

Neutrinoless double beta decay and Physics beyond the Standard Model

Yash Kaur Singh¹, Pooja Lohani², V. K. Nautiyal¹, R. Gautam¹, R. Chandra¹,* K. Chaturvedi³, P. K. Rath² and P. K. Raina⁴

¹ Department of Applied Physics, Babasaheb Bhimrao Ambedkar University, Lucknow - 226025, INDIA

² Department of Physics, University of Lucknow, Lucknow, 226007, INDIA

³ Department of Physics, Bundelkhand University, Jhansi - 284128, INDIA

⁴ Department of Physics, IIT Ropar, Nangal Road, Rupnagar, Punjab - 140001, INDIA

* email: ramesh.luphy@gmail.com

Introduction

The nuclear $\beta\beta$ decay, in which the charge Z of an even Z -even N nucleus is changed by two units while the mass number A remains the same, is a rare and spontaneous process of weak interaction in nature. The $\beta\beta$ decay can be categorized in mainly two modes, namely two neutrino double beta $(\beta\beta)_{2\nu}$ decay and neutrinoless double beta $(\beta\beta)_{0\nu}$ decay. The $(\beta\beta)_{2\nu}$ decay establishes the validity of different nuclear models employed for nuclear structure calculations by calculating the nuclear transition matrix elements (NTMEs) $M_{2\nu}$. The $(\beta\beta)_{0\nu}$ decay violates the lepton number conservation by two units and can be studied in any theory in which lepton number conservation is not exact. Hence $(\beta\beta)_{0\nu}$ decay is an excellent process to probe the new physics beyond the standard model (SM) of electroweak unification. Apart from the well studied left-right symmetric model, the $(\beta\beta)_{0\nu}$ decay can also be studied in Majoron models, R-parity violating as well as conserving supersymmetric (SUSY) models. Further, the $(\beta\beta)_{0\nu}$ decay can verify issues like compositeness, leptoquarks, sterile neutrinos and violation of weak equivalence principle. The detailed progress of experimental as well as theoretical studies on $\beta\beta$ decay in general and $(\beta\beta)_{0\nu}$ decay in particular can be found in references [1,2] and references there in.

The PHFB model in conjunction with pairing plus multipole type of two-body effective interaction has been successfully applied to study the $(\beta\beta)_{0\nu}$ decay [3-7]. In present work the same PHFB model has been applied to study the $(\beta\beta)_{0\nu}$ decay in various theories beyond the SM.

Theoretical framework

The details about the model space, single particle energies, PQQ type of effective two-body interaction and the procedure to fix its parameters have been given in Refs. [3,4]. The Hamiltonian of the effective two-body interaction used in the present work is given as

$$H = H_{s.p.} + V(P) + V(QQ) + V(HH) \quad (1)$$

where $H_{s.p.}$, $V(P)$, $V(QQ)$ and $V(HH)$ denote the single particle Hamiltonian, pairing, quadrupole-quadrupole and hexadecapole-hexadecapole parts of the effective two-body interaction. We use four different parametrizations of the interaction Hamiltonian, namely $PQQ1$, $PQQ2$, $PQQHH1$ and $PQQHH2$ [5]. Further, we use the Jastrow type of short range correlations with Miller-Spencer, Argonne V18 and CD-Bonn NN potentials [5,8]. The detailed theoretical formalism to study the $(\beta\beta)_{0\nu}$ decay in Majoron and SUSY models are given in refs. [9,10,11]. Further, the theory of compositeness and leptoquark in connection with $(\beta\beta)_{0\nu}$ decay is given in refs. [12,13].

Results and discussions

The NTMEs involved in $(\beta\beta)_{0\nu}$ decay in various theories stated above are calculated within PHFB model using pairing plus multipole type of two-body interaction. The NTMEs have been calculated by considering the finite size of nucleon (F) and Jastrow type of short range correlations (SRC) with Miller-Spencer, Argonne V18 and CD-Bonn NN potentials for $(\beta\beta)_{0\nu}$ decay of $^{94,96}\text{Zr}$, ^{100}Mo , $^{128,130}\text{Te}$ and ^{150}Nd isotopes for the $0^+ \rightarrow 0^+$ transition. At present, some results are presented for the case of ^{100}Mo

with *PQQI* parametrization. The detailed results will be presented in the symposium.

Table 1: NTMEs for the Majoron accompanied $(\beta\beta)_{0\nu}$ decay of ^{100}Mo for the *PQQI* parametrization.

| NTME | F | F+S | | |
|---------------------------------------|-------|-------|-------|-------|
| | | SRC1 | SRC2 | SRC3 |
| $M_{Fm_\nu}^{(\chi)}$ | 2.15 | 1.89 | 2.15 | 2.22 |
| $M_{GTm_\nu}^{(\chi)}$ | -5.52 | -4.72 | -5.43 | -5.66 |
| $M_{Tm_\nu}^{(\chi)}$ | 0.05 | 0.05 | 0.05 | 0.05 |
| $M_{CR}^{(\chi)}$ | -0.26 | -0.23 | -0.26 | -0.27 |
| $M_{F\omega^2}^{(\chi)} \times 10^3$ | 1.18 | 1.14 | 1.19 | 1.20 |
| $M_{GT\omega^2}^{(\chi)} \times 10^3$ | -5.79 | -5.60 | -5.86 | -5.91 |

Table 2: NTMEs of $(\beta\beta)_{0\nu}$ decay in SUSY models via exchange of gluinos for ^{100}Mo with in PHFB model using *PQQI* interaction.

| NTME | F | F+S | | |
|-----------------|--------|--------|--------|--------|
| | | SRC1 | SRC2 | SRC3 |
| M_F^N | 0.716 | 0.037 | 0.056 | 0.067 |
| M_F^P | 0.007 | 0.003 | -0.002 | 0.002 |
| M_{GT}^N | -0.210 | -0.108 | -0.164 | -0.196 |
| M_{GT}^P | -0.023 | 0.009 | 0.004 | -0.007 |
| $M_{GT}^{1\pi}$ | 5.422 | 2.169 | 3.80 | 4.831 |
| $M_{GT}^{2\pi}$ | 2.613 | 1.856 | 2.45 | 2.661 |
| M_T^P | 0.001 | 0.001 | 0.001 | 0.001 |
| $M_T^{1\pi}$ | 0.285 | 0.285 | 0.302 | 0.302 |
| $M_T^{2\pi}$ | 0.130 | 0.132 | 0.134 | 0.134 |

Table 3: NTMEs of $(\beta\beta)_{0\nu}$ decay in SUSY models via exchange of squark for ^{100}Mo with in PHFB model using *PQQI* interaction.

| NTME | F | F+S | | |
|--------------|--------|--------|--------|--------|
| | | SRC1 | SRC2 | SRC3 |
| M_F | -258.7 | -131.7 | -201.4 | -241.1 |
| M_{GT-MT} | 412.4 | 207.4 | 318.8 | 382.1 |
| M_{GT-AP} | -25.77 | -20.00 | -24.92 | -26.57 |
| $M_{GT-\pi}$ | 506.4 | 441.8 | 511.2 | 528.9 |
| M_{T-MT} | -10.33 | -10.36 | -10.85 | -10.83 |
| M_{T-AP} | -1.13 | -1.14 | -1.15 | -1.15 |
| $M_{T-\pi}$ | 17.90 | 18.06 | 18.09 | 18.07 |

In Table 1, 2 and 3, SRC1, SRC2 and SRC3 denote the Jastrow type of short range correlations (SRC) with Miller-Spencer, Argonne V18 and CD-Bonn NN potentials, respectively.

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