

Microscopic study of oblate to prolate shape transition at higher spins in neutron-rich $^{100-104}\text{Zr}$ isotopes

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Neutron-rich Zr nuclei in the A~100 mass region have attracted both theoretical and experimental attentions because of some interesting phenomena observed in these nuclei. The ground states of Zr isotopes undergo a shape transition from spherical to well deformed prolate (or oblate) deformations at neutron number (N)~60 [1-4]. However, this transition becomes steadily weaker as proton number (Z) is increased. It has been shown by the theoretical relativistic [5,6] and non-relativistic [7,8] studies of the nuclear structural evolution in this mass region that there is a rapid change in the equilibrium shape of the nucleus as a function of number of nucleons, and shape coexistence is present with competing spherical, axially symmetric prolate and oblate shapes.

The ground-state bands in case of $^{100-104}\text{Zr}$ isotopes have been extended up to high spins of 20^+ , 20^+ and 14^+ , respectively, by Hua et al.[9]. Besides yrast bands, the experimental data is available on electromagnetic quantities such as B(E2) transition probabilities and g-factors for neutron-rich $^{100-104}\text{Zr}$ [10-13]. Smith et al. [11] have calculated B(E2; $2_1^+ \rightarrow 0_1^+$) values in $^{100,102}\text{Zr}$ from the lifetime measurement of 2_1^+ state. Hwang et al. [10] have also measured the lifetime and hence B(E2) value of 2_1^+ state in ^{104}Zr via γ -ray spectroscopy. The latest B(E2) values extracted from the lifetime of 2_1^+ state have been reported in ^{104}Zr by Browne et al. [13] by using β - γ spectroscopy. The experimental information on B(E2) values in neutron-rich $^{100-104}\text{Zr}$ isotopes have been very rare until Smith et al. [12] presented new lifetime measurements for excited rotational states in these nuclei at spins I~10. These authors have measured the lifetimes for $8 \rightarrow 6$, $10 \rightarrow 8$, $12 \rightarrow 10$ transitions in the yrast bands of $^{100-104}\text{Zr}$ isotopes. Moreover, transitional quadrupole moments at nuclear spins I~10 have been extracted and compared with

various theoretical approaches [8,14]. It has been interpreted from the results that there is prolate oblate shape coexistence in Zr isotopes for spins less than 10^+ , which in turn indicates that the shapes of Zr isotopes are unchanged up to 8^+ . Another interpretation is that there is a gradual oblate to prolate shape transition at 10^+ in Zr isotopes. In the present work, the Projected Shell Model (PSM) [15] has been employed as a theoretical tool in order to examine the presence of prolate oblate shape coexistence in low-lying states of Zr isotopes as well as understand the mechanism underlying oblate to prolate shape transition at 10^+ as suggested by Smith et al. [12]. The PSM calculations have been performed for both prolate as well as oblate deformations for $^{100-104}\text{Zr}$ isotopes. The valence space of the model consists of three harmonic oscillator shells N=2,3,4 for protons and 3,4,5 for neutrons.

In Fig.1(a,b,c) the theoretical yrast spectra obtained in PSM framework are compared with the experimental data for $^{100-104}\text{Zr}$ isotopes, respectively. The results presented in Fig. 1(a) show that in case of ^{100}Zr , low-lying yrast states up to 6^+ are seen to be in agreement with the corresponding states arising from both prolate as well as oblate deformations with $\epsilon_2 = 0.29$ and -0.42 , respectively. Thus, the PSM calculations predict coexistence of prolate and oblate shapes for low-lying states in ^{100}Zr . However, the yrast states from 8^+ to 14^+ show a better agreement with the corresponding states arising from the prolate deformation. Above spin 14^+ , the prolate yrast spectrum is lower and the oblate yrast spectrum is higher than the observed yrast spectrum. In case of ^{102}Zr (Fig. 1(b)), the observed yrast spectrum up to 10^+ is reproduced by the corresponding states arising from both prolate as well as oblate deformations with $\epsilon_2 = 0.33$ and -0.47 , respectively, thereby suggesting shape coexistence in this nucleus up to 10^+ .

However, for the yrast states above 10^+ , the energy states arising from prolate deformation are in better agreement than those arising from oblate deformation. In case of ^{104}Zr (Fig. 1(c)), the observed yrast spectrum up to 10^+ is in agreement with the corresponding states arising from both prolate as well as oblate deformations with $\epsilon_2=0.32$ and -0.40 , respectively, thereby suggesting shape coexistence in this nucleus up to spin 10^+ . However, the observed yrast spectrum at spin 12^+ is well reproduced by the yrast spectrum arising from prolate deformation.

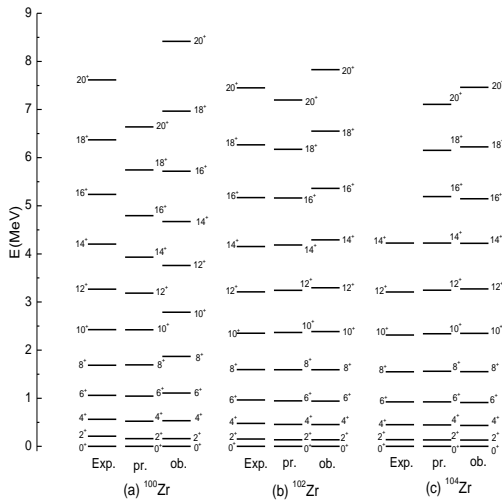


Fig. 1 Comparison of calculated yrast spectra with the experimental data for (a) ^{100}Zr (b) ^{102}Zr (c) ^{104}Zr , respectively

Acknowledgements

One of the authors (Prof. S.K. Khosa) is grateful to the Council of Scientific and Industrial Research, New Delhi for financial support under the Emeritus Scientist grant No. 21(0959)/13/EMR-II.

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