

Shape evolution of nuclei in $A \sim 55$ region

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Introduction

The nuclear structure of $A \sim 55$ region are governed by $1f_{7/2}$ and intruder $g_{9/2}$ orbitals across $N = Z = 28$ shell closure. The nature of this shell closure, which is the first shell closure originated due to the spin-orbit coupling term, is in debate from various theoretical and experimental findings. Further, the presence of shape driving orbitals like $1f_{7/2}$ and $1g_{9/2}$ can induce deformation in the system. There are limited informations regarding the observation of different nuclear shapes in this mass region with nuclei around doubly magic ^{56}Ni ($Z = 28, N = 28$).

In this work, we have calculated the deformations for even-even nuclei from $Z = 22$ (Ti) to $Z = 28$ (Ni) for different N to get an idea regarding the evolution of nuclear shapes with change in proton and neutron Fermi surfaces in this mass region. The transition rate and quadrupole moment, $B(E2)$ and Q_t , have been deduced from the deformation parameters, obtained from the calculation, and compared with the experimental ones. The experimental $B(E2)$ and Q_t values were deduced from the measured lifetimes (τ) reported for the 1^{st} (2^+) and using the equations given in [1]. These calculations have been performed for lower rotational frequencies corresponding to spin value of $2\hbar$.

Method of the calculation

The shape of a nucleus has been obtained from the Total Routhian Surface (TRS) calculations. In this method, the potential energies in the body-fixed frame are calculated in the Nilsson-Strutinsky formalism [2]. A deformed Woods-Saxon potential with universal param-

eters is used to calculate the single particle energies. The nuclear deformation is defined by the β_2 , β_4 and γ parameters. In the Lund convention, used here, $\gamma = 0^\circ$ (-60°) is for axially deformed prolate (oblate) shape and any other value of γ indicate triaxial shape. The detailed technical procedure has been given in [3, 4]. The total Routhians are obtained at different values of β_2, β_4, γ and are plotted in the β_2, γ surface after minimizing on β_4 . The calculations are done for different quasi-particle configurations and at several rotational frequencies ($\hbar\omega$). The shape of a nucleus for a particular $\hbar\omega$ at a particular configuration is obtained from the values of β_2 and γ corresponding to the minimum of the TRS at that $\hbar\omega$ for that configuration.

Results and Discussions

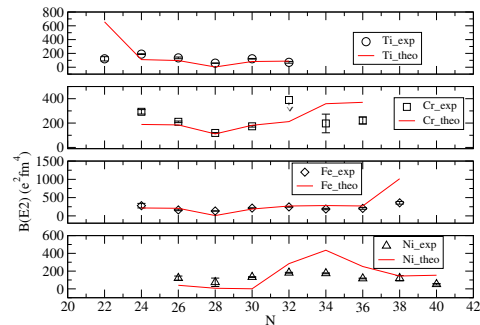


FIG. 1: The experimental values (symbols) of $B(E2)$ for the Ti, Cr, Fe and Ni isotopes are compared with those obtained from the β_2 and γ deformation parameters obtained from the TRS calculations (solid line) in this work.

The theoretical results are summarized in Table I. It contains the deformation parameters β_2 and γ , obtained from the present calculation, for various nuclei with even Z from Ti ($Z = 22$) isotopes to Ni ($Z = 28$)

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TABLE I: Deformation parameters (β_2 and γ) of the nuclei near the ground state obtained in this work.

| Z | N | Nucleus | β_2 | γ | $B(E2)(e^2fm^4)$ |
|----|----|-----------|-----------|----------|------------------|
| 22 | 22 | ^{44}Ti | 0.4 | -21.8 | 656.83 |
| 22 | 24 | ^{46}Ti | 0.19 | -2.7 | 110.55 |
| 22 | 26 | ^{48}Ti | 0.17 | -5.6 | 97.06 |
| 22 | 28 | ^{50}Ti | 0.03 | -30 | 3.49 |
| 22 | 30 | ^{52}Ti | 0.15 | -5.4 | 82.67 |
| 22 | 32 | ^{54}Ti | 0.14 | -21 | 88.76 |
| 24 | 24 | ^{48}Cr | 0.22 | -2.2 | 188.75 |
| 24 | 26 | ^{50}Cr | 0.21 | -3.6 | 184.96 |
| 24 | 28 | ^{52}Cr | 0.15 | -14.8 | 110.83 |
| 24 | 30 | ^{54}Cr | 0.2 | -2.8 | 182.07 |
| 24 | 32 | ^{56}Cr | 0.2 | -9.6 | 212.24 |
| 24 | 34 | ^{58}Cr | 0.26 | -3.9 | 359.08 |
| 24 | 36 | ^{60}Cr | 0.26 | -3.1 | 370.49 |
| 26 | 24 | ^{50}Fe | 0.21 | -3.4 | 216.32 |
| 26 | 26 | ^{52}Fe | 0.2 | -4 | 207.50 |
| 26 | 28 | ^{54}Fe | 0.04 | -18.6 | 9.30 |
| 26 | 30 | ^{56}Fe | 0.18 | -5.4 | 187.32 |
| 26 | 32 | ^{58}Fe | 0.2 | -12.4 | 269.95 |
| 26 | 34 | ^{60}Fe | 0.21 | -4.8 | 282.47 |
| 26 | 36 | ^{62}Fe | 0.2 | -6.1 | 271.44 |
| 26 | 38 | ^{64}Fe | 0.38 | -0.7 | 1016.86 |
| 28 | 26 | ^{54}Ni | 0.08 | -8.7 | 40.07 |
| 28 | 28 | ^{56}Ni | 0.03 | -30 | 6.58 |
| 28 | 30 | ^{58}Ni | 0.01 | -23.8 | 0.75 |
| 28 | 32 | ^{60}Ni | 0.18 | -36.8 | 284.07 |
| 28 | 34 | ^{62}Ni | 0.23 | -52.3 | 435.01 |
| 28 | 36 | ^{64}Ni | 0.17 | -47.9 | 251.91 |
| 28 | 38 | ^{66}Ni | 0.13 | -50.4 | 144.88 |
| 28 | 40 | ^{68}Ni | 0.18 | -77.8 | 153.61 |

isotopes. These shapes are calculated near the ground state at lower $\hbar\omega$.

The calculated values of the deformation parameters are used to calculate the other parameters like the quadrupole moment Q_t and the transition strengths $B(E2)$. The experimental values of these parameters are extracted from the measured lifetimes (τ) reported in the ENSDF [5]. The expressions used for determining these quantities in both the cases are taken from reference [1]. These

measured quantities are for the $1^{st} 2^+$ states in all these cases. The experimental and theoretical values of Q_t and $B(E2)$ are compared in Fig. 1. It can be seen that the experimental values are well reproduced by the calculations.

Summary and Conclusion

The calculated results predict different shapes for different N/Z. At $N = 28$, which is the magic number, in all the isotopes the deformation seems to decrease and shape tends towards triaxiality. For Ti and Cr isotopes, the shapes are mostly prolate except $N = 28$. In Fe isotope, $N=28$ and 32 are triaxial while other N are prolate in shape. Thus the evolution of shapes from prolate to triaxial and again to prolate are predicted for these isotopes from these calculations. The Ni isotope ($Z = 28$) is predicted to have evolution of shape from prolate to oblate via triaxial shape with increase in N. The experimental $B(E2)$ values are well reproduced by the theoretical calculations which validate the present calculations. Experimental observations of this shape evolution in these nuclei will be interesting.

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