

## Fusion hindrance for asymmetric systems at extreme sub barrier energies

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### Introduction

Recent measurements with medium-heavy nuclei accentuated phenomenon of fusion hindrance, observed as a steep change of slope in fusion excitation function and its logarithmic derivative ( $L(E)$ ) with respect to the coupled channels (CC) calculation at deep sub-barrier energies [1]. At present there are two successful models to explain the deep sub-barrier fusion data - model suggested by Mişicu and Esbensen [2] based on sudden approximation using M3Y potential with repulsive core and a dynamical two-step model proposed by Ichikawa *et al.* [3] based on an adiabatic picture. Currently experimental studies at these low energies have been restricted mainly to the measurement of fusion cross sections of symmetric systems with the exception of  $^{16}\text{O} + ^{204,208}\text{Pb}$  [4] and  $^6\text{Li} + ^{198}\text{Pt}$  [5]. Unlike the sharp change in slope of  $L(E)$  as observed in symmetric medium-heavy systems, a saturation in the slope of  $L(E)$  was observed for asymmetric  $^{16}\text{O} + ^{208}\text{Pb}$  system [4]. In case of very asymmetric system involving light weakly bound projectile  $^6\text{Li} + ^{198}\text{Pt}$  [5] absence of fusion hindrance was reported. In the present work we investigate whether absence of fusion hindrance in  $^6\text{Li} + ^{198}\text{Pt}$  system arises from the effect of weakly bound cluster structure or it is a property of the very asymmetric systems. For this purpose we selected one

weakly bound ( $^7\text{Li}$ ) and a tightly bound ( $^{12}\text{C}$ ) projectile on the same target. The data is analysed with standard coupled channels and the adiabatic model of fusion to understand the underlying mechanism at deep sub barrier energies.

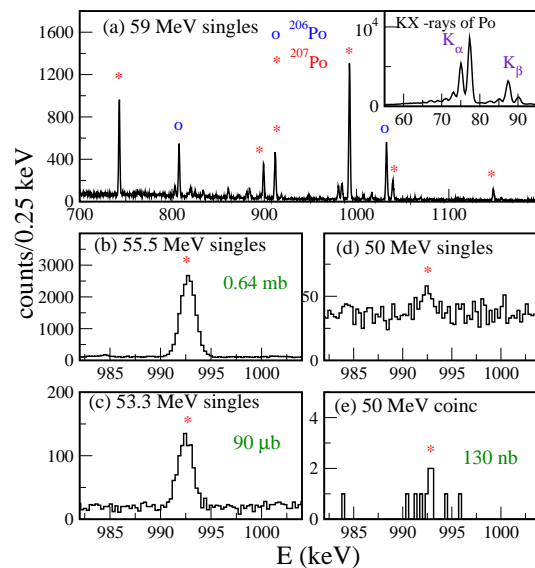


FIG. 1: Activation  $\gamma$ -ray spectra for the  $^{12}\text{C} + ^{198}\text{Pt}$  system (a) inclusive spectrum at  $E_{lab}$  of 59 MeV. Inset shows X-ray region of the spectrum. The dominant  $\gamma$ -rays arising from the evaporation residues are labeled (b)-(d) inclusive spectra showing photo-peak at 992.3 keV, corresponding to the residue  $^{207}\text{Po}$  (e) same as (d) but in coincidence with the  $K_{\alpha}$  X rays shown in the inset of (a).

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## Experimental Details

The experiment was performed at Pelletron Linac Facility-Mumbai, using beams of  ${}^7\text{Li}$  (20 - 35 MeV) and  ${}^{12}\text{C}$  (50 to 64 MeV) on a  ${}^{198}\text{Pt}$  target with beam current in the range of 10 to 35 pnA. The targets were self supporting rolled foils of enriched  ${}^{198}\text{Pt}$  ( $\sim 1.3 \text{ mg/cm}^2$  thick) followed by an Al catcher foil. Two efficiency calibrated HPGe detectors - one with Al window for detection of  $\gamma$ -rays and another with a Be window for detection of KX-rays were used for performing KX- $\gamma$ -ray coincidence. The measurements were performed in a low background setup with a graded shielding. The evaporation residues were uniquely identified by means of their characteristic  $\gamma$ -ray energies and half-lives. The residues populated are  ${}^{205-207}\text{Po}$  in case of  ${}^{12}\text{C}+{}^{198}\text{Pt}$  and  ${}^{200-202}\text{Tl}$  in case of  ${}^7\text{Li}+{}^{198}\text{Pt}$ . Typical inclusive  $\gamma$ -ray spectra resulting from the residues of  ${}^{12}\text{C}+{}^{198}\text{Pt}$  are plotted in Fig. 1(a)-(d) at different beam energies. The  $\gamma$ -ray yields at lowest energies were extracted by gating on their KX-ray transitions. Due to the increased sensitivity of the KX- $\gamma$ -ray coincidence method, cross-section down to a few nano-barns could be measured (Fig. 1(e)).

## Analysis and Summary

The data were analysed using the standard coupled-channels (CC) calculations and the adiabatic model that simulates a smooth transition between the two-body and the adiabatic one-body states by damping gradually the off-diagonal part of the coupling potential [6]. The standard CC calculations for both the systems were performed including the quadrupole excitation in  ${}^{198}\text{Pt}$  in the vibrational mode. Projectiles  ${}^7\text{Li}$  and  ${}^{12}\text{C}$  were coupled in the rotational mode. In case of  ${}^7\text{Li}+{}^{198}\text{Pt}$ , fusion hindrance was not observed as the CC calculations nicely reproduce the data for energies around and well below the barrier in the measured energy range. A change in slope in fusion excitation function and  $L(E)$  as compared to CC calculations was clearly observed for  ${}^{12}\text{C}+{}^{198}\text{Pt}$ , at lowest energies indicating onset of fusion hindrance. In order to explain the fusion data at energies

deep below the barrier in case of  ${}^{12}\text{C} + {}^{198}\text{Pt}$  system, calculations were performed using the adiabatic model of Ref. [6]. On inclusion of damping in the adiabatic framework an excellent agreement with the fusion and  $L(E)$  data was observed as shown in Fig. 2. These results are relevant to fusion hindrance at deep sub-barrier energies with respect to observations from  ${}^{6,7}\text{Li} + {}^{198}\text{Pt}$ ,  ${}^{16}\text{O} + {}^{208}\text{Pb}$  and systems with different mass asymmetry.

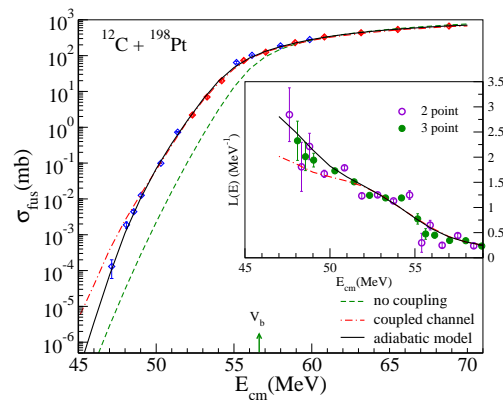


FIG. 2: Fusion excitation function and  $L(E)$  for  ${}^{12}\text{C} + {}^{198}\text{Pt}$ . CC calculations with and without inclusion of coupling along with adiabatic model calculations are shown.

## Acknowledgments

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## References

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