

## On how $\langle P_{\text{CN}} \rangle$ behaves with mass-asymmetry in heavy nuclei

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The heavy evaporation residue (ER) production cross-section can be factorised as

$$\sigma_{\text{ER}} = \sum_{J=0}^{J_{\text{max}}} \sigma_{\text{cap}} P_{\text{CN}} W_{\text{sur}} \quad (1)$$

where,  $\sigma_{\text{cap}}$  is the capture cross section at center of mass energy ( $E_{\text{c.m.}}$ ) and spin  $J$ ,  $P_{\text{CN}}$  is the probability of forming the compound nucleus (CN) and  $W_{\text{sur}}$  is the fission survival probability of the CN. Here,  $P_{\text{CN}}$  (or  $P_{\text{fus}}$ , the fusability) is the least known quantity.

Following have been reported about  $P_{\text{CN}}$ :

- $P_{\text{CN}}$  does not depend much on the charge of the CN, but mainly on the combination of fusing nuclei [1].
- $P_{\text{CN}}$  demonstrates a sharp energy dependence. It increases with the excitation energy  $E^*$  [1].
- For the more asymmetric hot combinations, the  $P_{\text{CN}}$  is usually much higher than for the cold combinations leading to the same elements [2].
- Systems with similar effective fissility, can have  $P_{\text{CN}}$  values differing by as much as an order of magnitude [3].

$P_{\text{CN}}$  can be extracted by analysing the shape of the fission-fragment (FF) angular distribution or the FF mass distribution. However, these methods are not robust and are yet to have predictive power. There are a few theoretical estimations and systematics of  $P_{\text{CN}}$

too. It has been reported that  $P_{\text{CN}}$  is approximately constant at energies above the barrier [4, 5]. A global systematics of  $P_{\text{CN}}$  will also aid super-heavy element (SHE) synthesis.

Sahm et al. [6] defined the angular momentum-averaged fusability as

$$\langle P_{\text{CN}} \rangle = \frac{\sigma_{\text{ER}}}{\sigma_{\text{ER}}^0} \quad (2)$$

where,  $\sigma_{\text{ER}}^0$  is the maximum limit of the ER cross section ( $\sigma_{\text{ER}}$ ) at a given  $E^*$  which corresponds to  $P_{\text{CN}} = 1$ . Therefore, it can be calculated by a statistical model without considering fusion hindrance. Sagaidak [7] also followed a similar approach for estimating  $P_{\text{CN}}$  for a number of heavy and super-heavy nuclei.

In the present work, we have calculated  $P_{\text{CN}}$  using Eq. (2) to study its behavior as a function of mass-asymmetry ( $\eta$ ). We have performed calculation for a number of systems forming three CNs having well separated mass numbers, namely  $^{200}\text{Pb}$ ,  $^{210}\text{Rn}$  and  $^{220}\text{Th}$ .

First, we calculated  $\sigma_{\text{cap}}$  for these systems using the coupled-channels code CCFULL [8]. Proper rotational and vibrational couplings to the targets and projectiles were considered in the calculation. Akjuz-Winther parametrization of Woods-Saxon potential was used in the input. The spin-distribution generated by the CCFULL was fed to the input of a statistical model (SM) code [9]. The code considers evaporation of neutrons, protons,  $\alpha$ -particles and statistical giant dipole resonance (GDR)  $\gamma$ -rays as the decay channels of an excited CN in addition to fission. Shell-effect in level-density and shell-correction to the barrier height were included in the calculation. Fission barrier was not scaled and dissipation strength was not taken into account.

While comparing with experimental  $\sigma_{\text{ER}}$ , calculated  $\sigma_{\text{ER}}$  is overestimated in all cases except for the system  $^{16}\text{O}+^{184}\text{W}$ . Shidling et al.

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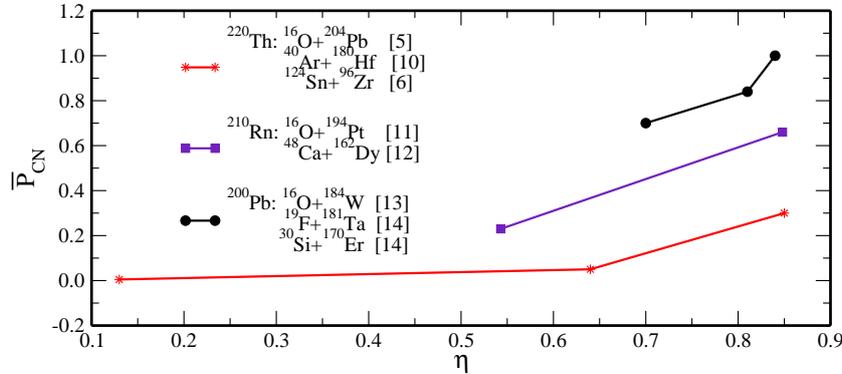


FIG. 1:  $P_{CN}$  as a function of  $\eta$ . Experimental  $\sigma_{ER}$  have been taken from Refs. [5, 6, 10–14].

[13] reported presence of dissipative forces in this system. Hence, we considered  $P_{CN} = 1$  for this reaction. For all other reactions,  $P_{CN}$  was estimated by scaling down the calculated  $\sigma_{ER}$  onto the experimental  $\sigma_{ER}$ . It is clear from Fig. 1 that  $P_{CN}$  decreases with decreasing  $\eta$ . In other words, fusability decreases for symmetric combinations. Also,  $P_{CN}$  goes down with increase in mass number of the CN.

Hinde et al. [5] obtained  $P_{CN}$  for a number of reactions, all leading to the CN  $^{224}\text{Th}$ , by comparing measured xn  $\sigma_{ER}$  without using model calculations. Sahm et al. [6], on the other hand, estimated  $P_{CN}$  by comparing experimental  $\sigma_{ER}$  with calculated  $\sigma_{ER}$ . One may argue that such calculations are susceptible to errors due to improper choice of parameters in the statistical model. Liang et al. [15] reported estimation of fusion probability for neutron-rich Sn-induced reactions, in which statistical model calculation was performed by the code PACE2 [16]. Statistical model parameters were varied within reasonable limits to study the uncertainties in the calculated  $\sigma_{ER}$ . It was demonstrated that the dependence of  $P_{CN}$  on the effective fissility parameter was not sensitive to the choice of parameters in the calculation. However, the magnitude of  $P_{CN}$  differed with the choice of statistical model parameters by a few percents.

We plan to carry out such calculations for more systems in the heavy mass region to gain a more complete understanding of  $P_{CN}$  with  $\eta$ .

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