

## Description of Long-Lived Isomers in $^{254}_{101}\text{Md}_{153}$

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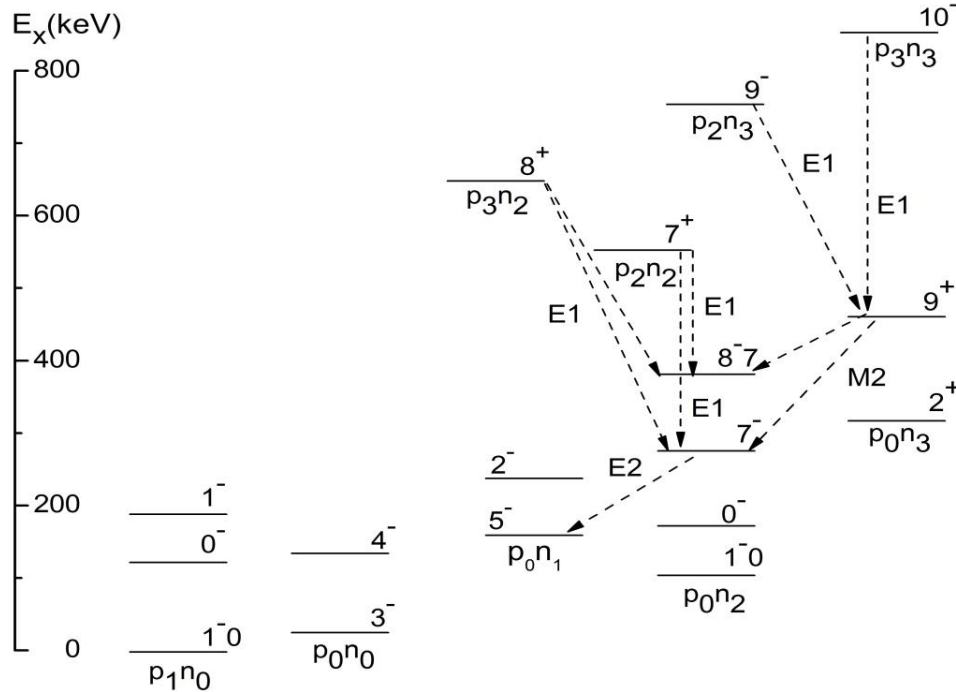
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Spectroscopic data in respect of Trans-Fermium nuclei beyond the N=152 deformed shell closure are very rare, and rarer still is a credible interpretation of such data [1]. Presently available experimental information [1,2] on the lightest such nucleus  $^{254}_{101}\text{Md}_{153}$  provides a typical instance illustrating this feature. This nucleus, first observed in 1970 [2], ‘was found to decay by EC with a  $t_{1/2}=10(3)$  m’; these experiments also ‘indicated another  $t_{1/2}=28(8)$  m isomer’. Even 45 years later no attempt to characterize these results has been reported anywhere [1]. The latest Nuclear Data Sheets [2] mentions 8 different assignments for  $^{254}\text{Md}$  ground state (gs) and low-lying states. Presently we seek to investigate the low-energy level structures in  $^{254}\text{Md}$  and thereby deduce its gs spin-parity ( $J^\pi$ ) and two-quasiparticle (2qp) configuration and also characterize its earlier indicated long-lived isomers. Our analysis employs the Two Quasiparticle Rotor Model (TQRm) which has been extensively and effectively used to describe the level schemes of various doubly odd deformed nuclei over the past three decades [3]. In particular, level structures of the isotonic  $^{252}\text{Es}$  [4] and isotopic  $^{250}\text{Md}$  [5] neighbours have been investigated earlier by us using this formalism.

In our 3-step TQRm, the relevant 1qp configuration space is mapped using the experimental n-orbital energies from a recent study [6] of N=153  $^{253}\text{Fm}$  levels and p-orbital energies from ( $Z\pm 2$ ) neighbours [1]. These data are respectively entered in the first column and top row of Table 1. Physically admissible  $K=\left|\Omega_p \pm \Omega_n\right|$  2qp bands arising from the coupling of these 1qp Nilsson orbitals form the subject of Table 1. Entries in each box therein are the 2qp band quantum numbers  $K_T$  (spins-parallel triplet) and  $K_S$  (spins-antiparallel singlet) according to the GM rule; the numbers within parentheses are summed ( $E_p + E_n$ ) energies in keV, which provide a zeroth order estimate of  $E_x(K_T)$ , with  $E_x(K_S)$  estimated  $\sim 100$  keV above it. Finally, the 2qp bandhead energies are evaluated in the TQRm formulation [3,4] using the model parameters obtained by fitting the experimental spectra of isotonic  $^{252}\text{Es}$  and  $^{250}\text{Bk}$  neighbours [3,4];  $E_{GM} \sim 100$  keV is taken as default value alternatively. Partial  $^{254}\text{Md}$  level scheme thus obtained is shown in Fig.1 in the context of characterizing the suggested two isomers. A critical examination of these data leads us to the following conclusions.

**Table 1:** Listing of physically admissible 2qp GM doublet bands in  $^{254}_{101}\text{Md}_{153}$  from coupling of the observed 1qp p-orbitals from ( $Z\pm 2$ ) neighbours (top row) and n-orbitals from  $^{253}\text{Fm}$  [1,4,6] (first column). Numbers within parenthesis in each box are the summed [ $E(p_i) + E(n_j)$ ] energies in keV.

$n_j \backslash p_i$	$p_0: 0$ 7/2-[514↓]		$p_1: 40$ 1/2-[521↓]		$p_2: \sim 400$ 7/2+[633↑]		$p_3: \sim 500$ 9/2+[624↑]	
	$K_T$	$K_S$	$K_T$	$K_S$	$K_T$	$K_S$	$K_T$	$K_S$
$n_0: 0$ 1/2+[620↑]	3 <sup>-</sup>	4 <sup>-</sup>	0 <sup>-</sup>	1 <sup>-</sup>	4 <sup>+</sup>	3 <sup>+</sup>	5 <sup>+</sup>	4 <sup>+</sup>
	(0)		(40)		(~ 400)		(~ 500)	
$n_1: 124$ 3/2+[622↓]	2 <sup>-</sup>	5 <sup>-</sup>	2 <sup>-</sup>	1 <sup>-</sup>	2 <sup>+</sup>	5 <sup>+</sup>	3 <sup>+</sup>	6 <sup>+</sup>
	(124)		(164)		(524)		(624)	
$n_2: 140$ 7/2+[613↑]	0 <sup>-</sup>	7 <sup>-</sup>	3 <sup>-</sup>	4 <sup>-</sup>	7 <sup>+</sup>	0 <sup>+</sup>	8 <sup>+</sup>	1 <sup>+</sup>
	(140)		(180)		(540)		(640)	
$n_3: 330$ 11/2-[725↑]	2 <sup>+</sup>	9 <sup>+</sup>	5 <sup>+</sup>	6 <sup>+</sup>	9 <sup>-</sup>	2 <sup>-</sup>	10 <sup>-</sup>	1 <sup>-</sup>
	(330)		(370)		(770)		(870)	



**Fig. 1:** Partial level scheme of  $^{254}\text{Md}$  including low-lying (< 250 keV) low-spin and all the high-spin ( $\geq 7$ ) bands. Arrows therein denote the lowest multipole decay path for each high-spin ( $K \geq 7$ ) state.

(a) The lowest-lying levels in  $^{254}\text{Md}$  spectrum correspond to the very closely spaced low-spin  $1^-0\{p_1:1/2^- \otimes n_0:1/2^+\}$  and  $3^- \{p_0:7/2^- \otimes n_0:1/2^+\}$  configurations. Our present set of input parameters places the  $1^-0$  level lower than the  $3^-$  level, and hence designates it as the  $^{254}\text{Md}$  gs.

(b) Existence of two isomers with comparable  $t_{1/2}$  values, each with  $\% \epsilon \leq 100$  decay mode, normally requires  $\Delta K \geq 5$ . Having identified a low spin gs, we examine location of all the physically admissible high-spin ( $K \geq 7$ ) 2qp structures (Fig.1). It is evident therein that each  $K \geq 7$  state admits of a low multipole ( $\Delta I \leq 2$ ) em decay path. Consequently we categorically rule out a high-spin structure as a long-lived  $^{254}\text{Md}$  isomer.

(c) Examining possible alternatives, we find that the pair of low-lying  $1^-0$  and  $3^-$  levels offer such a possibility. The two levels have same  $n_0$  constituent, while the proton orbitals, namely  $p_0 7/2^-$  &  $p_1 1/2^-$ , differ by  $\Delta K = \Delta \Omega = 3$ . The current data tables [1] reveal that in  $_{101}\text{Md}$  ( $A=245, 247, 249$ ) and  $_{103}\text{Lr}$  ( $A=253$  &  $255$ ) isotopes these p orbitals occur as close-lying isomers with comparable  $t_{1/2}$  in each case. Occurrence of  $\Delta K=3$

isomers with  $t_{1/2} \sim m$  in  $^{166}\text{Lu}$  has also been pointed out earlier by us [7].

Considering all these factors, we arrive at the following assignments for  $^{254}\text{Md}$  isomer pair:  
 10 m gs  $1^-0\{p_1:1/2^-[521] \otimes n_0:1/2^+[620]\}$   
 28 m 0+x  $3^-3\{p_0:7/2^-[514] \otimes n_0:1/2^+[620]\}$

Investigations of these and other distinctive features of  $^{254}\text{Md}$  level scheme are being pursued.

### References

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