

## Heavy quark energy loss and heavy meson spectra in pPb collisions at $\sqrt{s_{NN}} = 5.02$ TeV

Kapil Saraswat<sup>1</sup>, Prashant Shukla<sup>2,3</sup>, and H. C. Chandola<sup>1</sup>

<sup>1</sup> Department of Physics, DSB Campus,

Kumaun University, Nainital - 263001, INDIA.

<sup>2</sup> Nuclear Physics Division, Bhabha Atomic Research Centre, Mumbai - 400085, INDIA. and

<sup>3</sup> Homi Bhabha National Institute, Anushakti Nagar, Mumbai - 400094, INDIA.

### Introduction

At very high temperature or baryon density, QCD predicted a deconfined state of the quarks and gluons, which is known as the Quark Gluon Plasma (QGP). The heavy ion collisions (AuAu, CuCu, XeXe and PbPb) at relativistic energies are performed to create and characterize QGP. The properties of QGP are studied through variety of probes accessible in STAR, PHENIX, CMS, ATLAS and ALICE experiments. Heavy quarks are the best tools to probe the QGP. Due to their large mass, heavy quarks are mostly produced in early phase of collisions when parton momenta are very high before the formation of QGP and the production of heavy quarks is small making them special as a probe for QGP. When heavy quarks travel in the QGP medium, they lose energy either due to collisions with the plasma partons or by radiating gluons. Some models are available to calculate the collisional and radiative energy loss. These energy loss models are used to calculate the nuclear modification factor  $R_{AA}$ .

In this article, we study the transverse momentum ( $p_T$ ) spectra, energy loss and  $R_{AA}$  spectra of the heavy quark in proton - lead (pPb) collisions at  $\sqrt{s_{NN}} = 5.02$  TeV.

### Methodology

The differential cross-section of the heavy mesons in proton - proton (pp) collision is calculated using pQCD approach which is described in the Ref.[1]. In this calculation, we use the LHAPDF6.2.1 (CT14NLO) sets [2] to obtain the parton distribution functions (PDFs). The EPPS16 package [3] is used for the modification of PDFs inside the nucleus.

Peterson fragmentation function is used for the fragmentation of the heavy quarks into the heavy mesons. The generalised dead cone approach (Present) [4] and reactor operator formalism (DGLV) [5] are used to calculate the radiative energy loss while Peigne and Peshier (PP) formalism [6] is used to calculate the collisional energy loss. The energy gain due to the chromoelectric field fluctuation (CMT) is calculated using the Ref.[7]. The path length and initial temperature is obtained using the evolution model [1, 8].

### Results and discussions

Figure 1 shows the pQCD LO calculation of  $d\sigma/dp_T$  of  $B$  mesons as a function of  $p_T$  in pp collision at  $\sqrt{s} = 5.02$  TeV compared with the CMS measurements [9].

Figure 2 shows the energy loss of bottom quark as a function of quark energy in pPb collision at  $\sqrt{s_{NN}} = 5.02$  TeV calculated using PP, DGLV, Present and CMT formalisms. The radiative energy loss calculated by present approach is larger than that by DGLV. The collisional energy loss calculated by PP formalism is less than the radiative energy loss calculation. CMT calculation is very less compared to the other calculations.

Figure 3 shows  $R_{AA}$  of  $B$  mesons as a function of  $p_T$ , obtained by including shadowing energy loss (DGLV, Present, PP+DGLV and PP+Present calculations) in pPb collisions  $\sqrt{s_{NN}} = 5.02$  TeV. The calculations are compared with the CMS data [9]. The radiative energy loss by DGLV formalism describes the CMS data without adding collisional energy loss at high  $p_T$ . The other calculations such as PP, Present, PP + Present and PP + DGLV are below the CMS measurements.

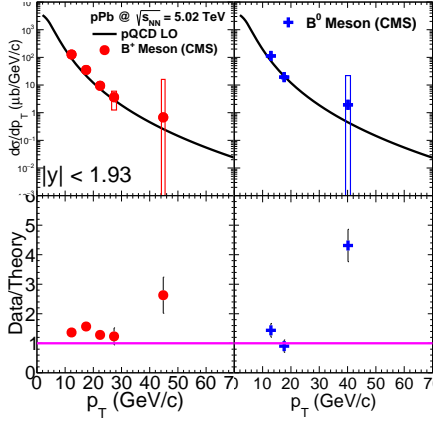
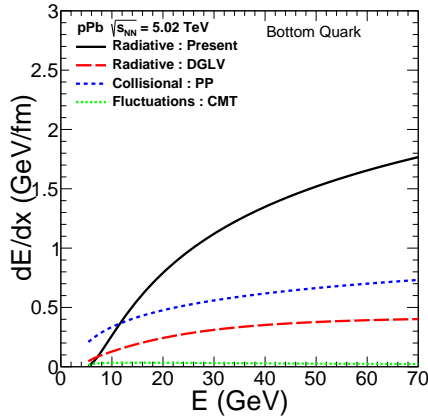
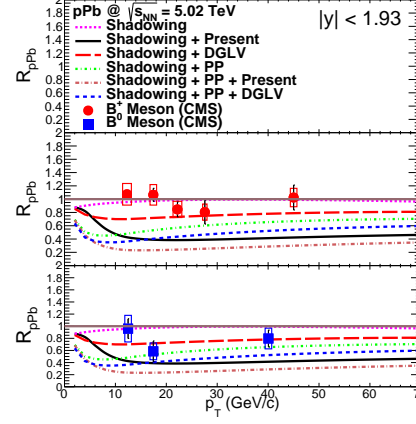
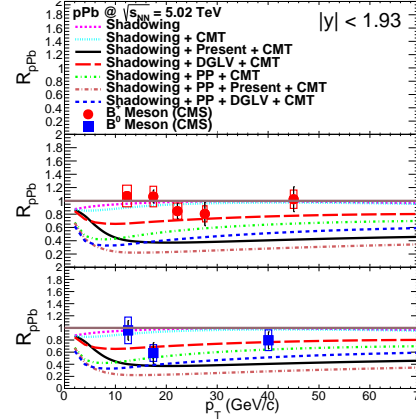

 FIG. 1:  $d\sigma/dp_T$  of  $B$  as a function of  $p_T$ .

 FIG. 2:  $dE/dx$  of Bottom as a function of  $E_T$ .

 FIG. 3:  $R_{AA}$  of  $B$  as a function of  $p_T$ .

 FIG. 4:  $R_{AA}$  of  $B$  as a function of  $p_T$ .

Figure 4 is same as Figure 3 but it includes the CMT formalism.

## Conclusion

We find that the radiative energy loss from DGLV formalism alone is sufficient to produce  $B$  meson  $R_{AA}$  at high  $p_T$ . The other calculations with or without including CMT are below the CMS measurements. The study of the energy loss and  $R_{AA}$  of  $D$  meson in pPb collisions is under process.

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