

# Radiative transition of $\psi(2S) \rightarrow \chi_{c0}(1P) + \gamma$ in light-front quark model

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## 1. Introduction

Radiative transitions offer a clear probe of internal structure of quark-antiquark bound states via electromagnetic probes, with varied resolution. These transitions are vulnerable to relativistic effects, highlighting recent differences between NRQCD and experimental data [1]. Radiative transition between  $1^{--}$  (vector) and  $0^{++}$  (scalar) mesons via the emission of a photon is characterized as the electromagnetic transition (E1) transition. In this work, we have calculated the decay width and transition form factors (TFFs) of the  $\psi(2S) \rightarrow \chi_{c0}(1P) + \gamma$  E1 transition using light-front quark model (LFQM).

The paper is organized as follows. In Sec. 2, we have discussed the LFQM. The TFFs and decay width for the transition of  $\psi(2S) \rightarrow \chi_{c0}(1P) + \gamma$  have been explained in Sec. 3. The results of the present work has been explained in Sec. 4.

## 2. Light-front quark model

The minimal hadron Fock-state wave function based on light-front quantization in the form of quark-antiquark is given by [2]

$$|M^{nS(nP)}(P, J, J_z)\rangle = \sum_{\lambda_q, \lambda_{\bar{q}}} \int \frac{dx d^2\mathbf{k}_\perp}{\sqrt{x(1-x)}16\pi^3} \psi_{\lambda_q, \lambda_{\bar{q}}}^{JJ_z}(x, \mathbf{k}_\perp) |x, \mathbf{k}_\perp\rangle,$$

where  $|M^{nS(nP)}(P, J, J_z)\rangle$  is the meson eigenstate with  $P = (P^+, P^-, P_\perp)$  being the four-vector total momentum of the hadron and  $(J, J_z)$  is the total angular momentum.  $nS(nP)$  stands for S-wave and  $P$ -wave,

$\mathbf{k} = (k^+, k^-, \mathbf{k}_\perp)$  is the momentum of the active quark and  $x = \frac{k^+}{P^+}$  is the longitudinal momentum fraction carried by the active quark. Here  $\lambda_{q(\bar{q})}$  is the helicity of quark ( $q$ ) and anti-quark ( $\bar{q}$ ).

The light-front wave function  $\psi_{\lambda_q, \lambda_{\bar{q}}}^{JJ_z}$  for the pseudo-scalar and vector mesons is defined as follows

$$\psi_{\lambda_q, \lambda_{\bar{q}}}^{JJ_z}(x, \mathbf{k}_\perp) = \phi_{nS(nP)}(x, \mathbf{k}_\perp) \mathcal{X}_{\lambda_q, \lambda_{\bar{q}}}^{JJ_z}(x, \mathbf{k}_\perp). \quad (1)$$

where  $\phi_{nS(nP)}(x, \mathbf{k}_\perp)$  is the radial (orbital) wave function and  $\mathcal{X}_{\lambda_q, \lambda_{\bar{q}}}^{JJ_z}$  is the spin wave functions [3]. The  $1P$ -wave radial wave function is taken as

$$\phi^{1P}(x, \mathbf{k}_\perp) = 4\sqrt{2} \left( \frac{\pi}{\beta^2} \right)^{3/4} \sqrt{\frac{dk_z}{dx} \frac{k_z}{\beta}} \exp\left(-\frac{k_z^2 + k_\perp^2}{2\beta^2}\right),$$

Similarly, for 2S radial wave form is given as

$$\phi^{2S}(x, \mathbf{k}_\perp) = 4 \left( \frac{\pi}{\beta^2} \right)^{3/4} \sqrt{\frac{\partial k_z}{\partial x}} \exp\left(-\frac{2^\delta k_z^2 + k_\perp^2}{2\beta^2}\right) \left(a'_2 - b'_2 \frac{k_z^2 + k_\perp^2}{\beta^2}\right).$$

Here  $k_z = (x - 1/2)M$  with bound state mass  $M = \sqrt{\frac{k^2 + m_c^2}{x(1-x)}}$  for charmonia ( $c\bar{c}$ ).  $\beta$  and  $m_c$  are the harmonic scale operator and quark mass, having a value of 0.566 GeV and 1.4 GeV respectively [4].

## 3. Transition form factors

Transition form factor  $E1(Q^2)$  for electric dipole decay of scalar meson  $S \rightarrow V\gamma^*$  ( $\gamma^*$  is virtual photon) is defined as

$$\langle M^{(1P)}(P', J, J_z) | J_{\text{em}}^+ | M^{2S}(P, J, J_z) \rangle = ie\epsilon^{\mu\nu\rho\sigma} \epsilon_\nu(P, h) q_\rho P_\sigma E1(Q^2),$$

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our work	BLFQ [1]	MRK1'76	Dudek'09 [5]
25.353	21.953 ±5.017	22.035 ±7.897	25.902 ±11.105
PDG [6]	CBAL'86	CLEO'04 [8]	BESIII'17[7]
28.781 ±0.987	29.11 ±2.879	27.053 ±1.645	27.547 ±1.233

TABLE I: Decay width of our LFQM result and other experimental data [1, 5–8] in unit of KeV

where  $q = P - P'$  is the four momentum of the virtual photon,  $\epsilon_\nu(P, h)$  is the polarization vector of the photon and helicity  $h$ .  $J^+$  is the longitudinal charge density.

The decay width is found to be

$$\Gamma(S \rightarrow V\gamma) = \frac{\alpha}{3} E1^2(Q^2 = 0) k_\gamma^3,$$

where  $\alpha$  is the strong coupling constant and  $k_\gamma = (M_S^2 - M_V^2)/2M_S$ .  $M_S$  and  $M_V$  are masses of  $\psi(2S)$  and  $\chi_{c0}(1P)$  respectively.

#### 4. Results and Discussions

For numerical predictions of transition predictions, we need only two input parameters  $\beta$  and  $m_q$ , which are taken from Ref. [4]. The  $E1(Q^2)$  TFFs for  $\psi(2S) \rightarrow \chi_{c0}(1P) + \gamma$  decay has been shown in Fig. 1 along with comparison with basic light-front quantization (BLFQ) [1] and lattice simulations results [5]. At  $Q^2 = 0$ ,  $E1(Q^2)$  is found to be 1.3  $\text{GeV}^2$ , which is nearly equal to the particle data group (PDG) data [6], where as in case of BLFQ it was found to be 1  $\text{GeV}^2$ . The behavior of  $E1(Q^2)$  is found to be in good agreement with BLFQ result along with lattice simulations. The  $E1(Q^2)$  decreases up to a energy scale  $Q^2 = 1.8 \text{ GeV}^2$ , then increases with increase in  $Q^2$ .

The decay width  $\psi(2S) \rightarrow \chi_{c0}(1P)$  of our calculations is found to be 25.353 KeV, which has been shown in Table I and compared with experimental results [6, 7] and lattice simulation results [1, 5, 6, 8]. Overall the LFQM results for the radiative decay  $\psi(2S) \rightarrow \chi_{c0}(1P)$

is in good agreement with other model predictions, lattice simulations and experimental results.

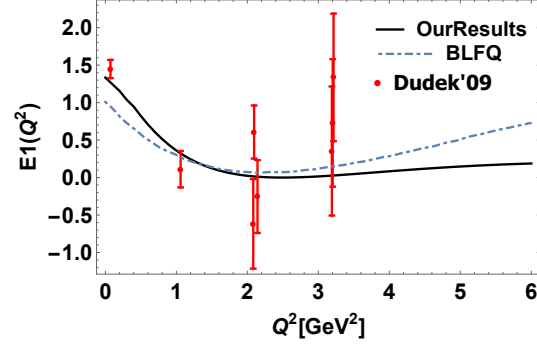


FIG. 1: Comparison of E1 radiative transition form factor of our work (LFQM) with BLFQ [1] and other lattice simulations [5].

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