

## De-excitation of $^{210}\text{Po}$

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De-excitation of the compound nucleus (CN)  $^{210}\text{Po}$ , formed in heavy ion-induced fusion reactions, has received considerable attention in the recent years [1–4]. The roles of shell effects and fission dynamics in reproducing the experimental fission cross sections ( $\sigma_{\text{fiss}}$ ) and pre-scission neutron multiplicities ( $\nu_{\text{pre}}$ ) have been examined in these works. Simultaneous reproduction of different observables though remains a challenge.

In this contribution, we present results of a statistical model (SM) calculation for decay of the CN  $^{210}\text{Po}$ . Shell effects in both the level density (LD) and the fission barrier ( $B_f$ ), the orientation degree of freedom of the CN spin ( $K_{\text{or}}$ ) and the collective enhancement in level density (CELD) are considered in the model. The effect of dissipation ( $\beta$ ) in fission dynamics is also included in the calculation.

We consider here population of  $^{210}\text{Po}$  via four different entrance channels, *viz.*, (a)  $^1\text{H}+^{209}\text{Bi}$  [5–8], (b)  $^4\text{He}+^{206}\text{Pb}$  [9–11], (c)  $^{12}\text{C}+^{198}\text{Pt}$  [1, 12, 13], and (d)  $^{18}\text{O}+^{192}\text{Os}$  [13–15]. Measured evaporation residue (ER) cross sections ( $\sigma_{\text{ER}}$ ),  $\sigma_{\text{fiss}}$  and  $\nu_{\text{pre}}$  are compared with our SM [16] predictions.

The fission barrier is obtained by including shell correction in the liquid-drop nuclear mass [3, 17]. The effect of  $K$ -degree (component of CN spin along the nuclear symmetry axis) of freedom [19], is added to the Bohr-Wheeler fission width [18] ( $\Gamma_{\text{BW}}$ ) as:

$$\Gamma_f(E^*, J) = \Gamma_{\text{BW}}(E^*, J, K=0) \frac{(K_0\sqrt{2\pi})}{2J+1} \text{erf}\left(\frac{J+1/2}{K_0\sqrt{2}}\right) \quad (1)$$

with  $K_0^2 = \frac{\tau_{\text{eff}}}{\hbar^2} T_{\text{sad}}$ , where  $\tau_{\text{eff}}$  is the effective moment of inertia.  $\frac{1}{\tau_{\text{eff}}} = \frac{1}{\tau_{\parallel}} - \frac{1}{\tau_{\perp}}$ ,  $\tau_{\perp}$  and  $\tau_{\parallel}$  are the moments of inertia of the nucleus

perpendicular to and about the nuclear symmetry axis, at the saddle.  $\text{erf}(x)$  is the error function. The effects of CELD is included following the work of Zagrebaev *et al.* [20] where the enhancement changes from vibrational ( $K_{\text{vib}}$ ) to rotational ( $K_{\text{rot}}$ ) type with increasing quadrupole deformation of a nucleus. The level density would then be  $\rho(E^*) = K_{\text{coll}}(E^*)\rho_{\text{intr}}(E^*)$  where  $K_{\text{coll}}$  is the enhancement factor and  $\rho_{\text{intr}}(E^*)$  is the intrinsic level density.

Results of our calculation (continuous black lines) are presented in Fig. 1. It is clear that, a suitable value of the dissipation coefficient has to be taken into consideration to reproduce  $\sigma_{\text{ER}}$  and  $\sigma_{\text{fiss}}$ . Hindrance in fission is introduced using the Kramers-modified fission width [21]:

$$\Gamma_K = \Gamma_f \left\{ \sqrt{1 + \left(\frac{\beta}{2\omega_s}\right)^2} - \frac{\beta}{2\omega_s} \right\} \quad (2)$$

where,  $\beta$  is the reduced dissipation coefficient and  $\omega_s$  is the local frequency of a harmonic oscillator potential.  $\Gamma_f$  is the Bohr-Wheeler fission width obtained with shell corrected level densities, CELD and  $K$ -orientation. A value of  $\beta = 3 \times 10^{21} \text{ s}^{-1}$  is required here to reproduce  $\sigma_{\text{fiss}}$  and  $\sigma_{\text{ER}}$  for all the reactions (dashed magenta lines). Measured  $\nu_{\text{pre}}$ , though, is not reproduced except for the most symmetric reaction.

The calculated values of  $\nu_{\text{pre}}$  include the neutrons during saddle-to-scission evolution ( $\nu_{\text{ss}}$ ) of the CN [22]. For a reaction induced by a very light projectile (*e.g.*  $^1\text{H}+^{209}\text{Bi}$ ), the CN spin is small and the saddle is close to the scission. Consequently,  $\nu_{\text{ss}}$  is a small fraction of  $\nu_{\text{pre}}$  and  $\nu_{\text{pre}}$  is reproduced well by the SM. In case of heavier projectiles, CN spin becomes larger and hence saddle-to-scission evolution takes place over a longer duration. This makes emission of more neutrons possible causing enhancement of  $\nu_{\text{ss}}$  and hence  $\nu_{\text{pre}}$ .

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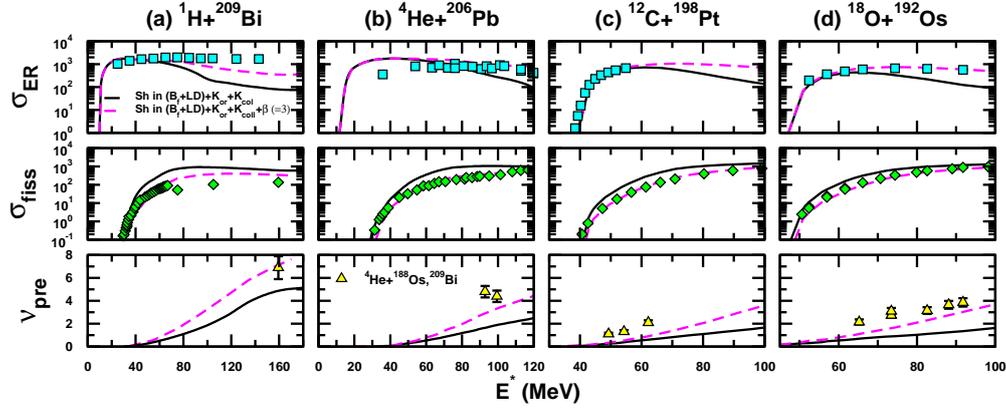


FIG. 1: Measured and calculated  $\sigma_{ER}$ ,  $\sigma_{fiss}$  and  $\nu_{pre}$  for (a)  $^1\text{H}+^{209}\text{Bi}$ , (b)  $^4\text{He}+^{206}\text{Pb}$ , (c)  $^{12}\text{C}+^{198}\text{Pt}$  and (d)  $^{18}\text{O}+^{192}\text{Os}$ . Continuous lines indicate predictions of SM without including dissipation ( $\beta = 0$ ). Dashed lines indicate SM predictions with a suitable value of  $\beta$ .

The present results indicate that a stronger dissipation in the saddle-to-scission region is required in order to reproduce the experimental  $\nu_{pre}$ . Such a shape-dependent dissipation has been reported earlier from both phenomenological studies [23] and also from theoretical considerations [24].

It has also been pointed out that neutrons can be emitted during the comparatively longer formation stage of the CN [25] for reactions induced by heavier projectiles. A suitable model for decay of the CN in which fission is treated dynamically appears to be more appropriate to describe such reactions.

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