

## Comparison of ground state $\Lambda$ -binding energies of light hypernuclei

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### Introduction

Earlier, we had obtained a semi-empirical formula [1] for ground state  $\Lambda$ -binding energies ( $B_\Lambda$ ) of light hypernuclei, using an exponential form of  $\Lambda$ -nucleus potential, which gives a fairly good account of  $B_\Lambda$  data. Although, the ground state  $B_\Lambda$  of light hypernuclei have been satisfactorily explained [2,3] using  $\Lambda$ -nucleus potential in the folding model with a suitable density-dependent effective  $\Lambda$ -nucleon ( $\Lambda N$ ) interaction, but these calculations have been done numerically. Hence, in order to obtain a broad insight into some of the important aspects, we do a comparative study of the ground state  $B_\Lambda$  of light hypernuclei obtained by us [1] and others [2-4].

### Semi-empirical formula and some $\Lambda$ -nucleus potentials

We solve the radial Schrödinger equation with the exponential form of the  $\Lambda$ -nucleus potential [1]:

$$V(r) = -V_0 e^{-r/a},$$

$$\text{where } a = r'_0 A_c^{1/3}.$$

Here,  $A_c$  is the core mass number of the hypernuclei and  $r'_0$  is taken as a free parameter. Further simplifying, we get a semi-empirical formula for the  $\Lambda$ -binding energy in bound s-states of light hypernuclei [1]:

$$B_\Lambda = \frac{\hbar^2 A_c^{-2/3}}{8m_\Lambda r_0'^2 ((\pi^2/2) - 1)^2} \{ (\xi_1 + \xi_2) + \left[ (\xi_1 + \xi_2)^2 - (\pi^2 - 2) \left( \frac{\xi_1^2}{\pi^2} + \left( \frac{1}{2} - 2n \right) \xi_2 + \frac{1}{8} \right) \right]^{1/2} \} + \frac{\beta' Z}{r_{rms}^3}, \quad (1)$$

where  $\xi_1 = \frac{\pi^2}{2} ((1/2) - 2n)$  and

$$\xi_2 = \frac{\pi r_0' A_c^{1/3}}{2} \left[ \frac{8m_\Lambda V_0}{\hbar^2} \right]^{1/2}.$$

This semi-empirical formula (eq. 1), gives a fairly good account of the ground state  $B_\Lambda$  of light hypernuclei, if very light nuclei are ignored. These  $B_\Lambda$  values are shown as 1 in Fig. 1. The  $B_\Lambda$  of  ${}^3_\Lambda\text{H}$ ,  ${}^4_\Lambda\text{H}$  and  ${}^4_\Lambda\text{He}$  are the predicted values. We have also plotted  $B_\Lambda$  values [1] obtained from  $\chi^2$ - fitting using eq.(1) in which the last term is replaced by  $\beta' \left( \frac{Z}{A_c} \right)$ . The results are quite reasonable when very light hypernuclei,  ${}^3_\Lambda\text{H}$ ,  ${}^4_\Lambda\text{H}$  and  ${}^4_\Lambda\text{He}$  and the known 'troublesome' nuclei,  ${}^7_\Lambda\text{Be}$  and  ${}^9_\Lambda\text{Be}$ , are excluded from the fit. The  $B_\Lambda$  of these excluded nuclei are the predicted values. These fitted and the predicted values are represented as 2 in Fig. 1.

We also carry out exact numerical fitting of ground state  $B_\Lambda$  data of light hypernuclei with core mass number ranging from 2 to 15, using exponential form of  $\Lambda$ -nucleus potential. The calculated  $B_\Lambda$  for hypernuclei with  $A_c = 3, 4$  and  $6$  are the predicted values. These results are shown as 3 in Fig. 1.

Ahmad et al. [2], have calculated the  $B_\Lambda$  values for light hypernuclei in the folding model with a density-dependent effective  $\Lambda N$  interaction. The  $\Lambda$ -nucleus potential [2] is

$$V_\Lambda(r) \approx \bar{V}_0 \tilde{\rho}(r) [1 - \beta \rho^{2/3}(r)],$$

where  $\tilde{\rho}(r)$  is assumed to be the same as the charge distribution of the core nucleus.  $\bar{V}_0$  and  $\beta$  are treated as parameters. The numerically calculated  $B_\Lambda$  values, plotted as 4 in Fig. 1, correspond to their potential parameter set B [2].

Mahmood Mian [3] using a three-parameter density-dependent  $\Lambda N$  force in the folding model, obtained the  $\Lambda$ -nucleus potential as

$$V_\Lambda(r) = \frac{\bar{V}_0}{2\pi^2} \int F(q) \exp \left[ -\frac{q^2 d^2}{4} \right] q^2 j_0(qr) dr,$$

where the symbols have their usual meaning. The numerically calculated  $B_\Lambda$  values [3] in the

ground state of light hypernuclei, corresponding to the potential parameter set A [3], are plotted as 5 in Fig. 1. The average size parameter of the nuclear density for  $N \neq Z$  core hypernuclei was also determined [3] by numerically fitting their experimental  $B_\Lambda$  data. The theoretical  $B_\Lambda$  values so obtained [3] are plotted as 6 in Fig. 1.

Employing a free-state two-body  $\Lambda N$  gaussian potential along with a  $\delta$ -function three-body ANN force, the  $\chi^2$ -fit to the experimental  $B_\Lambda$  data was performed by Ansari et al. [4]. The corresponding  $B_\Lambda$  values obtained for their set AII [4] are plotted as 7 in Fig. 1.

### Result and Discussion

The calculated ground state  $B_\Lambda$  values of light hypernuclei are plotted versus  $A_c$  in Fig. 1, along with the available experimental  $B_\Lambda$  data. As the error bar in the experimental  $B_\Lambda$  data of  ${}^{16}_\Lambda\text{O}$  is not available, we take a plausible value of 5% of the experimental  $B_\Lambda$  as the error in the datum. The  $B_\Lambda$  values calculated from our semi-empirical formula (eq. 1) as well as from the numerical fitting using an exponential potential, give a fairly good agreement with the experimental  $B_\Lambda$  data. Predicted  $B_\Lambda$  values, though slightly lower for  ${}^4_\Lambda\text{H}$  and  ${}^4_\Lambda\text{He}$  are also qualitatively acceptable. When the last term of  $B_\Lambda$  formula (eq. 1) is expressed as  $\beta' \left(\frac{Z}{A_c}\right)$ , the formula distinguishes between nuclei with different  $Z$  but same  $A_c$ , though the size effect is ignored. The results, however, are more or less similar except for some differences at higher  $A_c$ .

The  $B_\Lambda$  values of Ahmad et al. [2], where  $\tilde{\rho}(r)$  was assumed to be the same as the charge distribution of the core nucleus, are represented as 4 in Fig. 1. The calculated  $B_\Lambda$  values for  $A_c = 6, 7, 15$  and  ${}^9_\Lambda\text{B}, {}^{10}_\Lambda\text{B}, {}^{12}_\Lambda\text{C}$  are obtained with potential parameter set B as given in ref. [2]. The available ground state  $B_\Lambda$  is satisfactorily reproduced with slight difference in few cases, especially beyond  $A_c = 11$ .

The  $B_\Lambda$  values obtained by Mahmood Mian [3], taking into account the range of  $\Lambda$ -nucleon interaction explicitly, referred to as 5 in Fig. 1, show deviation at higher  $A_c$  ( $>12$ ). Another set of  $B_\Lambda$  values determined [3], to obtain average size parameter of the nuclear density for  $N \neq Z$  core hypernuclei, referred to as 6 in Fig. 1, are

somewhat on the lower side of the experimental data at higher  $A_c$ .

Ansari et al. [4] calculated  $B_\Lambda$  values from  $\chi^2$ -fit to the experimental  $B_\Lambda$  data of hypernuclei with core mass number ranging from 9 to 14. Their results, represented as 7 in Fig. 1, show slight over binding for  $A_c = 13$  and 14.

The results obtained in refs. [2-4] give a reasonable account of the data, except for some deviations at higher  $A_c$ . Fairly satisfactory reproduction of ground state  $B_\Lambda$  of light hypernuclei obtained by our semi-empirical formula [1] and from our numerical fitting of the  $B_\Lambda$  data indicates that even an exponential  $\Lambda$ -nucleus potential is quite adequate in reproducing  $B_\Lambda$  of many light hypernuclei. Moreover, it leads to a semi-empirical formula which gives fairly good results. However, more work is needed to analyze whether any deeper significance can be ascribed to the success of exponential potential.

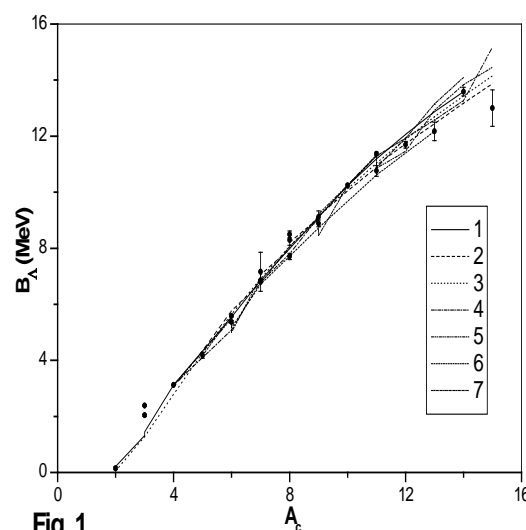


Fig. 1

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