

## Nucleosynthesis in decompressed Neutron stars crust matter

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### Introduction

It's being a challenging problem to understand the formation of heavy & super heavy nuclei in the universe [1]. The rapid neutron capture process (r-process) is believed be responsible for the synthesis of heavier elements in supernovae explosions & neutron stars (NS) crusts under extreme astrophysical conditions [2]. Some NSs have very high magnetic fields~ $10^{17}G$ . These are called as magnetars. The ground state properties of inner crusts of NSs are studied in the presence of strong magnetic fields [3]. We studied the r-process in the decompressed NSs crust matter in the presence of strong magnetic field using the calculations of [3] as input.

### Nuclear Statistical Equilibrium

During fusion reactions inside the stars, once Si burning stage is reached, the temperature rises up to a limit, when various nuclei reach one large equilibrium group stretching from proton (p), neutron (n),  $\alpha$ -particles to the iron peak nuclei. The system attains nuclear statistical equilibrium (NSE) with equal forward & backward reaction rates. We consider the hot and dense decompressed crust matter consists of free p, n & seed nuclei such as  $^{56}Fe$ .

At NSE, mass & charge conservation imply [4],

$$\sum X_i = 1 \text{ ----- (1)}$$

$$\sum \frac{Z_i}{A_i} X_i = Y_e \text{ ----- (2)}$$

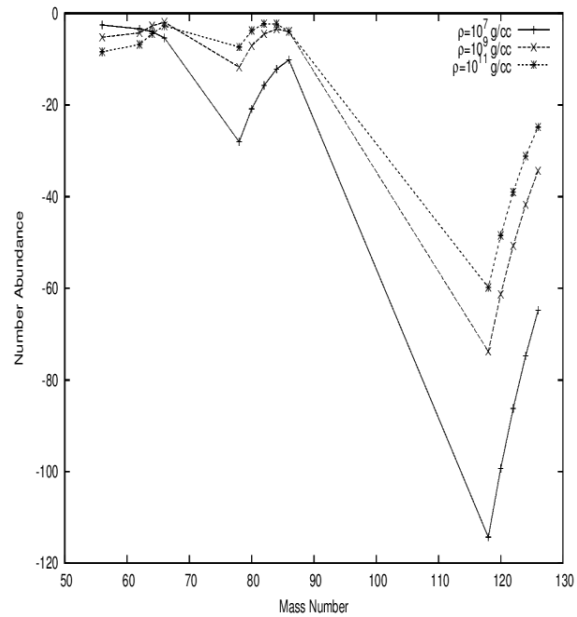
where,  $Y_e$  is the proton fraction,  $X_i$  is the mass abundance for i-th nuclei with atomic number  $Z_i$  and mass number  $A_i$  & is given by,

$$X_i = \frac{1}{2} G_i(T) \left( \frac{1}{2} \rho N_A \lambda^3 \right)^{A_i-1} A_i^{5/2} X_n^{A_i-Z_i} X_p^{Z_i} e^{Q_i/KT} \text{ ----- (3)}$$

Here,  $G_i(T)$  is the partition function,  $N_A$  is the Avogadro's number,  $Q_i$  is the binding energy,  $\rho$  is the mas density, T is the temperature, K is the Boltzmann constant,  $X_p$  is the proton fraction,

$X_n$  is the neutron fraction &  $\lambda$  is the thermal De' Broglie wavelength given by,  $\lambda = \frac{h}{\sqrt{2\pi m_H K T}}$  where h is the Planck's constant &  $m_H$  is the mass of a hydrogen atom.

We solve eqns.(1) & (2) using Newton-Raphson method for  $X_p$  &  $X_n$  & use these values to obtain the mass abundance of different nuclei. This leads to number abundances of different nuclei,  $Y_i = \frac{X_i}{m_i}$  [5].



**Fig.1** Abundance variation with mass number for  $T=9 \times 10^9 K$  &  $Y_e=0.4$  with varying  $\rho$

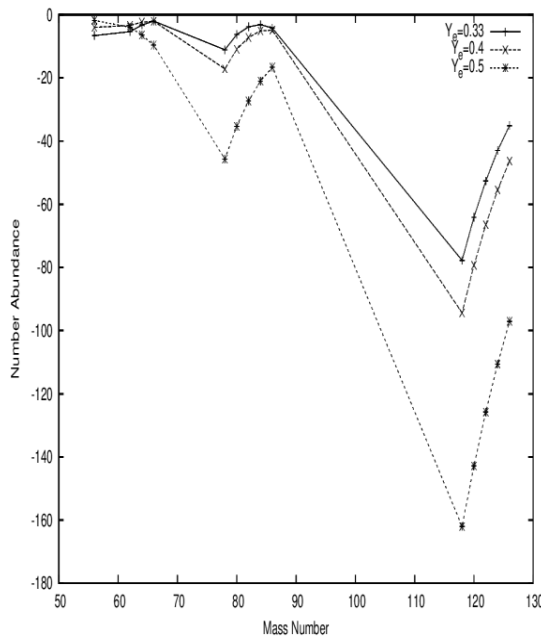
### Result

We have shown the abundance variation with mass number for NS crust nuclei. Fig.1. shows the abundance variation with mass number for constant T &  $Y_e$  with varying density ( $\rho$ ). It can be seen that abundance of nuclei

increases with increasing  $\rho$ .  $^{86}\text{Kr}$  is the most abundant nuclei &  $^{118}\text{Kr}$  is the least abundant nuclei.

Fig.2. shows the variation of abundance for constant  $\rho$  & T with varying  $Y_e$ . Abundance of nuclei increases with decreasing  $Y_e$ . More nuclei present in low  $Y_e$  region.

We have also studied the abundance variation with mass number for constant  $Y_e$  &  $\rho$  with varying T and found that the nuclei are more abundant for low temperature values.



**Fig. 2** Abundance variation with mass number for  $T=8 \times 10^9 \text{K}$  &  $\rho=10^8 \text{g/cc}$  with varying  $Y_e$

### Summary

We have investigated the abundance of neutron star crust in the presence of high magnetic field in the outer crust of decompressed NSs crust matter.

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### References

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