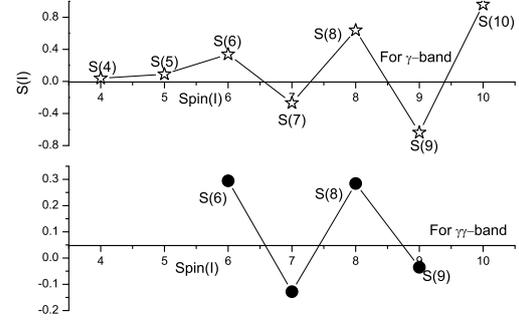


FIG. 2: Staggering indices  $S(I)$  is plotted versus spin  $I$  using experiment values for  $^{112}\text{Ru}$  nucleus for  $\gamma$  and  $\gamma\gamma$ -bands.

values are positive and increases slowly with increasing spin  $I$  and show no zigzag behaviour [8]. The  $^{112}\text{Ru}$  nucleus develop a staggering pattern, here the experimental values of  $S(I)$  are positive for even spin values and negative for odd spin values for  $\gamma$ -band. The  $^{112}\text{Ru}$  nucleus is predicted as  $\gamma$ -rigid triaxial nucleus in  $\gamma$ -band(see Ref.[7]). In case of  $\gamma\gamma$ -band,  $^{112}\text{Ru}$  nucleus show the similar alternating behaviour (see Fig.2). Hence, it is recommended that the neutron rich  $^{112}\text{Ru}$  nucleus is  $\gamma$ -rigid triaxial in nature in multiphonon  $\gamma\gamma$ -band.

## References

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