

High-spin structure in odd- A Hg isotopes

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Introduction

The study of high-spin states in the mass region $A = 190$ -200 is interesting due to the presence of varied structural features [1, 2]. In Hg isotopes, due to their oblate shape near the ground state, the Fermi level moves towards low- Ω $i_{13/2}$ neutron orbitals with increase in mass number. In lower mass odd- A isotopes, the positive-parity yrast states built on the $13/2^+$ spin isomer have been observed beyond the first backbending. Information on the nature of backbending is yet not available in ^{199}Hg , where the levels are known only up to the $I^\pi = 25/2^+$ level.

The nucleus ^{199}Hg may be considered to have a hole coupled to ^{200}Hg and is therefore expected to show similar properties in its excitation spectrum. In ^{200}Hg , there is a report [3] on the influence of $N = 120$ subshell gap on the excitation of rotation-aligned $i_{13/2}$ neutron quasiparticles, which results in a significant gap between the 10^+ and 8^+ levels in comparison to that observed in lower-mass isotopes. This effect also shows up in terms of a higher crossing frequency in this band. It is therefore important to have new information on high-spin states in ^{199}Hg .

Experimental details and data analysis

High spin states in mercury isotopes were populated via multinucleon transfer reactions between ^{197}Au and ^{209}Bi . The target was a gold foil of thickness 50 mg/cm^2 and the beam was ^{209}Bi with an energy of 1450 MeV. The experiment was performed using the Gamma-sphere facility at the Argonne National Laboratory, USA. Further details about the ex-

periment and data analysis may be found in Ref.[4, 5].

Results and Discussion

The previously reported level schemes of ^{197}Hg and ^{199}Hg [6] have been verified and extended. Preliminary findings of this work have been reported earlier [7]. A search was performed for isomers in the nanosecond to sub-microsecond time range however no such states were evident.

A three $i_{13/2}$ neutron quasiparticle (qp) band in ^{199}Hg has been observed for the first time. Spin-parity assignment has been possible up to the $41/2^+$ level. A similar 3-qp band in ^{197}Hg has been extended up to the $49/2^+$ level. As mentioned earlier, based on the expectation of similar properties of excitations in ^{199}Hg and ^{200}Hg , a large gap in the excitation energy is observed between the $33/2^+$ and $25/2^+$ levels in ^{199}Hg . This is due to the effect of the $N = 120$ subshell gap.

A more detailed understanding of the yrast positive-parity bands in Hg isotopes in this mass region can be obtained from the systematic behavior of aligned angular momentum as a function of rotational frequency (Fig. 1). It is evident that the band crossing frequency on an average is higher for the odd- A isotopes ($\hbar\omega = 0.22 \text{ MeV}$) than the even- A isotopes ($\hbar\omega = 0.16 \text{ MeV}$). Within the cranking model, this observation can be explained by the blocking of the AB $i_{13/2}$ neutron crossing in odd- A isotopes. The crossing frequency in both ^{199}Hg and ^{200}Hg is higher than observed for the lower mass isotopes, also a consequence of the $N = 120$ subshell gap.

Cranking calculations have been performed for these nuclei using the Ultimate Cranker code. These calculations provide a satisfactory description of the experimental value of alignment frequency for the BC crossing and angular momentum gain of $10 \hbar$. These cal-

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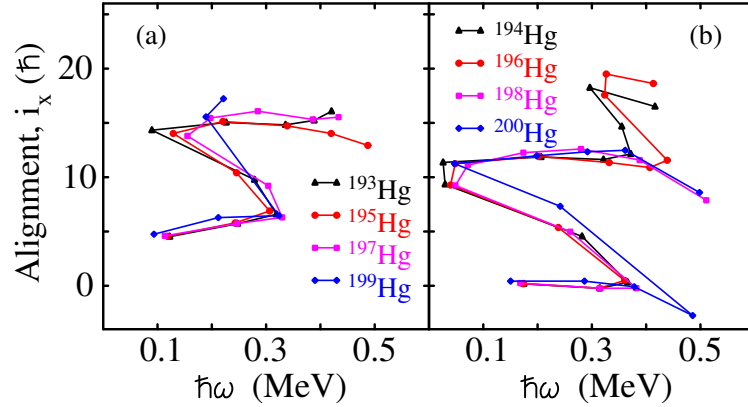


FIG. 1: Aligned angular momentum for the positive-parity bands in (a) $^{196,198,200}\text{Hg}$ and (b) $^{195,197,199}\text{Hg}$.

culations also suggest a considerably higher $h_{11/2}$ proton crossing frequency.

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References

[1] D. Proetal *et. al*, Nucl. Phys. **A226**, 237 (1974).

[2] D. Proetal *et. al*, Nucl. Phys. **A231**, 301 (1974).
 [3] A. G3rgen *et al.*, Eur. Phys. J. A **6**, 141 (1999).
 [4] S. K. Tandel *et. al*, Phys. Lett. B **750**, 225 (2015).
 [5] S. G. Wahid *et. al*, Phys. Rev. C **92**, 054323 (2015).
 [6] D. Mertin *et. al*, Nucl. Phys. **A301**, 365 (1978).
 [7] D. Negi *et. al*, Proc. of the DAE Symp. on Nucl. Phys. **62**, 266 (2017).