

## Systematic dependence of product $(E(2_{\gamma}^+) * B(E2) \uparrow)$ on asymmetry parameter $\gamma_0$

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### Introduction

Davydov and Filippov [1] introduced the Asymmetric Rotor Model (ARM) to determine the energy level spacing and transition probabilities of excited states in few even-even nuclei. The Hamiltonian operator  $H$  used in ARM is given by:

$$H = \frac{1}{2} \sum J_{\lambda}^2 \text{Sin}^{-1}(\gamma - \lambda \frac{2\pi}{3})$$

where  $J_{\lambda}$  are the projections of the operators of the total angular momentum along the axes of coordinate system fixed in the nucleus. In mass region of deformed nuclei,  $150 \leq A \leq 190$  and  $A \geq 222$ , the theoretical calculations were found to be in good agreement with experimental values. In mass region of spherical nuclei,  $70 \leq A \leq 130$ , the ARM model also showed good agreement with experimental data for some nuclei. The asymmetry parameter  $\gamma_0$  of ARM varies between 0 and  $\pi/3$  and it mainly determines the deviation of shape of the nucleus from axial symmetry. In ARM model, the  $B(E2) \uparrow$  value is related with  $\gamma_0$

$$B(E2, 2_1^+ \rightarrow 0_1^+) = \frac{1}{2} (1 + \frac{3-2\text{sin}^2(3\gamma_0)}{(9-8\text{sin}^2(3\gamma_0))^{1/2}})$$

in units of  $\frac{e^2 Q_0^2}{16\pi}$ .

The reduced excitation strength,  $B(E2) \uparrow$  value is also of great interest in nuclear physics, as the collective effects in low-lying energy states are quadrupole in nature. Recently, a systematic study of  $B(E2) \uparrow$  value with asymmetry parameter  $\gamma_0$  have been presented by R. Kumar et al., [2]. Earlier, Gupta and Sharma [3] and Mittal et al., [4] gave correlation between the  $B(E2)$  ratio and the asymmetry parameter  $\gamma_0$  for medium and light nuclei. Grodzins [5] illustrates a close

relationship between energy of first excited state  $E(2_1^+)$  and reduced excitation strength,  $B(E2) \uparrow$  values. In nuclear energy spectra, with increasing the valence neutron and proton pairs, the collectivity also increases, which yields an decrease in  $E(2_1^+)$  and increase in  $B(E2) \uparrow$ , Grodzins gave this relation

$$(E(2_1^+) * B(E2) \uparrow) \sim \text{constant}(Z^2/A)$$

Recently, Gupta [6] concluded that the constancy of Grodzins product rule breaks down in the combined effect of the  $Z=64$  subshell effect and the shape transition.

In the present work, we replace the first excited state energy  $E(2_1^+)$  with second excited  $E(2_{\gamma}^+)$  energy state, and now the equation becomes

$$(E(2_{\gamma}^+) * B(E2) \uparrow)$$

and study the systematics of the product  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  with asymmetry parameter  $\gamma_0$  for the first time. This study provides a new insight in nuclear physics. The values of  $B(E2) \uparrow$  are taken from Ref. [7].

### Method of determining $\gamma_0$

There are many methods to find  $\gamma_0$ , but the most relevant way [1] is by using  $R_{\gamma}(= E_{2\gamma}/E_{2g})$  in equation:

$$\gamma_0 = \frac{1}{3} \text{Sin}^{-1}(\frac{9}{8}(1 - (\frac{R_{\gamma}-1}{R_{\gamma}+1})^2))^{1/2}$$

The values of  $E_{2g}$  and  $E_{2\gamma}$  are taken from National Nuclear Data Centre website [8].

### Results and discussion

Here, we use a simple rule of dividing the major shell space ( $Z=50-82$ ,  $N=82-126$ ) based on the hole and particle boson subshell space

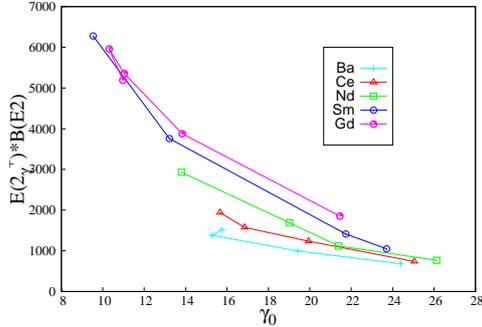


FIG. 1: The plot of energy  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  vs. asymmetry parameter  $(\gamma_0)$  for Ba-Gd nuclei.

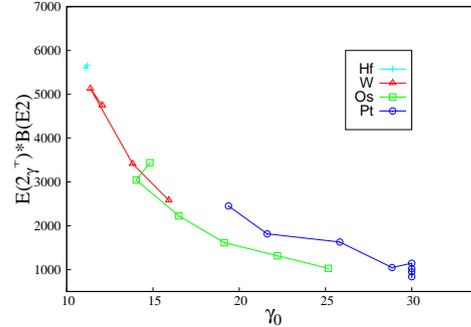


FIG. 3: The plot of  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  vs. asymmetry parameter  $(\gamma_0)$  for Hf-Pt nuclei.

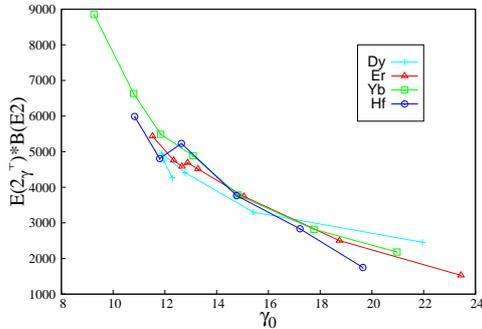


FIG. 2: The plot of  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  vs. asymmetry parameter  $(\gamma_0)$  for Dy-Hf nuclei.

given earlier by Gupta et al., [9]. The variation of  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  with asymmetry parameter  $\gamma_0$  are shown in Figs.1-3. In quadrant-I, the graph of  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  vs.  $\gamma_0$  is shown in Fig.1 for Ba-Gd nuclei. The plot shows monotonic fall of  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  with increasing  $\gamma_0$ , which reflects the smooth decrease of nuclear deformation. For quadrant-II, the plot is shown in Fig. 2 for nuclei Dy-Hf and it indicates that the  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  values show linear dependency on the asymmetry parameter  $\gamma_0$ . Fig.3 indicates that the value of  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  decreases with increasing value of asymmetry parameter  $\gamma_0$  for Hf-Pt nuclei.

### Conclusion

The product  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  provides a good measure of deformation in mass region

A=120-200. In three quadrants, the product  $(E(2_{\gamma}^+) * B(E2) \uparrow)$  shows systematic dependence on the asymmetry parameter  $\gamma_0$ .

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