

# Beam energy and centrality dependence of particle production in heavy ion collision at RHIC and LHC

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## Introduction

The primary goal of colliding heavy-ions at ultra-relativistic energies is to study nuclear matter under extreme conditions, in which hadronic matter is expected to undergo a phase transition to a new state of matter, the Quark-Gluon Plasma (QGP). In this paper is to find a function which can predict the charged particle multiplicity at any higher energy at a given centrality. There are two main mechanism for particle production that explain well the multiplicity density: Glauber:

$$\frac{dN_{ch}}{d\eta}_{AA} = \frac{dN_{ch}}{d\eta}_{pp} [xN_{coll} + (1-x)\frac{N_{part}}{2}] \quad (1)$$

&  
CGC:

$$\frac{dN_{ch}}{d\eta} \approx N_{part}^\alpha \sqrt{s_{NN}}^\gamma \quad (2)$$

## Motivation

The charged particle multiplicity depends on the beam energy  $\sqrt{S}$  as well as the number of participants nucleon [1]  $N_{part}$ .

$$\frac{2}{N_{part}} \frac{dN_{ch}}{d\eta} = f(S) * g(N_{part}) \quad (3)$$

The energy and centrality dependence of multiplicity normalised with the number of participant pair are given by[1],

$$f(S) = 0.0147 \ln(S)^2 + 0.6 \quad (4)$$

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$$g(N_{part}) = 1 + 0.095 N_{part}^{1/3} \quad (5)$$

for Au+Au collision. For Cu+Cu, the data are seen to be well fitted by,

$$g(N_{part}) = 1 + 0.129 N_{part}^{1/3} \quad (6)$$

f(S) is same as previous. The plot is shown in Fig.1 for Au+Au, Cu+Cu and Pb+Pb. It is seen that the Pb+Pb data do not match well with such functional form. As we go towards higher and higher  $N_{part}$ , the data deviate from the fit function more and more.

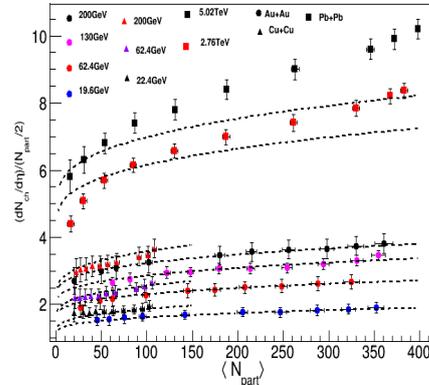


FIG. 1: Particle multiplicity normalised with participant pair as a function of number of participants

## Analysis

We have to find a functional form that contains both energy and centrality, that matches with all the energies to a good extent. The Central Limit Theorem predicts that, Mean multiplicity  $\mu$  i.e.  $\frac{dN_{ch}}{d\eta} \propto N_{part}$

Therefore we assume that the mean multiplicity has the form,

$$\frac{dN_{ch}}{d\eta} = \alpha * N_{part} + \beta \quad (7)$$

where,  $g(N_{part}) = N_{part}$  and  $\alpha$  &  $\beta$  contain the beam energy. Now we have to find the energy dependence of these two coefficients. Particle production mechanism is of two types: soft scattering or thermal part & hard scattering. The thermal part behaves logarithmically with  $\sqrt{s}$  and hard scatter part shows power law dependency on  $\sqrt{s}$ . We have fitted the  $dN_{ch}/d\eta$  vs.  $\langle N_{part} \rangle$  with Eqn.1 and plot  $\alpha$  &  $\beta$  as a function of beam energies to find their dependency on  $\sqrt{s}$ .

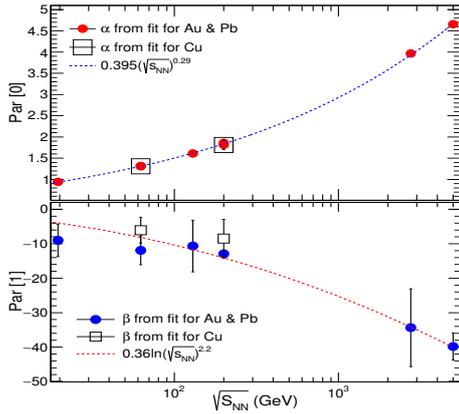


FIG. 2: Energy dependence of the parameters  $\alpha$  and  $\beta$

From the plot, we see that the data can be well described by:

$$\alpha = 0.395(\sqrt{S_{NN}})^{0.29} \quad (8)$$

&

$$\beta = -0.36(\ln\sqrt{S_{NN}})^{2.2} \quad (9)$$

Collision system	$\sqrt{S_{NN}}$ (GeV)	a (func)	a (fit)	b (func)	b (fit)
Au+Au	19.6	0.936	0.937	-3.96	-9.12
Au+Au	62.4	1.309	1.310	-8.17	-11.9
Au+Au	130	1.620	1.604	-11.7	-10.67
Au+Au	200	1.836	1.848	-14.10	-12.94
Cu+Cu	62.4	1.309	1.310	-8.17	-6.06
Cu+Cu	200	1.836	1.809	-14.10	-8.50
Pb+Pb	2760	3.931	3.957	-34.18	-34.45
Pb+Pb	5020	4.675	4.652	-40.12	-39.84

Table shows the comparison of the values of the parameters obtained from the function and fit.

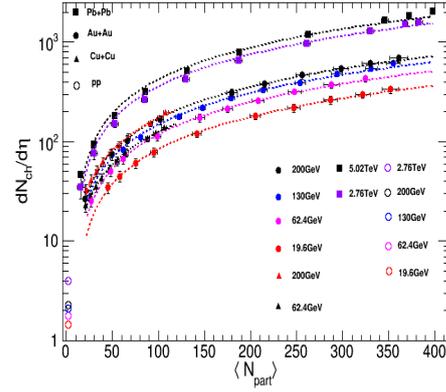


FIG. 3:  $dN_{ch}/d\eta$  vs  $N_{part}$

Now we find the parameters for different energy and fit the  $dN_{ch}/d\eta$  vs  $\langle N_{part} \rangle$  with eqn.(1). The plot is shown in Fig.3.

## Summary

This work could be verified in the d-Au and p-Pb collisions for 200 GeV and 5020 GeV. It could be studied further that eqn [7-9] for finding multiplicity density and Bjorken energy density.

## References

- [1] B. Alver *et al.* [PHOBOS Collaboration]
- [2] A. Adare *et al.* [PHENIX Collaboration], Phys. Rev. C **78** (2008) 044902
- [3] S.Basu *et al.* Phys. Rev. C **93** (2016) 064902