

Prediction of temperature limit for hot rotating magic SHN

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Introduction

The survival of the nuclei at different physical conditions is one of the very basic criteria in the field of nuclear physics. Many nuclei especially far from stability are expected to be formed during supernova explosions either through r- or rp-process. After the discovery of $Z \leq 118$, formation and detection of nuclei with bigger Z is one among the main objectives of nuclear physicists. It is known that nuclei featuring magic nos. of nucleons are more stable than their neighbors. The relativistic Hartree-Bogolyubov theory determined the possible proton magic nos., as $Z=106, 114, 120, 126, 132$ & 138 and neutron magic nos. as $N=138, 172, 184, 198, 216, 228, 238, 252, 258$ & 274 [1]. In this paper we discuss about the possible existence of the predicted superheavy nuclei with proton magic numbers 132 and 138, in the vicinity of stellar nucleosynthesis, i.e., the maximum temperature of nuclear existence.

Methodology

The nuclei in the SHN region may be formed during supernova explosions through r-process and hence the formed system must be treated as a thermo dynamical one. So the statistical theory incorporating temperature and rotation is more appropriate one. In this work we have followed the statistical theory by treating the nuclei as a hot system. The statistical code developed by us is executed for the numerical values, such as, $T=0.1-2.0$ MeV (0.1); spin $J=0-60\hbar$; $\gamma = -180^\circ$ to -120° (10°); and $\delta=0.0-0.6$. The BE/A is calculated using the droplet model [2] and obtained a higher value for $N=200$ for $Z=132$ and $N=210$ for $Z=138$ as $E_B=6.77158$ MeV and 6.64661 MeV respectively. Comparing the said two nuclei $Z=132$ is always having higher BE/A for the range of $N=170-188$, but high E_B alone does not provide stability to

the nucleus as for as the nuclei synthesized in the chemical evolution of stellar nuclear processes.

Results and Discussion

Z=132: At $T=0.5$ MeV the neutron separation energy (S_n) decreases smoothly with increases of spin for the nuclei with $N=172-178$, but when $N=180$, the S_n shows an increasing trend. S_n has the maximum value of 16.15 MeV at $0\hbar$ in this range of nuclei. While increasing the temperature S_n value get decreases as neutron exhibits local minima at $N=164$ and around $1-1.8$ MeV at $T=1.0$ MeV, and 2 MeV at higher temperatures. Such a fluctuation of S_n at low spins at $T=1.0$ MeV pronounces the maximum possible temperature of the system may be 1.0 MeV.

The l_{dp} is a smooth Gaussian for $N=172-186$ at $T=0.5$ MeV, but at $N=188$, a shift at $44\hbar$ and another at $62\hbar$ is observed. At higher temperatures, i.e., $1.0-2.0$ MeV a shift at $J>20\hbar$ is obtained for $N=172-178$. Nucleus with $N=180$ shows a shift at $J=10\hbar$ and the corresponding shift at a particular temperature increases slowly for higher N , Such an effect is observed in E_x also, and which resembles in the shape of the nucleus studied.

It is oblate deformed ($\gamma=-180^\circ$, $\delta=0.1$) at low spins for $N=172-178$ at $T \leq 1.0$ MeV. At higher spins the deformation is increased from $\delta=0.1$ to 0.2 . But at $T>1.0$ MeV the system behaves like a spherical one at low spins ($J<30\hbar$) and oblate deformed at higher spins. If N is added i.e., $N=180-186$, a spherical shape is predicted at $T=1.0$ MeV. At $T=0.5$ MeV, an oblate shape is predicted irrespective of the spins but the deformation increases with increase of spin. At higher temperatures, $T>1.0$ MeV, all the nuclides studied show a spherical shape at lower spins, ($J=0-30\hbar$), and become oblate deformed.

At $J=0\hbar$, a drastic difference in shape is observed for $N>160$ when $T \leq 0.5$ MeV, i.e., for

$N < 210$: $\gamma = -180^\circ$, $\delta = 0.2$ and $\gamma = -120^\circ$, $\delta = 0.2$ for $N < 232$. For $N = 234$: $\gamma = -180^\circ$, $\delta = 0.1$ and $210 < N < 230$: $\gamma = -160^\circ / -140^\circ$, $\delta = 0.1$. Spherical nuclei are observed for $N = 180-198$, $N = 182-206$, and $N = 180-208$ at $T = 0.7 \text{ MeV}$, 0.9 MeV and 1.0 MeV respectively. At $T > 1.0 \text{ MeV}$ a spherical shape is observed irrespective of N .

Z=138: At $T = 1.0 \text{ MeV}$, a drop in S_n is observed at low spins and which increases further for $N > 198$. But a continuous decreasing effect is observed when $238 < N < 198$. And at $T > 1.0 \text{ MeV}$ the S_n seems to be constant. i.e., the change in S_n is too small. Such a trend in l_{dp} also observed at $T > 1.0 \text{ MeV}$, the trend of the l_{dp} is entirely different at $T < 1.0 \text{ MeV}$ and at this temperature a fluctuating trend is observed when $238 > N > 198$ and which is a smooth Gaussian for $238 < N < 198$.

The nucleus $Z=138$ shows an oblate shape ($\gamma = -180^\circ$; $\delta = 0.1/0.2$) at $T = 0.5 \text{ MeV}$ when $N = 170-190$ from spin $J = 0-60\hbar$. Increasing N shows a shape fluctuation with spin, i.e., $N = 192, 194$ is spherical at low spin and becomes oblate. $N = 196, 198$ are oblate at $J < 12\hbar$ and become spherical at $14\hbar < J < 30\hbar$ then to oblate. A triaxial shape is predicted at low spins for $N = 200-236$ and its shape fluctuates from prolate to oblate via triaxial with spin. An oblate shape with deformation $\delta = 0.2$ is predicted for $N > 238$ at all spins. When $T = 1.0 \text{ MeV}$, $N = 160-164$ are prolate at lower spins and become oblate deformed via triaxial, but $N = 166-176$ are oblate deformed and $N = 178-208$ are spherical. It is interesting to note that the shape transition of the highest BE/A nucleus, $^{348}_{138}210$, is quite different than the neighbouring isotopes. The spherical shape obtained at $N = 178$, may be correlated with the proton shell correction exhibiting shallow local minima at $Z = 120$ and 138 and the shell correction for $N = 164$ and 178 , reported by Denisov[3]. At higher temperature $T = 1.5 \text{ MeV}$ and 2.0 MeV , there is no shape dependence with spin and temperature for all isotopes. Hence the limiting temperature for the nuclei with such a heavy mass formed in supernova explosions may be $\approx 1.0 \text{ MeV}$ and which is the ideal temperature in statistical theory to understand the characteristics of such nuclei with very less life time. The E_x Vs N plot (fig.1) for $Z = 132$ and 138 at $T = 1.0 \text{ MeV}$ is plotted and which shows the most probable shape of the isotopes at spin $J = 0\hbar$.

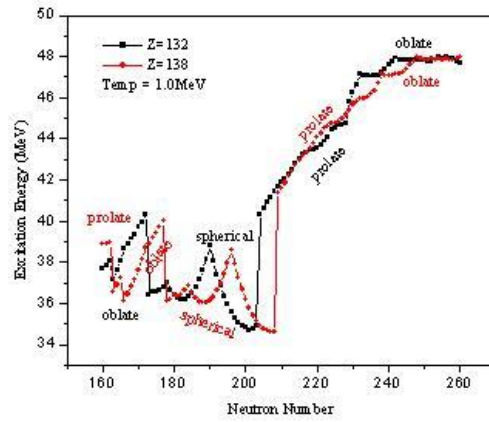


Fig. 1 Probable shape of nuclides at spin $0\hbar$

To find the most probable nuclei with spherical shell closure, the l_{dp} , n_{ld} are calculated and the l_{dp} Vs N plot (fig.2) for different temperature show minima at $N = 166, 174, 200, 220, 234$ and 244 which are very close to the predicted next neutron magic numbers[1]. But at $T > 1.0 \text{ MeV}$ the l_{dp} value is very less and is almost constant for all the isotopes considered. Hence a thermal evolution beyond $T = 1.0 \text{ MeV}$ is hardly possible for these nuclei since at higher temperatures the nucleus may be in plasma state.

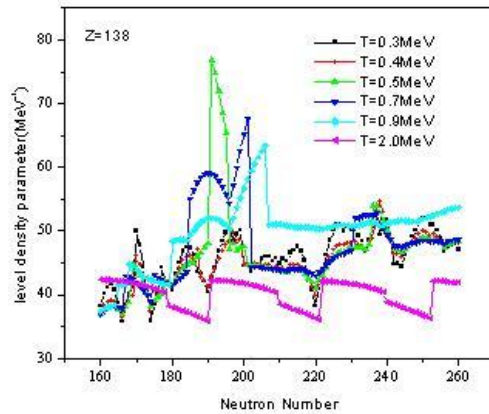


Fig. 2 Variation of 'a' against n^0 addition

References

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