

## Study of s-state $\Xi$ -hypernuclei in the Hyperspherical Harmonics Expansion approach

M. Alam, M. Hasan, S. H. Mondal, and Md. A. Khan\*

*Department of Physics, Aliah University,  
IIA/27, Newtown, Kolkata-700160, INDIA*

### Introduction

Hypernuclei are the subatomic systems with a remarkable strangeness degree of freedom as compared with conventional nuclei. Such exotic nuclei are formed when one or more strange exotic hyperon(s) (like  $\Lambda, \Sigma, \Xi$  etc) are injected into an atomic nucleus so as to replace one or more nucleon(s) of the nucleus and the injected particles are coupled to the nuclear core[1]. Hypernuclei provide a unique source of information on hyperon-nucleon, hyperon-hyperon interactions and the behaviour of strange particles in baryonic matter which opened an avenue for many significant phenomenon in particle physics and nuclear-astrophysics[2].

Nakazawa et al., 2015[3] first reported the bound state of the p-state  $\Xi^- - ^{14}\text{N}$  hypernuclei with a  $\Xi^-$  separation energy of  $4.38 \pm 0.25$  MeV. Recently, Yoshimoto et al., 2021[4] first identified the nuclear 1s state of the hypernucleus  $^{15}\text{C}$  having  $B_{\Xi^-}$  of IRRAWADDY ( $6.27 \pm 0.27$  MeV) and KINKA ( $8.0 \pm 0.77$  MeV) events. Some theoretical investigations were done to estimate whether the  $\Xi^-$  occupies on 1s state or 1p state in KISO event, such as relativistic mean-field (RMF) model, Skyrme-Hartree-Fock (SHF) model[5] and quark mean-field (QMF) model[6].

The main objective of the present work is to study the ground state in 1s state of  $^{15}\text{C}$  and we will predict  $^{16}_{\Xi^- \Xi^-}\text{C}$  hypernuclei. We will calculate the one- and two- $\Xi^-$  separation energies and some relevant geometrical observables of these systems. In our two-body model, we will first solve Schrödinger equation

numerically to reproduce the single  $\Xi^-$  separation energy of  $^{15}\text{C}$  hypernuclei by adjusting the parameter of the effective core-hyperon potential. We use the  $\Xi$ -nucleus potential together with  $\Xi\Xi$  potential to predict the ground state energy of the  $^{16}_{\Xi^- \Xi^-}\text{C}$  hypernuclei. For the three-body model calculation of double- $\Xi$  system, we employ the hyperspherical harmonics expansion (HHE) method.

### Theoretical formalism

In our HHE approach[7], we label the relatively heavier core nucleus  $^{14}\text{C}$  as the particle “i” and two valence  $\Xi$ -hyperons as particles “j” and “k” form the interacting pair are:

$$\left. \begin{aligned} \vec{\xi}_i &= a_{jk}(\vec{r}_j - \vec{r}_k) \\ \vec{\chi}_i &= a_{(jk)i} \left( \vec{r}_i - \frac{m_j \vec{r}_j + m_k \vec{r}_k}{m_j + m_k} \right) \\ \vec{R} &= \sum_{i=1}^3 \frac{m_i \vec{r}_i}{M} \end{aligned} \right\} \quad (1)$$

where  $\rho = \sqrt{\xi_i^2 + \chi_i^2}$ ;  $\phi_i = \tan^{-1}(\chi_i/\xi_i)$ . The motion of the three-body system in relative coordinates can be described by the Schrödinger’s equation

$$\left[ -\frac{\hbar^2}{2\mu} \left\{ \frac{\partial^2}{\partial \rho^2} + \frac{5}{\rho} \frac{\partial}{\partial \rho} + \frac{\hat{K}^2(\Omega_i)}{\rho^2} \right\} \right] \psi(\rho, \Omega_i) + [V(\rho, \Omega_i) - E] \psi(\rho, \Omega_i) = 0 \quad (2)$$

where  $V(\rho, \Omega_i)$  is the total interaction potential in the partition  $i$  and  $\hat{K}^2(\Omega_i)$  is the square of hyper angular momentum operator satisfying the eigenvalue equation.

### Choice of potential

For the core- $\Xi$  subsystem and  $\Xi\Xi$  pair, we used a two-range Gaussian Isle-type potential and Yukawa-type potential [8] of the form re-

\*Electronic address: drakhan.rsm.phys@gmail.com

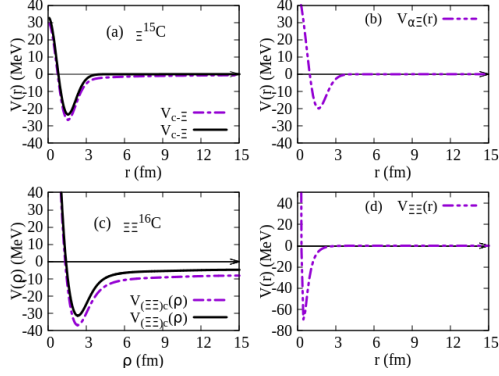


FIG. 1: Representative plot of effective core- $\Xi$  potential components with radial distance are (a) with coulomb and without coulomb (solid line) interaction, (b)  $\alpha - \Xi$  (d)  $\Xi\Xi$  pair and (c) for three-body effective potential with hyper-radial distance  $\rho$ .

spectively

$$V_{c-\Xi}(r) = 450.4 \exp(-r/1.269)^2 - 404.9 \exp(-r/1.41)^2 \quad (3)$$

$$V_{\Xi\Xi}(r) = (-155.0 \exp(-1.75r) + 490.0 \exp(-5.6r))/r \quad (4)$$

A Coulomb interaction is added for the charged pair of interacting particles

$$V_{ij}^{(C)}(r_{ij}) = \begin{cases} 1.44 \left( \frac{Z_i Z_j}{2R_c} \right) \left( 3 - \frac{r_{ij}^2}{R_c^2} \right) ; r_{ij} \leq R_c \\ 1.44 \left( \frac{Z_i Z_j}{r_{ij}} \right) ; r_{ij} > R_c \end{cases} \quad (5)$$

## Results and Discussions

In this work, we first investigate the 1s bound state of  ${}^{15}_{\Xi^-}\text{C}$  and predict double- $\Xi$  hypernucleus  ${}^{16}_{\Xi^-\Xi^-}\text{C}$  using HHE method which is an essentially exact method, where calculations can be carried out up to any desired precision by gradually increasing the basis.

## Acknowledgments

Authors would like to express an appreciation of necessary facilities from the Aliah university and UGC.

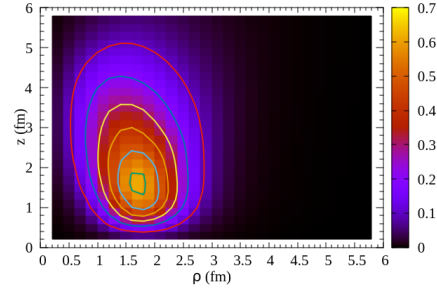


FIG. 2: Contour plot reflects probability density distribution of nucleons and  $\Xi^-$  hyperons in the ground state of  ${}^{16}_{\Xi^-\Xi^-}\text{C}$ .

TABLE I: (I) Parameters of the attractive component of the core- $\Xi^-$  potential and  $B_{\Xi^-}$  separation energy of  ${}^{15}_{\Xi^-}\text{C}$ . (II) Double- $\Xi^-$  separation energy and r.m.s raddi of  ${}^{16}_{\Xi^-\Xi^-}\text{C}$ . The bold face parameters obtained by switch off the coulomb interaction term. Energies and potentials are in MeV and raddi are in fm unit.

(I)	$V_{att}$	$V_{c-\Xi^-}$ <i>Min.</i>	$B_{\Xi^-}$ Cal.	$B_{\Xi^-}$ Expt.
${}^{15}_{\Xi^-}\text{C}$	-417.18	-26.3556	8.0012	$8.0 \pm 0.77$
		<b>-23.2811</b>	<b>5.1671</b>	
(II)	$B_{\Xi^-\Xi^-}$	$R_A$	$R_{c\Xi^-}$	$R_{\Xi^-\Xi^-}$
${}^{16}_{\Xi^-\Xi^-}\text{C}$	-20.543	2.581	2.198	2.716
	<b>-15.280</b>	<b>2.585</b>	<b>2.229</b>	<b>2.741</b>

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