

Ignatyuk damping coefficient - from the haze of fitting to an expression

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Introduction

Level density of a nucleus is widely extracted from transmission data [1] or gamma energy spectrum of the nucleus [2]. It is an important experimental quantity which helps in understanding transitions and reaction cross sections. This extraction goes through lot of rigorous model fittings from codes like CASCADE [3]. An excitation energy-dependent form for fitting the level density parameter was proposed by Ignatyuk [4] as:

$$a(E^*) = \tilde{a} \left[1 + \frac{\delta W}{E^*} \{1 - \exp(-\gamma E^*)\} \right] \quad (1)$$

Here, \tilde{a} is the asymptotic value of nuclear level density parameter, γ is the damping coefficient and δW is the shell correction term which is the difference in the experimental binding energy of nucleus and that calculated from the liquid drop model [9]. This damping parameter has no exact expression and is often employed as a fitting parameter.

In this work we give an exact expression for this damping coefficient, γ , with no adjustable parameters. This is done by using an expression for the level density based on the semiclassical trace formula [5].

Level density parameter and the damping coefficient

For a spherically symmetric harmonic oscillator potential including spin-orbit interactions, finite-temperature single-particle level density $\tilde{g}_T(E)$ has been obtained analytically [5]:

$$\begin{aligned} \tilde{g}_T(E) = & \frac{1 + 3\kappa^2 \hbar^2 \omega_0^2}{\hbar^3 \omega_0^3} \left(E^2 + \frac{\pi^2 T^2}{3} \right) \\ & - \frac{(1 + 5\kappa^2 \hbar^2 \omega_0^2)}{4\hbar \omega_0} + 2\kappa^3 \hbar \omega_0 T \ln \left[1 + \exp \left(\frac{E}{T} \right) \right] \\ & + \dots \end{aligned} \quad (2)$$

where κ is the spin-orbit interaction strength parameter. For nuclei with $N \neq Z$, the chemical potential $\tilde{\mu}^{(n,p)}$ is fixed from proton and neutron number [6]. If we utilize the total finite-temperature level density $g_T(E)$ [5] to find the chemical potential μ 's, then the total level density parameter a_T is given as:

$$a_T = \frac{\pi^2}{6} [g_T(\mu^n) + g_T(\mu^p)]. \quad (3)$$

However, the link between the temperature and the excitation energy $E^* = \sum_{n,p} E_{n,p}^*$ for a nucleus can be found by defining it as difference between the internal energies at finite temperature and at zero temperature [6] i.e.

$$\begin{aligned} E_{n,p}^* &= E_{n,p}(T) - E_{n,p}(T=0) \\ &= \int_0^\infty E g(E) \left[1 + \exp \left(\frac{E - \mu}{T} \right) \right]^{-1} dE \\ &\quad - \int_0^{\epsilon_F^{n,p}} E g(E) dE. \end{aligned}$$

where $g(E)$ is the zero-temperature single-particle level density [7] and $\epsilon_F^{n,p}$ are Fermi energies found using neutron number and proton numbers [6], respectively. Thus, excitation energy dependent level density parameter $a(E^*)$ is found using, $\hbar \omega_0 = \frac{41}{A^{1/3}}$ MeV for nuclei with $N = Z = A/2$ and for $N \neq Z$, we have used [8],

$$\hbar \omega_0(n, p) = \hbar \omega_0 \left(1 \pm \frac{(N - Z)}{A} \right)^{1/3}. \quad (4)$$

Also, the spin-orbit interaction strength parameter κ and the repetitions over the classical periodic orbits k are used as given in [5, 8].

Damping coefficient following from (6),

$$\gamma = -\frac{1}{E^*} \log \left[1 + \frac{E^*}{\delta W} \left(1 - \frac{a(E^*)}{\tilde{a}} \right) \right], \quad (5)$$

is determined and investigated for various nuclei in the present work.

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Results and Discussion

For $^{56,78}\text{Ni}$, $^{208,210}\text{Pb}$, ^{82}Ge and ^{84}Se nuclei, we find the excitation energy dependent level density parameter and then employing (7), we find damping coefficient γ using the parameters as given in Table I.

Nucleus	δW (MeV)	$\kappa_p(\hbar\omega_p)^{-1}$	$\kappa_n(\hbar\omega_n)^{-1}$	k_p	k_n
^{56}Ni	-5.300	$\kappa = -0.05$	$k=7$ ($N=Z$)		
^{208}Pb	-9.923	-0.151	-0.134	10	9
^{210}Pb	-6.945	-0.151	-0.134	10	9
^{78}Ni	-10.842	-0.075	-0.060	5	5
^{82}Ge	-3.803	-0.090	-0.060	8	5
^{84}Se	-1.880	-0.090	-0.060	8	5

TABLE I: Shell correction term δW , spin-orbit interaction strength parameters κ 's and the repetitions over the classical periodic orbits k 's used in our calculations are given in this Table.

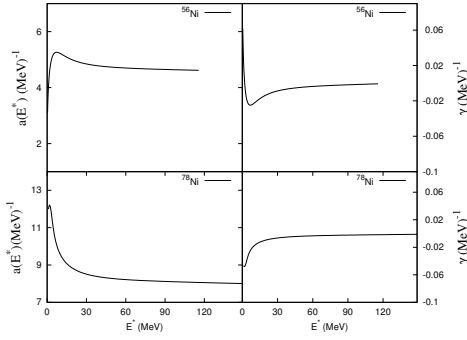


FIG. 1: Excitation energy dependent level density parameter $a(E^*)$ and damping coefficient γ are plotted with respect to the excitation energy for ^{56}Ni and ^{78}Ni nuclei.

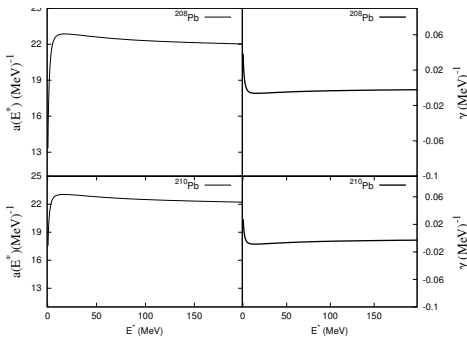


FIG. 2: Excitation energy dependent level density parameter $a(E^*)$ and damping coefficient γ are plotted with respect to the excitation energy for ^{208}Pb and ^{210}Pb nuclei.

It is known [5, 10] that the level density parameter may increase (like in ^{56}Ni and $^{208,210}\text{Pb}$) or decrease (like in ^{78}Ni , ^{82}Ge and ^{84}Se) with the excitation energies, until it attains a saturated value. This leads

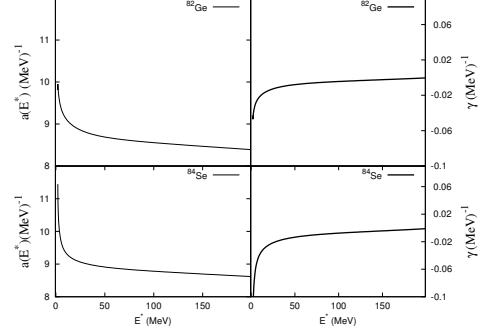


FIG. 3: Excitation energy dependent level density parameter $a(E^*)$ and damping coefficient γ are plotted with respect to the excitation energy for ^{82}Ge and ^{84}Se nuclei.

to both positive and negative values of the damping coefficient until it approaches to zero at saturation of level density. The values of γ obtained here are of the order of experimentally deduced values. Here the values of \tilde{a} are taken from the average semiclassical level density [5], which almost follow a trend of $A/10$. In practice, this behaviour is taken as A/k , where k is a parameter which can vary from 8 to 12 [3]. If we consider this variation of the k parameter as well, the experimental values in the range $0.04 \sim 0.06$ can also be obtained.

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