

# Cross Section Measurement of $^{nat}\text{Ni}(\alpha, \text{xp})^{61}\text{Cu}$ Reaction: Implications for Reactor Safety and Secondary Particle Production

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## Introduction

The radiation-induced transmutation reactions, hydrogen and helium bubble formation in the reactor structural materials must be analysed for understanding the structural stability of reactors. Stainless steels with Fe, Ni and Cr as main constituents (SS316 composed of Fe~65%, Ni~14%, Cr~17%) are frequently used as reactor structural materials. Moreover, a class of nickel-based alloys has become prominent in Gen IV reactor systems due to their outstanding temperature sensitivity and corrosion resistance [1].

In fusion reactor environment there is high fluence neutrons, and their interaction with structural material will cause (n,p) and (n, $\alpha$ ) channels [2]. In fission reactors, there will be a production of alpha as a result of ternary fission. The reactor structure will interact with energetic particles inside the reactor, resulting in the production of secondary charged particles and will lead to degradation of structural stability.

The experimental cross section data available for  $^{nat}\text{Ni}(\alpha, \text{xp})$  in the IAEA-EXFOR [3] library is compared with the cross sections available in evaluated nuclear data library TENDL-2023, as shown in FIG. 1. Literature data shows significant cross section discrepancies across energy ranges, despite using the same stacked foil activation method. As a result, we decided to investigate the alpha-

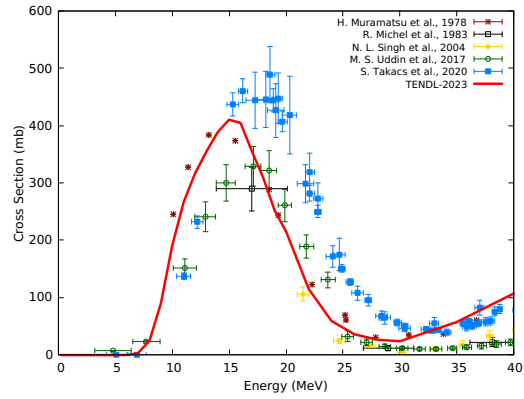


FIG. 1: Comparison of data from evaluated nuclear data library-TENDL-2023, with Previous experimental data sets taken from EXFOR.

induced reaction on natural nickel, focussing primarily on the residue  $^{61}\text{Cu}$ , because its production leads to an increase in proton yield.

## Experimental Setup and Data Analysis

The experiment was performed at the K-130 Cyclotron facility, VECC, Kolkata, India. Natural nickel foils ( $\sim 45 \text{ mg/cm}^2$ ) were prepared by rolling. Two separate stacks were irradiated with 30 MeV and 40 MeV alpha beams at a current of  $\sim 100 \text{ nA}$ . The stopping ranges of alpha particles through the stacks were simulated using the SRIM code for both energy configurations.

After irradiation, the evaporation residues formed in the  $\alpha + ^{nat}\text{Ni}$  were identified by all the characteristic  $\gamma$ -rays and the gamma-ray

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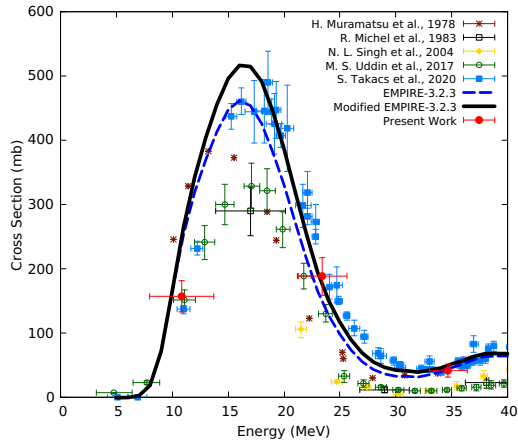


FIG. 2: Comparison of experimental cross sections with EMPIRE-3.2.3 predictions, including modifications accounting for the production cross section of  $^{61}\text{Zn}$ .

activity of the residue was investigated using the HPGe detector (CANBERRA). The energy and efficiency calibration of the detector were done using a standard  $^{152}\text{Eu}$  source. The spectrum was analysed using IUAC data acquisition and analysis software CANDLE.

The cross section calculation was done using the activation equation.

$$\sigma(E) = \frac{C_\gamma \lambda e^{\lambda t_{cool}}}{N \phi I_\gamma \epsilon_\gamma (1 - e^{-\lambda t_{irrad}}) (1 - e^{-\lambda t_{count}})} \quad (1)$$

where  $C_\gamma$  is the net counts under the photopeak of the radionuclide,  $\lambda$  is decay constant of the radioisotope (in  $s^{-1}$ ),  $N$  is the number of target atoms per  $cm^2$ ,  $\phi$  is the beam flux (particles/s),  $\epsilon_\gamma$  is the absolute efficiency of the detector at the measured gamma ray,  $I_\gamma$  is the absolute intensity of the measured gamma ray,  $t_{irrad}$ ,  $t_{cool}$  and  $t_{count}$  are the irradiation time, cooling time and counting time (in s) respectively.

## Results and Discussion

The experimentally obtained cross section presented here agrees well with the experimental data available in IAEA EXFOR [3] (FIG. 2). The nuclear reaction code EMPIRE-3.2.3 is used to optimize experimental data. Different level density models available in EMPIRE-3.2.3 have been used in simulations, out of which the Gilbert-Cameron nuclear level density model is found to be reproducing the experimental data up to some extent. The production of  $^{61}\text{Cu}$  involves both direct and indirect pathways. The direct pathway is through the  $^{nat}\text{Ni}(\alpha, xp)$  reaction. Indirectly,  $^{61}\text{Cu}$  is generated via  $\beta^+$  decay of  $^{61}\text{Zn}$ , which itself is formed predominantly from  $^{58}\text{Ni}(\alpha, n)$  (Q-Value = -9.53 MeV) and  $^{60}\text{Ni}(\alpha, 3n)$  (Q-Value = -29.91 MeV) reactions. Incorporating the production probability of  $^{61}\text{Zn}$  into computational models with optimized parameters improves their agreement with experimental data to some extent, as illustrated in FIG. 2.

## Acknowledgments

The authors gratefully acknowledge financial support from DAE-BRNS (Sanction Order 36(6)/14/30/2017-BRNS/36204) and UGC-DAE CSR. Swapna Balakrishnan thanks IUAC, Delhi, for financial assistance. We also thank Dr. Chandana Bhattacharya, A.A. Mallick, and the operating staff of the K-130 Cyclotron facility at VECC, Kolkata, for their support during the experiment.

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