

Fractality and phase transition studies for π^+ and π^- in $p - p$ collisions at $\sqrt{s} = 13$ TeV

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Introduction

From relativistic heavy ion collisions under extreme conditions of temperature and/or density, we get the valuable informations about the the spatiotemporal progression of the multiparticle production process. The existence of QGP like state of nuclear matter for nucleus–nucleus collisions at relativistic and ultra-relativistic energies was confirmed by the experiment at SPS at CERN, the RHIC at BNL and LHC at CERN. Now, $p - p$ collisions can be considered as the promising probe to know the underlying dynamics of the heavy-ion collisions. In this study we

different properties of charged particle multiplicities. It is also very true that the UrQMD simulation has gained huge popularity to generate the events in $p - p$ collisions. The creation of pions for $p - p$ collisions were investigated through the mean multiplicity, rapidity spectra and particle ratios as a function of collision energy. We have considered the UrQMD model produced data to observe the intermittency type particle dynamical fluctuations of π^+ and π^- separately for $p - p$ collisions, at $\sqrt{s} = 13$ TeV energy in the form of scale factorial moment. To justify the utility of UrQMD model, we have compared p_T spectra of pions for UrQMD events with ALICE experimental data for $p - p$ collisions at $\sqrt{s} = 13$ TeV [2].

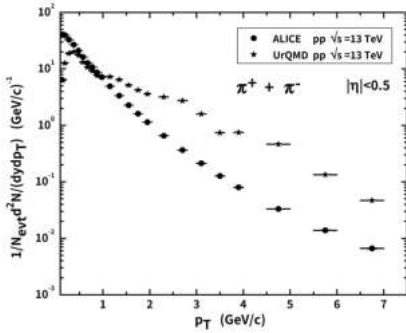


FIG. 1: Comparison of the p_T spectra of pions between the ALICE data and the UrQMD simulated data for p - p interactions at $\sqrt{s} = 13$ TeV.

have used Ultra-relativistic Quantum molecular Dynamics (UrQMD) model [1] to generate the events in $p - p$ collisions and studied the

Method of Analysis

According to Bialas and Peschanski [3] in one-dimensional phase-space which is divided into M number of bins, the factorial moment F_q of order q will be given by
$$F_q = M^{q-1} \sum_{m=1}^M \frac{\langle n_m(n_m-1)\dots(n_m-q+1) \rangle}{n_m^q},$$
 where n_m represents the number of net pions (i.e. sum of all π^+ or π^-) in m^{th} bin and q is the order of moment. Here the bracket $\langle \dots \rangle$ represents the event average.

The anomalous fractal dimension “ d_q ” on intermittency exponent “ α_q ”, can be written as $d_q = \frac{\alpha_q}{q-1}$, where d_q is also known as Renyi Co-dimension. Now the general fractal dimension “ D_q ” is also given by the equation $D_q = 1 - d_q$.

In order to find the degree of multifractality, we have consider the following eqn.

$$\frac{d_q}{d_2} = (1 - r) + \frac{qr}{2},$$

where “ r ” implies the degree of multifractality of the colliding particles. A non-zero values of “ r ” indicate the presence of multifractal behaviour [4]. The existence of non-thermal

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TABLE I: Value of different parameters depending on the scale factorial moment analysis in one dimensional η space for π^+ & π^-

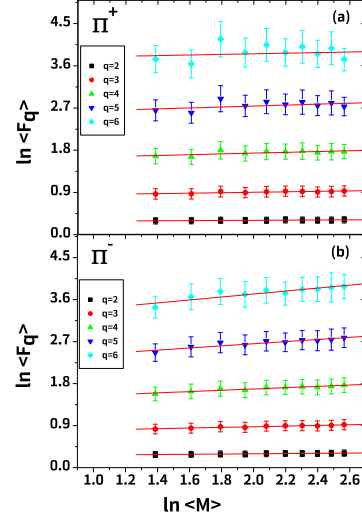
Parameters	π^+	π^-
r	0.622 ± 0.064	0.786 ± 0.066
ν	1.507 ± 0.046	1.582 ± 0.047
μ	1.630 ± 0.104	1.798 ± 0.016
c	0.057 ± 0.037	0.118 ± 0.036

phase transitions can be verified through a function λ_q , where $\lambda_q = \frac{\alpha_q + 1}{q}$. In order to observe the phase transitions, a function β_q is defined as $\beta_q = \frac{d_q}{d_2}(q-1)$. Ginzburg and Landau have also derived a critical exponent ν , which can be calculated by the equation $\beta_q = (q-1)^\nu$. According to Levy law, the ratio of anomalous dimension depends on μ through the relation $\beta_q = \left(\frac{q^\mu - q}{2^\mu - 2}\right)$. Now the multifractal specific heat (c) can be derived through the eqn. $D_q = (a - c) + c\left(\frac{\ln q}{q-1}\right)$, where “ a ” is a constant and “ c ” indicates the transition between multifractality and monofractality for relativistic heavy ion collisions.

Discussions

We have estimated the values of factorial moment for different order q in one dimensional η space using the relation $\langle F_q \rangle \propto M^{\alpha_q}$ which are shown in Fig. 2. We get the intermittency exponent α_q for different order q from the best fitted linear behaviour curve in Fig. 2 for π^+ & π^- separately. We have derived the values of “ r ”, “ ν ”, “ μ ” & “ c ” by using the values of anomalous fractal dimension α_q . From the variation of β_q vs $(q-1)$ we have derived the value of critical exponent “ ν ” and we have observed that for one dimensional η space some of the values of “ ν ” are close to 1.304 and some are not which in accordance to the framework of GL theory agrees quark-hadron phase transitions may exists for the values close to 1.304 and remaining values of “ ν ” doesn’t as they are not close to the given value of “ ν ”. It has been observed that π^- is more multifractal than π^+ .

From the D_q vs $\left(\frac{\ln q}{q-1}\right)$ graph we have de-


 FIG. 2: Variations of $\ln \langle F_q \rangle$ with respect to $\ln M$ in one-dimensional η space (a) for π^+ and (b) π^-

derived the value of multifractal specific heat (c). For 1D η space, the value of “ c ” for π^+ and π^- being greater than zero which signals the presence of the transformations from the multifractality to monofractality for p-p collisions at $\sqrt{s}=13$ TeV.

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References

- [1] M. Bleicher *et al.*, *J. Phys. G* **25**, 1859 (1999).
- [2] ALICE collaboration, <https://arxiv.org/pdf/2003.02394v2.pdf>
- [3] A. Bialas, R. Peschanski, *Nucl. Phys. B*, **308**, 857 (1988).
- [4] S. Bhattacharjee *et al.* *IJMPE* <https://doi.org/10.1142/S0218301322500793>