

## Low-lying nuclear structure of stable even-A Zn-isotopes

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### Introduction

From last few decades, significant efforts were made to study the evolution of nuclear shell structure, ranging from stable nuclei to the exotic nuclei. Experiments performed at the newly developed radioactive ion beam facilities have indicated that some of the traditional magic numbers should be replaced with new magic numbers. Zn-isotopes are of particular interest, being two protons away from the  $Z = 28$  magic number and near neutron midshell ( $N=28-50$ ) along with a harmonic subshell closure for  $^{70}\text{Zn}$  at  $N = 40$ . Generally, the collective vibrational motion is attributed to the structure of even-A Zn-isotopes on the basis of typical energy pattern of first few excited states viz. the two-phonon triplet states ( $0_2^+$ ,  $2_2^+$ ,  $4_1^+$ ) at about twice the energy of first excited  $2_1^+$  state. However such a general picture fails while considering other electromagnetic properties of Zn-isotopes [1]. Moreover the excitation energy of  $0_2^+$  in these isotopes is also surprising as it follows altogether a different behaviour as compared with the other low-lying excited states. A myriad of nuclear phenomenon like shape coexistence, shape transition, etc. have been observed and predicted in the neighbouring nuclei of Zn-isotopic chain. While no clear evidences of such phenomenon are experimentally observed in Zn-isotopes, a number of predictions are made theoretically and few hints available from some of the previous experiments.

### Shell model calculations

Shell model calculations have been performed in order to understand the low-

lying electromagnetic nuclear structure of Zn-isotopes. The calculations were taken up using KSHELL code of Shimizu et al.[2], and the available high performance computing facility at IUAC New Delhi.

The single particle space was provided by the  $f_5pg_9$  model space viz.  $1p_{3/2}$ ,  $0f_{5/2}$ ,  $1p_{1/2}$ ,  $0g_{9/2}$  and the matrix elements of the interaction, by using two different interaction hamiltonian viz. JUN45 and jj44b, in pn format. An inert  $^{56}\text{Ni}$  core was considered. In JUN45, the initial Hamiltonian is based on an effective interaction derived from the realistic Bonn-C potential in a microscopic way. It was modified iteratively using a set of 400 experimental binding and excitation energy data from 69 nuclei in the mass region  $A=63-96$  on 45 well-determined linear-combinations of the 133 TBME and four SPE. Single-particle energies (SPE) are taken to be -9.8280, -8.7087, -7.8388, and -6.2617 MeV for the  $p_{3/2}$ ,  $f_{5/2}$ ,  $p_{1/2}$  and  $g_{9/2}$  orbits, respectively. Similarly the jj44b interaction was also derived from a realistic interaction based on the Bonn-C potential which has been developed by fitting 600 binding energies and excitation energies of nuclei with  $Z = 28-30$  and  $N = 48-50$ . The single particle energies are taken to be -9.6566, -9.2859, -8.2695 and -5.8944 MeV for the  $p_{3/2}$ ,  $f_{5/2}$ ,  $p_{1/2}$  and  $g_{9/2}$  orbits, respectively.

### Results and discussions

#### 1. Excitation energy

The excitation energy spectrum is shown in Fig.1.

#### 2. Reduced transition strength

The reduced transition strengths for even-even  $^{62}\text{Zn} - ^{70}\text{Zn}$  are calculated using both the JUN45 and jj44b effective interactions and employing an effective charge of  $e_\pi = 1.8e$ ,  $e_\nu = 0.8e$  [3]. Fig. 2 shows the reduced electric quadrupole transition

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