# Evaporation residue cross-section measurements for compound nuclei around $Z_{CN}=82$ region

Priya Sharma<sup>1</sup>,\* B.R. Behera<sup>1</sup>, N. Madhavan<sup>2</sup>, I. Mazumdar<sup>3</sup>, Ruchi Mahajan<sup>1</sup>,

Meenu Thakur<sup>1</sup>, Gurpreet Kaur<sup>1</sup>, Kushal Kapoor<sup>1</sup>, Kavita Rani<sup>1</sup>, J. Gehlot<sup>2</sup>, S. Nath<sup>2</sup>, M. Dhibar<sup>4</sup>, S.M. Patel<sup>3</sup>, M.M. Hosamani<sup>5</sup>, Khushboo<sup>6</sup>, R. Dubey<sup>2</sup>, A. Shamlath<sup>7</sup>, N. Kumar<sup>6</sup>, G. Mohanto<sup>8</sup>, and Santanu Pal<sup>9</sup>

Department of Physics, Panjab University, Chandigarh - 160014, INDIA

<sup>2</sup>Inter University Accelerator Centre, Aruna Asaf Ali Marg, New Delhi - 110067, INDIA

<sup>3</sup> Tata Institute of Fundamental Research, Mumbai - 400005, INDIA

<sup>4</sup>Department of Physics, Indian Institute of Technology, Roorkee - 247667, INDIA

<sup>5</sup> Department of Physics, Karnatak University, Dharwad - 580003, INDIA

<sup>6</sup>Department of Physics and Astrophysics, University of Delhi - 110007, INDIA

<sup>7</sup>Department of Physics, Central University of Kerala, Kasaragod - 671314, INDIA

<sup>8</sup>Nuclear Physics Division, Bhabha Atomic Research Centre, Mumbai - 400085, INDIA and

<sup>9</sup>CS-6/1 Golf Green, Kolkata-700095, INDIA (Formerly with VECC, Kolkata)

## Introduction

For understanding the reaction mechanism involving heavy compound nuclei (CN), we have performed measurements of the evaporation residue (ER) cross-section for the  $(i)^{48}\text{Ti}+^{150}\text{Nd}\rightarrow^{198}\text{Pb}$ ,  $(ii)^{48}\text{Ti}+^{142}\text{Ti}+^{160}\text{Nd}\rightarrow^{198}\text{Pb}$  $^{142}$ Nd $\rightarrow$   $^{190}$ Pb and (iii)  $^{48}$ Ti+ $^{144}$ Sm $\rightarrow$   $^{192}$ Po systems leading to CN around  $Z_{CN}=82$ region. It is conjectured that proton shell closure  $Z_{CN}=82$  may lead to enhanced ER production [1] and may help in the synthesis of heavy nuclei. Here, we investigate the effect of proton shell closure  $Z_{CN}=82$  and deformation through the ER cross-section measurements.

The experiment was carried out in the gasfilled mode of HYbrid Recoil mass Analyzer (HYRA) [2] using 15 UD Pelletron+LINAC accelerator facility at IUAC, New Delhi. More details of the experimental set up is available in Ref. [3].

#### Experimental analysis

For the determination of ER cross-section, extraction of transmission efficiency plays a key role. For the efficiency measurement, <sup>48</sup>Ti+<sup>122</sup>Sn was used as a calibration reaction. Using the already existing experimental ER cross-section [4] and the Monte Carlo simulation code TERS [5], experimentally extracted average transmission efficiency value for  ${}^{48}\text{Ti} + {}^{122}\text{Sn}$  reaction is  $28.71 \pm 5.6\%$ . The efficiency value of the calibration reaction was further normalized for other systems to obtain the ER cross-section for  ${}^{48}$ Ti+ ${}^{142,150}$ Nd. <sup>144</sup>Sm systems which are shown in Fig. 1.



FIG. 1: ER cross-section for  ${}^{48}\text{Ti}+{}^{142,150}\text{Nd}$ . <sup>144</sup>Sm systems as a function of  $E_{cm}/V_b$ .

Here, ER cross-section for deformed <sup>150</sup>Nd target ( $\beta_2=0.27$ ) is higher in comparison to the spherical <sup>142</sup>Nd ( $\beta_2=0.09$ ) and <sup>144</sup>Sm  $(\beta_2=0.087)$  targets.

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<sup>\*</sup>Electronic address: priya.apr250gmail.com



FIG. 2: Experimental ER cross-section (filled circle) versus energy in lab frame  $(E_{lab})$  along with theoretical results obtained using statistical model calculations after varying  $K_f$  for a)  ${}^{48}\text{Ti}+{}^{150}\text{Nd}$  b)  ${}^{48}\text{Ti}+{}^{142}\text{Nd}$  and c)  ${}^{48}\text{Ti}+{}^{144}\text{Sm}$  systems including shell-correction in the level density and fission barrier.

# Theoretical calculations

Statistical model (SM) calculations are performed using Bohr-Wheeler (BW) formalism [6] for fission width  $(\Gamma_f)$  including shell correction in the level density and fission barrier. A scaling factor  $(K_f)$  for the FRLDM barrier is introduced and treated as an adjustable parameter to fit the experimental ER cross-sections. Coupled channel calculations were performed to obtain CN spin distributions from the CCFULL code [7] and used as an input in the SM code. SM calculations were subsequently performed and  $K_f$  was adjusted to reproduce the ER cross-sections in the measured energy range. The final theoretical cross-section values are compared with the experimental ones after varying  $K_f$  values are shown in Fig. 2. It is observed that a single value of  $K_f$  cannot fit the ER crosssections over the entire range. For  $^{142,150}$ Nd targets,  $K_f = 0.80$  and for <sup>144</sup>Sm  $K_f > 1$  gives fit to the ER cross-section. The smaller values of scaling factor corresponds to the higher contribution of non compound nuclear (NCN)

processes. The results shows that quasi-fission (QF) is either absent or very small for <sup>144</sup>Sm target and it supports the earlier measurements mentioned in Ref. [8]. In case of Nd targets, smaller  $K_f$  signifies the contribution from QF.

## References

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