

## Freeze-out Conditions in proton-proton Collisions From SPS to LHC Energies

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### Introduction

Chemical freeze-out dynamics of the particles produced in high-energy heavy-ion collisions (HICs) can be studied from the particle multiplicities. At chemical freeze-out, inelastic collision cease and the particle yields get fixed. The statistical model of non-interacting gas of hadrons and resonances at some volume  $V$ , temperature  $T$  and conserved charge chemical potentials  $\mu_B, \mu_S$  and  $\mu_Q$  corresponding to three conserved charges namely baryon number  $B$ , strangeness  $S$  and electric charge  $Q$  have been remarkably successful in providing a good qualitative description of the mean hadron yields in heavy-ion collisions[1] as well as in p+p collisions[2].

Here we will present the application of thermal model to p+p collisions at  $\sqrt{s_{NN}}=17.3$  (SPS), 200 (RHIC), and to new data from LHC at 900 and 7000 GeV. We will study the systematics of the extracted thermal parameters by employing different ensembles and freeze-out schemes. We have analysed the data in three different ensembles - Grand Canonical Ensemble (GCE), Strangeness Canonical Ensemble (SCE) and Canonical Ensemble (CE). Two different freeze-out schemes have also been studied. One is single freeze-out (1CFO) where all hadrons freeze-out together. The other is double freeze-out (2CFO) where strange and non-strange hadrons freeze-out separately[3]. A comparison with the corresponding results from heavy ion collisions have also been done.

### Thermal Model

In 1CFO scheme, the particle multiplicity within GCE can be written as:

$$N_i = \frac{g_i V}{2\pi^2} \sum_{k=1}^{\infty} (\pm 1)^{k+1} \frac{m_i T^2}{k} K_2 \left( \frac{km_i}{T} \right) \times \exp(\beta k \mu_i) \gamma_S^{k|S_i|} \quad (1)$$

Where  $V$  is the fireball volume,  $g_i$  is the degeneracy,  $m_i$  is the particle mass,  $K_2$  is the second order Bessel function.  $\beta = \frac{1}{T}$ , where  $T$  is the chemical freeze-out temperature. The hadron chemical potential  $\mu_i$  can be written as

$$\mu_i = B_i \mu_B + S_i \mu_S + Q_i \mu_Q \quad (2)$$

It is a standard practice to extract  $\mu_S$  and  $\mu_Q$  from the following constraints: (1) Net  $S = 0$ , (2) Net  $B$ /Net  $Q = 1$ . The remaining parameters ( $V, T, \mu_B$ ) are extracted from fits to hadron yields. The total hadron yield  $N_i^{tot}$  of the  $i$ th hadron include primordial yields  $N_i^{prim}$  and secondary yields, which are the feed-down from decays of heavier resonances and can be written as :

$$N_i^{tot} = N_i^{prim} + \sum_{\text{states } j} N_j^{prim} \text{B.R.}(j \rightarrow i), \quad (3)$$

where B.R. is the branching ratio of  $j$  to  $i$  through all possible channels. We have used the publicly available THERMUS[4] code for 1CFO analysis.

### Results

The upper left plot of Fig. 1 shows chemical freeze-out temperature  $T_{ch}$  as a function

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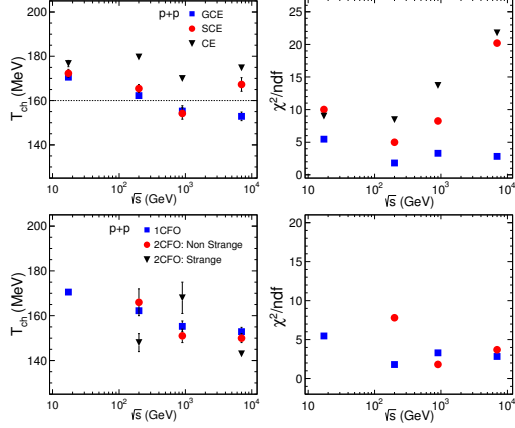


FIG. 1:  $T_{ch}$  and  $\chi^2/ndf$  vs  $\sqrt{s}$  extracted from a statistical fit using particle yields in different ensembles (upper panel) and compared between 1CFO and 2CFO schemes in GCE (lower panel).

of center of mass energy  $\sqrt{s}$  obtained from 3 different ensembles in 1CFO approach. The freeze-out  $T_{ch}$  decreases monotonically from SPS to LHC energies. Also we have observed from the  $\chi^2/ndf$  (upper right plot) of the fits that, GCE describes data best at all the energies with a comparatively lower value of  $\chi^2/ndf$ .

The comparison between temperatures of 2CFO and 1CFO scheme as a function of  $\sqrt{s}$  is shown in the lower left plot of Fig. 1. The 1CFO temperature almost agree with that of non-strange 2CFO temperature. Unlike A+A collisions where 2CFO scheme describes the data better than 1CFO, in p+p the improvement is not significant as the  $\chi^2/ndf$  (lower right plot) is similar in both 1CFO and 2CFO.

Finally, we have also compared the freeze-out parameters extracted in heavy-ion collision with that of p+p in 1CFO as shown in Fig. 2. At lower energies the p+p freeze-out  $T_{ch}$  is higher than A+A, whereas at higher energies it is in agreement with HICs.  $\mu_B$  observed in p+p is similar to that of A+A. Whereas the radius is 5-10 times smaller compared to HICs. Also in contrast to HICs, the value of  $\gamma_S$  is consistently lower in p+p. Which indicates a significant strangeness sup-

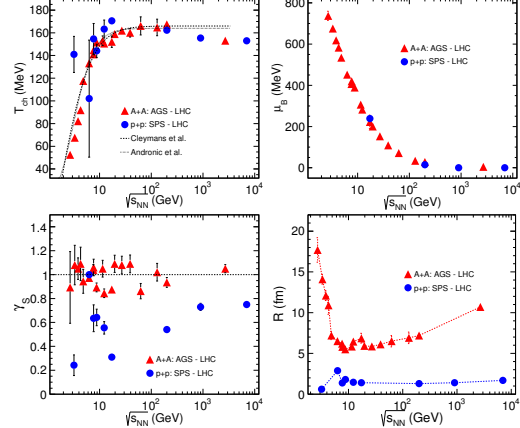


FIG. 2: A compilation of  $T_{ch}$  (top left),  $\mu_B$  (top right),  $\gamma_S$  (bottom left),  $R$  (bottom right) vs  $\sqrt{s_{NN}}$  in p+p collision and compared with A+A collisions in GCE.

pression in p+p as compared to heavy ion even at LHC energies.

## Summary

The chemical freeze-out conditions for p+p collisions have been studied at  $\sqrt{s_{NN}} = 17.3, 200, 900, 7000$  GeV in three different ensembles and two freeze-out schemes. We have also compared the results with that of A+A collisions. The main difference arises in the freeze-out of strange hadrons in p+p collisions with a lower value of strangeness suppression factor as compared to HICs. Which could be due to the expected shorter lifetime of the fireball in p+p collisions.

## Acknowledgments

DM acknowledges the support from DAE and DST/SERB projects of Govt. of India.

## References

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