

## On nuclear shapes of $^{170}\text{Er}$

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### Introduction

In recent past extremely rich experimental data has come in light in low – lying nuclear spectroscopy. The basic property of nucleus is its geometric shape and it is quantified in terms of geometric deformation parameters  $\beta$  and  $\gamma$ . The possibility of static triaxial shape is a long-standing problem in nuclear structure physics. The  $\gamma$  – unstable and  $\gamma$  – rigid models predict the similar values of energy levels in ground state band but a significant difference is found in the  $\gamma$  – band. The  $\gamma$  – unstable model group the  $\gamma$  – band energy levels as  $2+$ ,  $(3+, 4+)$ ,  $(5+, 6+)$ , ... while  $\gamma$  – rigid model group these energy levels as  $(2_2^+, 3_1^+)$ ,  $(4_2^+, 5_1^+)$ ,  $5_1^+$  ... respectively. The relative displacement of odd spin levels with respect to even spin levels that is odd – even staggering (OES) be taken as a signature of nucleus being  $\gamma$  – soft,  $\gamma$  – rigid or axial.

The staggering indices  $S(I)$  for experimental as well as theoretical energy levels of  $\gamma$  – band is expressed as –

$$S(I) = \frac{(E_I + E_{I-2}) - (2E_{I-1})}{E_{2_1^+}}$$

For axially symmetric rotor,  $S(I)$  does not show any variation in phase and remain small in magnitude. The pattern of  $S(I)$  versus spin  $(I)$  in experiment if found similar in phase with that of  $\gamma$  – rigid model [1] then nucleus is said to be rigid in nature while if the experimental energy staggering pattern is similar to that of  $\gamma$  – soft model [2] in phase, the nucleus is said to be  $\gamma$  – soft.

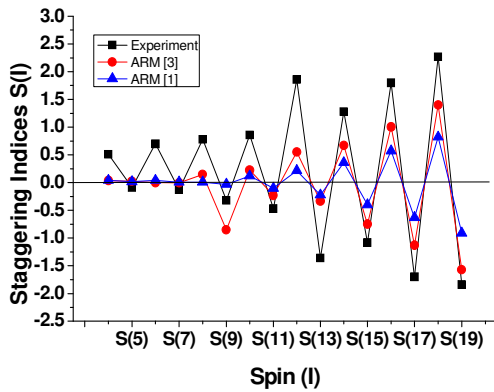
We undertake the study of  $^{170}\text{Er}$  nucleus in the present work. The  $\gamma$  – band energies for this nucleus have appeared in literature showing many high spin states.

**Table – 1**

Experimental staggering indices  $S(I)$  alongwith ARM for  $\gamma = 11.5^\circ$  and  $\gamma = 13^\circ$  calculated from  $E_{2_2^+}/E_{2_1^+}$  [1] and  $B(E2; 2_1^+ \rightarrow 0_1^+)$  [3] values for  $^{170}\text{Er}$  nucleus

S (I)	Exp. Value	ARM [1] $\gamma=11.5^\circ$	ARM [3] $\gamma=13^\circ$
S (4)	0.510	0.041	0.036
S (5)	-0.090	0.017	0.026
S (6)	0.700	0.039	0.0001
S (7)	-0.130	0.008	-0.006
S (8)	0.780	0.009	0.150
S (9)	-0.320	-0.030	-0.850
S (10)	0.860	0.120	0.225
S (11)	-0.470	-0.100	-0.230
S (12)	1.860	0.217	0.550
S (13)	-1.360	-0.224	-0.334
S (14)	1.280	0.360	0.670
S (15)	-1.080	-0.400	-0.75
S (16)	1.800	0.570	1.005
S (17)	-1.700	-0.630	-1.130
S (18)	2.270	0.820	1.400
S(19)	-1.84	-0.910	-1.570

According to ref. 1 violation of axial symmetry of even nuclei the energy states for  $I = 2$  ( $2_1^+$ ,  $2_2^+$ ) one for  $I = 3$  ( $3_1^+$ ) three for  $I = 4$  ( $4_1^+$ ,  $4_2^+$ ,  $4_3^+$ ), two for  $I = 5$  ( $5_1^+$ ,  $5_2^+$ ) etc are generated. Their energies are evaluated using asymmetric rotor relations. The staggering indices  $S(I)$  feeding experimental and asymmetric rotor model (ARM) values are evaluated and are listed in table 1. The values of asymmetric parameter  $\gamma$  are calculated using experimental energy ratio  $E2_2^+/E2_1^+$  [1] and  $B(E2; 2_1^+ \rightarrow 0_1^+)$  [3]. The values of  $S(5)$  and  $S(7)$  in experiment and theory are nearly zero, therefore we can consider all values of  $S(I)$  {table – 1} in same phase for experiment and theory both up to the spin  $I = 19$ . Thus, the nucleus is rigid in nature in beginning and remains as such at higher spin also. This is also reflected from relative magnitude the parameter  $\Delta E1$  and  $\Delta E2$  [4]. The value of asymmetric parameter ( $\gamma$ ) is slightly improved in ref. 3. and these ARM values at higher spin give  $S(I)$  values nearer to that of experiment including similarity in phase of  $S(7)$  in experiment and ARM. The plot between staggering indices and spin is shown in fig. 1.



**Fig. 1**  
Plot between Staggering Indices  $S(I)$  and spin ( $I$ ) in  $\gamma$  - band for  $^{170}\text{Er}$  nucleus

The parameters  $\Delta E1$  and  $\Delta E2$  are defined as –

$$\Delta E1 = E3_1^+ - (E2_1^+ + E2_2^+) \approx 0$$

(for  $\gamma$  – rigid – rotor shape) and

$$\Delta E2 = E3_1^+ - (2E2_1^+ + E4_1^+) \approx 0$$

(for  $\gamma$  – soft shape).

We get these values for  $^{170}\text{Er}$  feeding experimental energies as –

$$\Delta E1 = 2 \text{ eV and } \Delta E2 = 592 \text{ eV.}$$

Thus, our calculations in the framework of rigid triaxial rotor model inferred that  $^{170}\text{Er}$  may be rigid at low spins and remains as such at higher spins also. In support of our results on  $^{170}\text{Er}$  nucleus, it will not be out of place to mention that McCutchen [5] refers to special solutions of the Bohr – Mottelson Hamiltonian that gave predictions for a triaxial structure in respect of some nuclei including  $^{170}\text{Er}$ . In addition, Liao – Ji – Zhi have predicted triaxial nature of fourteen nuclei including  $^{170}\text{Er}$ .

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