

Novel Insights into Deformed Nuclear Structures via γ -Band Energy Variations

Pooja Sharma¹, Amit Bindra^{1*} and Mudasir Ahmad^{1#}

¹Department of Physics, School of Chemical Engineering and Physical Science, Lovely Professional University, Jalandhar-144411, INDIA

Introduction

In this study, the link between rotational momentum ($J=6$) and transition probability ($B(E2)$) is analyzed using empirical data to discover a first-order phase change in nuclei, indicating a quantum shift from spherical to deformed nuclear structure. The shift at $N \sim 90$ is explained, highlighting the groundbreaking investigation of staggering parameters in $B(E2)$. In this article, the term "axially rigid region" refers to the region between the vibrator and the axially symmetric rotor where γ remains rigid, axially aligned with a harmonic oscillator potential, and the lowest in gamma is very close to zero. As part of the Interacting Boson Model (IBA), this region shows the change from the $U(5)$ symmetry to the $SU(3)$ symmetry. This is called a first-order phase transition [2]. The critical point symmetry $X(5)$ has also been linked to this location [3]. In the deformed rare earth region, this work intends to introduce and examine a novel signature by comparing it to an extensive dataset. In the end, we will show that studying energy staggering along an isotopic chain can be a quick and easy way to figure out how nuclear structure has changed over time and maybe even give us clues about how phase transitions happen. Researchers Kumari and Mittal [4] studied the Grodzins product and found a strong association with $B(E2)$. Bindra and Mittal [5] also looked at the ESF and ROTe, both multiplied by $B(E2)$ values, in the context of shape-transitional isotopes and noted several noticeable anomalies in their study. Many nuclear data points have been obtained through experimental study, and theoretical calculations have been revised as a result of contemporary technology developments. Changes in $S(J)$ characteristics with $B(E2)$ and $R_{4/2}$ within the mass range of Gd-W nuclei are the focus of our study. The region between 150 and 180 mass is our main focus since it shows a lot of different features.

Formalism

For γ -bands, we will use the amount [1] to investigate the odd-even staggering:

$$S(J) = \frac{\{E(J^+_{\gamma}) - E[(J-1)^+_{\gamma}]\} - \{E[(J-1)^+_{\gamma}] - E[(J-2)^+_{\gamma}]\}}{E(2^+_{\gamma})} \quad (1)$$

This amount, normalized to the energy of the ground band's initial excited state, 2^+_{γ} , quantifies the displacement of the $(J-1)^+_{\gamma}$ level with respect to the average energy of its surrounding levels, J^+_{γ} and $(J-2)^+_{\gamma}$. $S(J)$ is highly adaptable to changes in configuration since it takes a confound derivative form. The γ -band levels are shown in Figure 1 together with the associated $S(6)$ values.

Results and Discussion

- **γ -Band Energy Staggering, $S(6)$, in relation with Transition Probability from Gd-W Isotopes**

It is useful to investigate the evolution of a particular characteristic, $S(6)$, in response to structural modifications after examining the variation of $S(J)$ with regard to spin for various nuclear structures. We will now turn our attention to $S(6)$, which is the displacement of the 3^+_{γ} state with respect to the average energy of its adjacent states, 2^+_{γ} and 4^+_{γ} . The 2^+_{γ} state's energy is used to standardize this quantity. By means of transition probabilities, we examine the changes in $S(6)$ for discrete structures. Figures 1 (left) and 1 (right) show how the staggering of Gd-W isotopes differs. $S(6)$ and $B(E2)$ show comparatively stable and minimal values for the heavier isotopes, ^{152}Gd , ^{156}Gd , ^{154}Dy , and ^{160}Er . But in $S(6)$, every single one of them exhibits a more noticeable distorted pattern. It is a number greater than 82 ($N > 82$). In fig. 1 (left) at $N = 90$, ^{64}Gd shows a sudden kink, followed by ^{66}Dy , which then exhibits that the heavy rare-earth nuclei display axially symmetric properties, particularly those with neutrons, with a gradually declining trend beyond $N \geq 90$. This indicates a shift from vibrator to rotor nuclei, serving as strong evidence for the abrupt start of deformation, consistent with the findings of Caikrili and Casten [6]. And also, we can say that $N = 90$ is the point where, as shown by the Gd nuclei and the first-order phase transition after $N = 92$ (at $N = 94$), Since this study's absolute $B(E2)$ values cover both excitation

energies and relative and absolute transition probabilities, they allow for a more thorough analysis of the X(5) model predictions. Also mentioned as potential possibilities for exhibiting the traits of the X(5) critical point symmetry are ^{154}Gd [7] and ^{156}Dy [8].

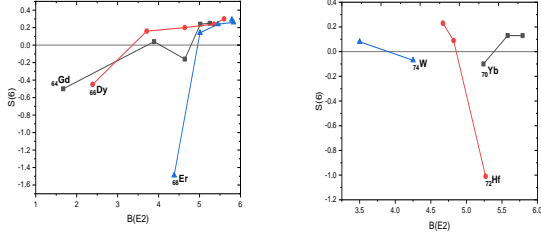


Figure 1: A comparative analysis of $S(6)_\gamma$, in relation with $B(E2)$ for Gd-W isotopes.

While Fig. 1 (right) comprises isotopes from Yb-W, The findings for Yb-W exhibit notable distinctions compared to the previous calculations that assigned less significance to the more elevated excited levels, and they become more pronounced when examining the Yb-W nuclides. As the value of Z rises and the atomic nuclei become increasingly gamma-soft, the smoother increase in $S(6)$ and that corresponds to a decrease in $B(E2)$ in the case of Hf and W nuclides indicates that the shape transition is less sudden in its manifestation and a transformation from a globular to a distorted configuration occurs abruptly.

Table 1: Transitional isotopes within the rare earth along with Empirical $S(6)_\gamma$ values and $B(E2)$ for Gd-W isotopes.

Nuclei	$S(6)$	$B(E2)$
^{152}Gd	-0.5	1.67
^{154}Gd	0.04	3.89
^{156}Gd	-0.16	4.64
^{158}Gd	0.24	5.02
^{160}Gd	0.25	5.25
^{154}Dy	-0.45	2.39
^{156}Dy	0.16	3.71
^{158}Dy	0.20	4.66
^{162}Dy	0.24	5.35
^{164}Dy	0.3	5.60
^{160}Er	-1.49	4.38
^{162}Er	0.14	5.01
^{164}Er	0.24	5.45
^{166}Er	0.26	5.83
^{168}Er	0.29	5.79
^{166}Yb	-0.10	5.24
^{168}Yb	0.13	5.58
^{170}Yb	0.13	5.79

^{176}Hf	-1.01	5.27
^{178}Hf	0.09	4.82
^{180}Hf	0.23	4.67
^{180}W	-0.07	4.25
^{186}W	0.08	3.50

Conclusion

To sum up, the exploration of the development of the $S(6)$ feature in response to structural alterations offers important information about the behaviour of nuclear structures, especially when looking at the staggering patterns in Gd-W isotopes. Heavier isotopes such as $^{152,156}\text{Gd}$, ^{154}Dy , and ^{160}Er show stable and minimal values for $S(6)$ and $B(E2)$, with $S(6)$ showing a more prominent distorted pattern when the neutron number goes beyond 82 ($N > 82$). The rapid kink at $N = 90$ in ^{64}Gd , and the equivalent behaviour in ^{66}Dy , highlight the onset of deformation and confirm the observations of Caikrili and Casten. This implies a shift from vibrator to rotor nuclei. The abrupt start of deformation, as predicted by the X(5) model, is consistent with this transition, which becomes especially noticeable at approximately $N = 90$. The analysis also indicates that ^{154}Gd and ^{156}Dy might show the X(5) critical point symmetry. A less abrupt shape transition towards atomic nuclei that are becoming more gamma-soft is suggested by the opposing results in Yb-W isotopes, where $S(6)$ values climb slowly with a commensurate reduction in $B(E2)$.

References

- [1] N. V. Zamfir and R. F. Casten, *Phys. Lett. B* 260, 265 (1991).
- [2] F. Iachello, *Phys. Rev. Lett.* 87, 052502 (2001).
- [3] R. F. Casten and N. V. Zamfir, *Phys. Rev. Lett.* 85, 3584 (2000).
- [4] P. Kumari and H. M. Mittal, *Open Phys.* 13, 305 (2015)
- [5] A. Bindra and H. M. Mittal, *Nuclear Physics A*, 975: 48–58 (2018)
- [6] R. B. Cakirli and R. F. Casten, *Physical Review Letters*, **96**, 132501 (2006),
- [7] F. Iachello and A. Arima, *The Interacting Boson Model* (Cambridge University Press, Cambridge, 1987).
- [8] G. Gneuss and W. Greiner, *Nucl. Phys.* A171, 440 (1971).