

## Thermalized hadron gas of different sizes approach common mean free path at high temperature

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### Introduction

After the discovery of thermalized partonic medium at in  $AuAu$  collisions at  $\sqrt{s_{NN}} = 130$  and 200 GeV at the Relativistic Heavy Ion Collider (RHIC), collisions at lower  $\sqrt{s_{NN}}$  and among comparatively lighter nuclei, also, revealed the QGP-like signals. On the other hand, the Large Hadron Collider (LHC), by unrevealing the long-range two-particle angular correlations in high-multiplicity  $pp$  events at  $\sqrt{s} = 7$  and 13 TeV [1–4] indicates the formation of collective medium in the high-multiplicity  $pp$  events. The possibility of formation of the collective medium in high-multiplicity  $pp$  events is yet to emerge as a consensus view and the main issue of disagreement lies with the question of thermalization in a small system of short lifetime that is supposed to be formed in the  $pp$  events. In this scenario, we present a study on thermalized hadron gas of different finite sizes, which may represent thermalized systems of final state particle production, cre-

ated in ultra-relativistic collisions of different combinations of projectiles and at different energies. We adopt Hagedorn mass spectrum [5] over and above the experimentally measured hadrons and resonance states in the ideal HRG model [6] and include the finite size of the constituents hadrons [7] and study the effect of the finite system-size [8].

### HRG Model with Hagedorn mass spectrum and EV effect

Considering the grand canonical system, with Boltzmann approximation, the thermodynamical variables, pressure ( $P(T)$ ), energy density ( $\epsilon(T)$ ) of hadron resonance gas including the Hagedorn mass spectrum, given by [5]:

$$\rho(m) = C \frac{\theta(m - M_0)}{(m^2 + m_0^2)^a} \exp\left(\frac{m}{T_H}\right) \quad (1)$$

can be written [6], at zero chemical potential, as:

$$P^H(T) = \frac{T}{2\pi^2} \int dm \int_0^\infty p^2 dp \exp\left(-\frac{\sqrt{m^2 + p^2}}{T}\right) \left[ \sum_i g_i \delta(m - m_i) + \rho(m) \right] \quad (2)$$

$$\epsilon^H(T) = \frac{1}{2\pi^2} \int dm \int_0^\infty p^2 dp \sqrt{m^2 + p^2} \exp\left(-\frac{\sqrt{m^2 + p^2}}{T}\right) \left[ \sum_i g_i \delta(m - m_i) + \rho(m) \right] \quad (3)$$

Implementation of the excluded volume effect (EV) [7] in terms of hard core description of constituent hadrons, resonances and the Hagedorn states results into following expressions for the thermodynamic variables [6]:

$P_{EV}^H(T) = \kappa^H P^H(T)$ ,  $\epsilon_{EV}^H(T) = \frac{\kappa^H \epsilon^H(T)}{1 + v \kappa^H n^H(T)}$ , where  $\kappa^H$  is the excluded volume suppression factor and is given by  $\kappa^H = \exp(-v p^{EV}/T)$ .

We implement the finite size effect down to  $R = 2.5$  fm., by introducing finite non-zero

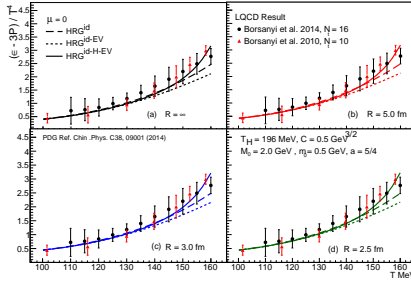


FIG. 1: The trace anomaly as a function of  $T$  for different system sizes of hadron gas with different options, as described in the text.

lower limit of momentum of the particles [9], given by  $p_{cutoff}(in MeV) = 197 \pi/R(in fm.)$ , where  $R$  is the size of a cubic volume. The trace anomaly,  $I(T) = \frac{(\epsilon-3P)}{T^4}$  calculated for different size of the hadron gas, corresponding to  $R = \infty, 5.0, 3.0$  and  $2.5$  fm, have been contrasted with Lattice QCD (LQCD) results for [9, 10] different options of HRG in the figure 1. We have used the mass table provided in the [11]. It is clear from the figure 1, for the hadron gas down to the size corresponding to  $R = 2.5$  fm., follow the LQCD EoS satisfactorily and at high temperature ( $T > 145$  MeV) the LQCD results are better matched with the inclusion of Hagedorn states.

We calculate the mean free path ( $\lambda = 1/n\sigma$ )

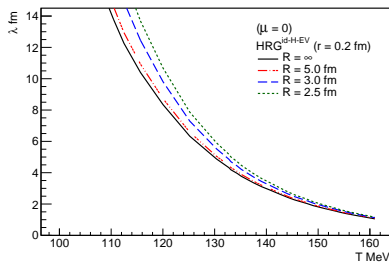


FIG. 2: Temperature dependent pion mean free path in different sizes of hadron gas.

of pions in the thermalized hadron gas of the considered sizes, as a function of temperature

for hadron resonance gas, including the Hagedorn states and corrected with EV effect from the temperature dependent number density and obtaining the pion-pion total cross-sections (from Ref. [12]), which vary from  $\sim 12$  to  $23$  mb in the temperature range from  $T = 100 - 160$  MeV.

## Conclusion

As can be seen in the figure 2, the mean free path for all the considered system-sizes of hadron gas, following the LQCD EoS, approach the universal pion freeze-out value of  $1$  fm., as has been observed [13] in relativistic heavy-ion collisions.

## References

- [1] V. Khachatryan et al., CMS Collaboration, J. High Energy Phys. **09**, 091 (2010).
- [2] G. Aad et al., ATLAS Collaboration, Phys. Rev. Lett. **116**, 172301 (2016).
- [3] V. Khachatryan et al., CMS Collaboration, Phys. Rev. Lett. **116**, 172302 (2016).
- [4] V. Khachatryan et al., CMS Collaboration, arXiv: 1606.06198v1 [nucl-ex] (2016).
- [5] R. Hagedorn, Nuovo Cimento Suppl. **3**, 147 (1995).
- [6] V. Vovchenko, D. V. Anchishkin, and M. I. Gorenstein, Phys. Rev. **C91**, 024905 (2015).
- [7] J. Cleymans and H. Satz, Z. Phys. **C57**, 135 (1993).
- [8] A. Bhattacharyya, R. Ray, S. Samanta and S. Sur, Phys. Rev. **C91**, 041901(R) (2015).
- [9] S. Borsanyi, et. al, J. High Energy Phys. **077**, 1011 (2010).
- [10] S. Borsanyi, et. al, Phys. Lett. **B730**, 99 (2014).
- [11] K. A. Olive et al. (Particle Data Group), Chin. Phys. **C38**, 090001 (2014).
- [12] F. S. Navarra, M. C. Nemes, U. Ornik and S. Pavia, Phys. Rev. **C45**, R2552 (1992).
- [13] D. Adamova et al., CERES Collaboration, Phys. Rev. Lett. **90**, 022301 (2003).