

Phase-Equivalent Potential Construction by SUSY Transformations

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INTRODUCTION

It is known that by the judicious use of the algebra of supersymmetry, a potential can be constructed with bound states at arbitrary energies. The SUSY transformation, a mathematical transformation which falls under SUSY algebra, is used to generate phase equivalent potentials with some desired alterations. There are 4 different types of SUSY transformations, using which, a supersymmetric partner to the Hamiltonian of a given radial Schrödinger equation can be formed. Using three pairs of successive SUSY transformation (T1T2, T1T3, and T1T4) on Manning-Rosen potential [2], three new phase-equivalent potentials namely V12; V13; and V14 are successfully constructed. As a model calculation we shall apply them to n-p, p-p and α -p systems.

Construction of Phase equivalent potential:

The three parameter short ranged Manning-Rosen potential [2]

$$V_m = b^{-2} \left[\frac{\alpha(\alpha-1)}{(1-e^{-r/b})^2} e^{-2r/b} + \frac{l(l+1)}{(1-e^{-r/b})^2} e^{-2r/b} - A \frac{e^{-r/b}}{(1-e^{-r/b})} \right] \quad (1)$$

which is generally used for the study of molecular systems, have been parameterized to fit nuclear scattering phase parameters. The Jost solution for the potential under consideration reads

$$f_l(kr) = (1 - e^{-r/b})^{1+p} e^{-wr} \quad (2)$$

Under the pairs of SUSY transformations, three new potentials are obtained as

$$V_{12}(r) = V_m(r) - \frac{d^2}{dr^2} \ln \left(1 + \lambda \int_0^r f_l^2(x) dx \right) \quad (3)$$

$$V_{13}(r) = V_m(r) - \frac{d^2}{dr^2} \ln \left(\int_0^r f_l^2(x) dx \right) \quad (4)$$

and

$$V_{14}(r) = V_m(r) - \frac{d^2}{dr^2} \ln \left(\int_r^\infty f_l^2(x) dx \right). \quad (5)$$

Where b, A and α are three adjustable parameters of Manning-Rosen potential. The quantity λ is an adjustable parameter for transformation operations, $w = -ikr$; and $p =$

$$\frac{1}{2} \left[1 \pm \sqrt{1 + 4\{\alpha(\alpha-1) + l(l+1)\}} \right] - 1.$$

For p-p and α -p systems the electromagnetic interaction part is provided by addition of atomic Hulthén potential which acts as a Screened Coulomb potential:

$$V_H(r) = \frac{V_0 \exp(-r/a)}{(1 - \exp(-r/a))}, a > 0, \quad (6)$$

where V_0 is strength and a is screening radius.

We generate phase shifts for all the aforementioned states via phase function method, and compare the findings with standard data.

Results and Discussion:

For n-p, p-p system we have used
 $\frac{\hbar^2}{2m} = 41.47 \text{ Mev fm}^2$ and for α -p
 systems $\frac{\hbar^2}{2m} = 25.92 \text{ Mev fm}^2$

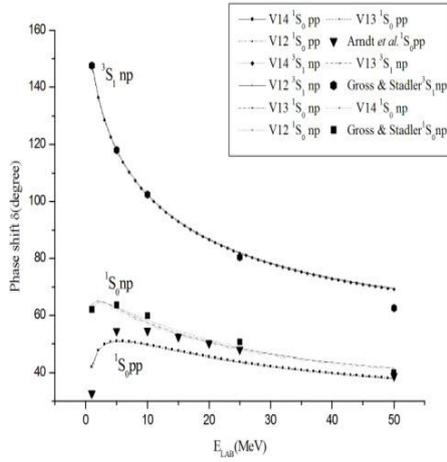


Fig 1. Graph for n-p; p-p Systems for V14, V13 and V12 potentials with reference data (1S_0 & 3S_1 states).

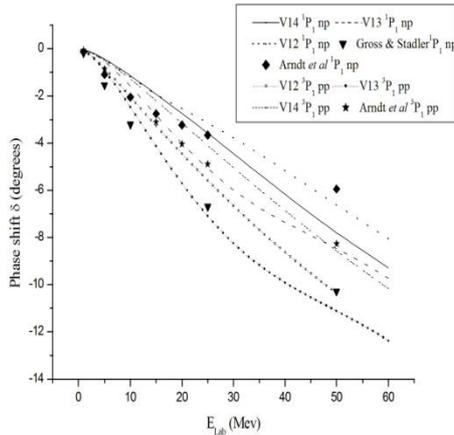


Fig 2. Graph for 1P_1 n-p & 3P_1 p-p phase shifts.

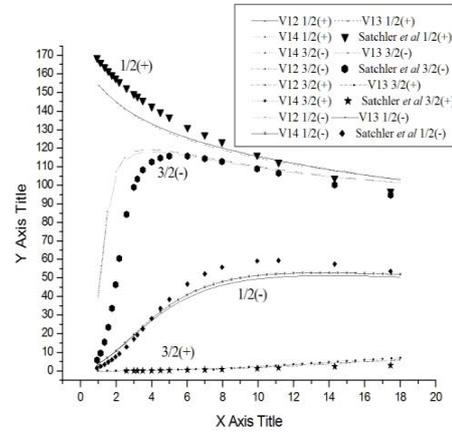


Fig 3. Graph for $1/2(+)$, $1/2(-)$, $3/2(+)$ and $3/2(-)$ α -p phase-shifts.

The 1S_0 & 3S_1 states for n-p system and 1S_0 state for p-p system are in good agreement with standard data [3] & [4] up to 50 MeV. For α -p system $1/2(+)$, $1/2(-)$, $3/2(+)$ and $3/2(-)$ States are in good agreement with experimental data [5]. The phase parameters for 1P_1 n-p and 3P_1 p-p states follow similar trends with reference data [3] and [4] with slight deviations in their numerical values.

Thus, we conclude that the SUSY transformations work to some good extent for generation of various phase equivalent potentials.

Reference :

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