

## Superdeformed rotational band in framework of three parameters rotational formula

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### Introduction

The discovery of superdeformation is one of the significant advances in nuclear structure physics. Superdeformed (SD) bands were first observed in fission isomers in the actinide region [1]. The discrete line SD shapes were found in  $^{152}\text{Dy}$  nucleus [2]. Since then vast experimental and theoretical studies have been undertaken. At present numerous SD bands have been observed in various mass region  $A=30,60,80,130,150$  and  $190$  [3-4].

Although a general understanding of SD bands has been achieved, there are still some striking features which remain partially understood. The excitation energy and firm assignment of spin-parity are not known for most of the SD bands because of near absence of information linking transitions between normal deformed (ND) and SD bands except in a few cases.

Several phenomenological formulae have been proposed to fit the transition energies and assign spin angular momentum to the observed levels in SD bands. The transition energies, spins n identical phenomenon for SD bands in the mass 150 and 190 regions have been predicted by various two and three parameters viz. variable moment of inertia,  $ab$  formula, Harris expansion [1-6].

In present work, we aim to describe the nuclear properties and SD bands of  $^{192}\text{Hg}$  (SD-1),  $^{194}\text{Hg}$ (SD-1) and  $^{194}\text{Hg}$ (SD-2) nuclei. We have employed the two-parameters  $ab$  formula and power law to calculate the transition energies of above mentioned bands.

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### Formulae used

The energies  $E(I)$  of the SD nuclear rotational bands as a function of the unknown spin ( $I$ ) can be expressed as

$$E(I) = E_0 + a \left[ \sqrt{1 + bI(I+1)} - 1 \right] + cI(I+1). \quad (1)$$

The rotational frequency  $\hbar\omega$ , defined as the derivative of energy  $E$ , with respect to the angular momentum  $J = I(I+1)$  is

$$\hbar\omega = \frac{dE}{dJ} = \left[ 2c + ab [1 + bI(I+1)]^{1/2} \right] (I(I+1))^{-1/2}. \quad (2)$$

The two types of possible nuclear moments of inertia have been suggested which reflects two different aspects of nuclear dynamics.

Kinematics moment of inertia:

$$J^1 = \hbar^2 I(I+1) (dEdI)^{-1} = \frac{\hbar^2}{ab} [1 + bI(I+1)]^{3/2} + \frac{1}{2c} \quad (3)$$

and the dynamic moment of inertia:

$$J^2 = \hbar^2 \left( \frac{d^2E}{dI^2} \right) = \frac{\hbar^2}{ab} [1 + bI(I+1)]^{3/2} + \frac{1}{2c} \quad (4)$$

Also the bandhead moment of inertia is

$$J_0 = \hbar^2 / (ab + 2c). \quad (5)$$

### A. Power Law

By replacing the concept of the arithmetic mean of the two terms used in Bohr-Mottelson expression by the geometric mean, Gupta [10] introduced a two-parameter formula called power law. In general the single-term power law is expressed as

$$E_I = aI^b. \quad (6)$$

## Results and Discussion

Theoretical energies so obtained for the SD-1 bands of  $^{192}\text{Hg}$  have been compared with the corresponding experimental values in Table 1 and Table 2. It shows that the theoretical transition energies obtained for SD band of  $^{194}\text{Hg}$  by using empirical formula, power law and *ab* formula show good agreement with the experimental values.

TABLE I: Comparison of theoretical (*ab* and Power law) and experimental results on transition energies of SD-1 band of  $^{192}\text{Hg}$  nuclei.

| Spin | $^{192}\text{Hg}(\text{SD1})$ | <i>ab</i> | Power Law |
|------|-------------------------------|-----------|-----------|
| 12   | 300.1                         | 300.1     | 301.2     |
| 14   | 340.4                         | 340.1     | 339.0     |
| 16   | 381.6                         | 382.5     | 384.4     |
| 18   | 421.1                         | 423.1     | 425.4     |
| 20   | 458.8                         | 455.6     | 459.1     |
| 22   | 496.0                         | 496.4     | 498.3     |
| 24   | 532.1                         | 531.5     | 534.3     |
| 26   | 567.4                         | 569.3     | 568.2     |
| 28   | 601.7                         | 604.5     | 603.2     |
| 30   | 634.9                         | 632.2     | 629.2     |
| 32   | 668.1                         | 666.2     | 667.2     |

TABLE II: Comparison of theoretical (*ab* and Power law) and experimental results on transition energies of SD-2 band of  $^{194}\text{Hg}$  nuclei.

| Spin | $^{194}\text{Hg}(\text{SD2})$ | <i>ab</i> | Power Law |
|------|-------------------------------|-----------|-----------|
| 12   | 283.1                         | 282.1     | 280.1     |
| 14   | 323.4                         | 324.0     | 325.0     |
| 16   | 363.1                         | 360.32    | 366.3     |
| 18   | 402.0                         | 401.2     | 405.2     |
| 20   | 440.3                         | 443.2     | 443.2     |
| 22   | 477.7                         | 480.1     | 479.1     |
| 24   | 514.2                         | 513.3     | 516.4     |
| 26   | 549.9                         | 547.3     | 548.3     |
| 28   | 584.9                         | 582.3     | 586.4     |
| 30   | 619.3                         | 616.3     | 623.5     |
| 32   | 652.3                         | 650.1     | 655.1     |

## 1. Conclusion

The present work provides a new insight to understand the nuclear structure of the SD-1 band of  $^{192}\text{Hg}$  and  $^{194}\text{Hg}$  nuclei. In the discussion above, we compared the four formulae: power law, *ab*, SRF and AHV. The power law and AHV show good accuracy in SD-1 and SD-2 bands.

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